Optimal control of stochastic phase-field models related to tumor growth *

CARLO ORRIERI⁽¹⁾ E-mail: carlo.orrieri@unitn.it

ELISABETTA ROCCA⁽²⁾ E-mail: elisabetta.rocca@unipv.it

LUCA SCARPA⁽³⁾ E-mail: luca.scarpa@univie.ac.at

(1) Department of Mathematics, University of Trento Via Sommarive 14, 38123 Povo (Trento), Italy

(2) Department of Mathematics, University of Pavia, and IMATI - C.N.R. Via Ferrata 5, 27100 Pavia, Italy

> ⁽³⁾ Faculty of Mathematics, University of Vienna Oskar-Morgenstern-Platz 1, 1090 Vienna, Austria

> > August 2, 2019

Abstract

We study a stochastic phase-field model for tumor growth dynamics coupling a stochastic Cahn-Hilliard equation for the tumor phase parameter with a stochastic reaction-diffusion equation governing the nutrient proportion. We prove strong well-posedness of the system in a general framework through monotonicity and stochastic compactness arguments. We introduce then suitable controls representing the concentration of cytotoxic drugs administered in medical treatment and we analyze a related optimal control problem. We derive existence of an optimal strategy and deduce first-order necessary optimality conditions by studying the corresponding linearized system and the backward adjoint system.

AMS Subject Classification: 35R60, 35K55, 49J20, 78A70.

Key words and phrases: stochastic systems of partial differential equations, Cahn-Hilliard equation, optimal control, first-order necessary conditions, tumor growth.

^{*}Acknowledgments. This research was supported by the Italian Ministry of Education, University and Research (MIUR): Dipartimenti di Eccellenza Program (2018–2022) – Dept. of Mathematics "F. Casorati", University of Pavia. In addition, it has been performed in the framework of the project Fondazione Cariplo-Regione Lombardia MEGAsTAR "Matematica d'Eccellenza in biologia ed ingegneria come acceleratore di una nuova strateGia per l'ATtRattività dell'ateneo pavese". The present paper also benefits from the support of the GNAMPA (Gruppo Nazionale per l'Analisi Matematica, la Probabilità e le loro Applicazioni) of INdAM (Istituto Nazionale di Alta Matematica) through the project "Trasporto ottimo per dinamiche con interazione". LS is also funded by Vienna Science and Technology Fund (WWTF) through Project MA14-009.

Contents

1	Inti	roduction	2
2	Ger	neral setting and main results	6
3	Well-posedness of the state system		12
	3.1	The approximated problem	12
	3.2	Uniform estimates	16
	3.3	Continuous dependence on the data	19
	3.4	Stochastic compactness and passage to the limit	20
4	Refined well-posedness		23
	4.1	Refined existence	23
	4.2	Refined continuous dependence	25
5	Optimal control problem		27
	5.1	Existence of an optimal control	27
	5.2	The linearized system	28
	$5.2 \\ 5.3$	The linearized system	28 31

1 Introduction

In the recent years phase-field systems have been used to describe many complex systems, in particular related to biomedical applications and specially to tumor growth dynamics. In this paper we consider a version of a phase-field model for tumor growth recently introduced in [29], where we have neglected the effects of chemotaxis and active transport. The new feature of the present work consists in adding stochastic terms in both the PDEs ruling the tumor-dynamic and then studying a related optimal control problem.

This model describes the evolution of a tumor mass surrounded by healthy tissues by taking into account proliferation of cells via nutrient consumption and apoptosis. In particular, the model under consideration in this paper fits into the framework of *diffuse interface* models for tumor growth. In this setting the evolution of the tumor is described by an order parameter φ which represents the local concentration of tumor cells; the interface between the tumor and healthy cells is supposed to be represented by a narrow transition layer separating the pure regions where $\varphi = \pm 1$, with $\varphi = 1$ denoting the tumor phase and $\varphi = -1$ the healthy phase.

We consider here the case of an incipient tumor, i.e., before the development of quiescent cells, when the equation ruling the evolution of the tumor growth process is often given by a Cahn-Hilliard (CH) equation [6] for φ coupled with a reaction-diffusion equation for the nutrient σ (cf., e.g., [13,29,34,35]). We just mention here that more sophisticated models have been also developed, possibly including different tumor phases (e.g., proliferating and necrotic), or incorporating the effects of fluid flow in the system evolution. In this direction, multiphase Cahn-Hilliard-Darcy systems [1, 14, 24, 28, 29, 59] have been analyzed in the deterministic case. Further studies on tumor growth modelling are presented in [43,49] with a particular emphasis to emergence of resistance to therapy.

All the references listed above deal with deterministic representations of the tumor growth. However, it is widely accepted that tumor-dynamics can be regarded as a random evolution due to stochastic

proliferation and differentiation of cells (cf. [56]). Tumor metastases, for example, are generally activated randomly by biological signals originating at the cellular level and these processes influence the critical proliferation rates that directly govern the evolution of tumor cells. In [7,45] the authors deal with stochastic angiogenesis models with the aim of generating more realistic structures of capillary networks. Other interesting contributions in the bio-medical literature (cf., e.g., [37]) are concerned with stochastic avascular models, taking into account the uncertainty of the (most influential) parameters. Finally, many studies have used stochastic perturbations in an attempt to model the effects of treatments on tumor growth. In this way, additive noises could represents environmental disruptions caused by therapeutic procedures in cancer treatment: indeed, these are delivered to the entire tumor tissue, and most likely to the healthy surrounding tissues as well. Although this appears far too simplified to represent real clinical procedures, it can be of acceptable accuracy when applied to the evolution of tumors at early stages. We refer to [44] for the modelling of treatment in the logistic tumor growth dynamic by means of a suitable (white) noise.

General PDE-models of tumor growth that account for randomness are rare in literature. For what concerns diffuse-interface descriptions, up to our knowledge, in this paper we give a first contribution in this direction. The stochastic perturbation that we are considering has a twofold motivation.

On one hand, we adopt a statistical approach to deal with the (too much) complicated free energy describing the bio-medical properties of the system. To do it, we directly add to the CH-equation a additive noise taking into account all the microscopical fluctuation affecting the evolution of the phase parameter. A natural choice for such a perturbation would be a space-time white noise. However, from the mathematical point of view, space-time whites noises are difficult to handle in higher space-dimension, and one usually considers smoothed-in-space noises by introducing a suitable covariance operator.

On the other hand, we introduce a multiplicative noise in the reaction diffusion equation (which actually depends on the nutrient concentration) with the aim of modelling the effects of angiogenensis. A stochastic forcing of this type is indeed related to the oxygen received by cancerous cells: this may result in enhancing its effectiveness, and therefore its contribution, to the total growth process of the tumor.

More precisely, including the two stochastic terms into the deterministic system we end up with the following stochastic Cahn-Hilliard-reaction-diffusion model for tumor growth:

$$d\varphi - \Delta \mu \, dt = (\mathcal{P}\sigma - a - \alpha u)h(\varphi) \, dt + G \, dW_1 \qquad \text{in } (0, T) \times D \,, \tag{1.1}$$

$$\mu = -A\Delta\varphi + B\psi'(\varphi) \qquad \text{in } (0,T) \times D, \qquad (1.2)$$

$$d\sigma - \Delta\sigma dt + c\sigma h(\varphi) dt + b(\sigma - w) dt = \mathcal{H}(\sigma) dW_2 \qquad \text{in } (0, T) \times D,$$

$$\partial_{\mathbf{n}}\varphi = \partial_{\mathbf{n}}\mu = \partial_{\mathbf{n}}\sigma = 0 \qquad \text{in } (0, T) \times \partial D,$$
(1.3)

$$\partial_{\mathbf{n}}\varphi = \partial_{\mathbf{n}}\mu = \partial_{\mathbf{n}}\sigma = 0 \qquad \text{in } (0,T) \times \partial D , \qquad (1.4)$$

$$\varphi(0) = \varphi_0, \quad \sigma(0) = \sigma_0 \qquad \text{in } D. \tag{1.5}$$

Here, $D \subset \mathbb{R}^3$ is a smooth bounded domain with smooth boundary, T > 0 is a fixed final time, W_1, W_2 are independent cylindrical Wiener processes on separable Hilbert spaces U_1 and U_2 , respectively, defined on a stochastic basis $(\Omega, \mathscr{F}, (\mathscr{F}_t)_{t \in [0,T]}, \mathbb{P}), G$ is a stochastically integrable operator with respect to W_1 and \mathcal{H} is a suitable Lipschitz-type operator. Notice that the initial configuration of the system can be chosen random as well.

The parameters $\mathcal{P}, a, \alpha, A, B, c, b$ are assumed to be strictly positive constants. Namely, \mathcal{P} denotes the tumor proliferation rate, a the apoptosis rate, α the effectiveness rate of the cytotoxic drugs, c the nutrient consumption rate, and b the nutrient supply rate, while A and B are related to the tickness of the interface between the pure phases. The function h is assumed to be monotone increasing, nonnegative in the "physical" interval [-1,1], and normalized so that h(-1) = 0 and h(1) = 1. The term $\mathcal{P}\sigma h(\varphi)$ models the proliferation of tumor cells, which is proportional to the concentration of the nutrient, the term $ah(\varphi)$ describes the apoptosis (or death) of tumor cells, and $c\sigma h(\varphi)$ represents the consumption of the nutrient by the tumor cells, which is higher if more tumor cells are present. The control variables are u in (1.1) and w in (1.3), which can be interpreted as a therapy (chemotherapy and antiangiogenic therapy, respectively, for example) distribution entering the system, either via the mass balance equation or the nutrient (cf. also [30], [10], and [11]) for similar choices for the control variables). Finally, ψ' stands for the derivative of a double-well potential ψ . A typical example of potential, meaningful in view of applications, has the expression

$$\psi_{pol}(r) = \frac{1}{4}(r^2 - 1)^2, \quad r \in \mathbb{R}.$$
(1.6)

We may observe that in the first part of our analysis (related to well-posedness of the system) we can allow for more general regular potentials having at least cubic and at most exponential growth at infinity.

The mathematical literature on the stochastic Cahn-Hilliard and Allen-Cahn equations is quite developed: we refer for example to [3,47] for the stochastic Allen-Cahn equation, and to [12,15,16,20,31,50,51]) for the stochastic Cahn-Hilliard equation. For completeness, let us quote also [21,22] for a study on a stochastic diffuse interface model involving the Cahn-Hilliard and Navier-Stokes equations. Nevertheless, similar results for coupled stochastic Cahn-Hilliard reaction-diffusion systems were not previously studied, up to our knowledge: in this sense, this contribution can be seen as a first work in this direction.

After proving the well-posedness of the SPDE-system above, we are interested here in the study of the following optimal control problem:

(CP) Minimize the cost functional

$$\begin{split} J(\varphi, u, w) &:= \frac{\beta_1}{2} \mathbb{E} \int_Q |\varphi - \varphi_Q|^2 + \frac{\beta_2}{2} \mathbb{E} \int_D |\varphi(T) - \varphi_T|^2 + \frac{\beta_3}{2} \mathbb{E} \int_D (\varphi(T) + 1) \\ &+ \frac{\beta_4}{2} \mathbb{E} \int_Q |u|^2 + \frac{\beta_5}{2} \mathbb{E} \int_Q |w|^2 \,, \end{split}$$

subject to the control constraint $(u, w) \in \mathcal{U}$, where

$$\mathcal{U} := \left\{ (u, w) \in L^2(\Omega; L^2(0, T; H))^2 \text{ progr. measurable: } 0 \le u, w \le 1 \text{ a.e. in } \Omega \times (0, T) \times D \right\},$$

and to the state system (1.1)–(1.5).

Let us comment on the optimal control problem (CP). The function φ_Q indicate some desired evolution for the tumor cells and φ_T stands for a desired final distribution of tumor cells (for example suitable for surgery). The first two terms of J are of standard tracking type, while the third term of J measures the size of the tumor at the end of the treatment. The fourth and fifth terms penalize large concentrations of the cytotoxic drugs through integral over the full space-time domain of the squared nutrient and drug concentrations. As it is presented in J, a large value of $|\varphi - \varphi_T|^2$ would mean that the patient suffers from the growth of the tumor during the treatment, and a large value of $|u|^2$ (or $|w|^2$) would mean that the patient suffers from high toxicity of the drugs. The nonnegative coefficients β_i indicate the importance of conflicting targets given in the strategy to avoid unnecessary harm to the patient, and at the same time increment the quality of the approximation of φ_Q , φ_T . By the optimal control problem (CP), we aim at searching for a medical strategy (u, w) such that the corresponding tumor concentration is as close as possible to the targets φ_Q and φ_T , the tumoral size at the end of the treatment is minimal, and the total amount of nutrient or drug supplied (which is restricted by the control constraints) does not inflict any harm to the patient.

We shall prove, under suitable regularity on the data, the existence of a *relaxed* optimal control and derive first-order necessary optimality conditions by exploiting a stochastic counterpart of the maximum principle à la Pontryagin. To do so, we introduce a backward (stochastic) system and we formulate necessary conditions for optimality through a variational inequality. The optimal control problem is studied in the particular case of the classical double well potential (1.6) and in case $\mathcal{H} \equiv 0$, meaning that we are neglecting the stochastic perturbation entering in the nutrient equation and that the only source of randomness is contatined in the CH-equation: see Remark 2.5 for further comments on this topic.

In the recent mathematical literature a number of results related to optimal control for deterministic tumor growth models has appeared. We can quote [8–10, 30, 53] for models coupling different variants of Cahn-Hilliard equations with reaction-diffusion equations and [18, 19, 55] for the case when also the velocity dynamic has been taken into account. Moreover, an optimal control problem taking into account also damages to healthy tissues was studied in [49] in the context of integro-differential modelling. For what concerns optimal control problems and necessary condition for optimality in the stochastic setting we can refer to the monography [60] for a general overview. The infinite dimensional case is tackled e.g. in [25, 26, 46], dealing with stochastic heat equation and reaction-diffusion systems. Optimal control for stochastic Cahn-Hilliard equations is discussed in the recent paper [52] where the control variable is contained in the chemical potential equation.

C. Orrieri, E. Rocca, L. Scarpa

Let us comment on the mathematical difficulties of this work. First of all, we stress that the stochastic system (1.1)-(1.3) does not fall in any available framework for studying stochastic evolution equations: this is due to both the growth of ψ , which may be of any polynomial order or first-order exponential, and more importantly to the coupling terms appearing in (1.1) and (1.3), which prevent from having a global monotonicity property for the pair (φ, σ) in a reasonable product space. Indeed, (1.1)-(1.2) have a monotone behaviour in the dual space of $H^1(D)$, while (1.3) is monotone on $L^2(D)$. Such issue does not depend on the stochastic setting and represents a crucial difficulty also in the deterministic case, which also requires an *ad-hoc* analysis of the system. Let us explain now the main differences when dealing with the stochastic case in comparison with the deterministic case.

Let us start by mentioning that the standard compactness techniques used in the deterministic setting cannot be directly applied. This mainly results from the lack of compactness of the embedding from $L^r(\Omega; \mathcal{X})$ into $L^r(\Omega, \mathcal{Y})$, $1 \leq r \leq +\infty$, even if $\mathcal{X} \hookrightarrow \mathcal{Y}$ in a compact fashion. To bypass the problem, we rely on the combination of the Skorohod representation theorem with the well-known Gyöngy-Krylov's criterium [32]. The crucial ingredient for the all machinery to work is a pathwise uniqueness result for the system, which is not a straightforward consequence of its deterministic counterpart.

The main issue when dealing with the well-posedness of the system (1.1)-(1.5) is the presence of proliferation terms in the Cahn-Hilliard equation, resulting in the non-conservation of the spatial mean of φ . In the deterministic setting, this difficulty is usually overcome by estimating directly the mean of φ and by keeping track of it in the estimates on the solutions: in particular, such procedure hinges on the Lipschitz-continuity of h and on proving the boundedness of $\sigma \in [0, 1]$ through a maximum principle for the reaction-diffusion equation. By contrast, in the stochastic setting, proving the boundedness of σ in [0, 1] is much more delicate, due to a lack of a maximum principle in a general framework for (1.3). In our setting, we are still able to show that σ remains bounded in [0, 1] during the evolution thanks to suitable assumptions on the operator \mathcal{H} : roughly speaking, we require that the noise in equation (1.3) "switches off" whenever σ touches the values 0 or 1. This procedure is very common from the mathematical point of view, and is also pretty reasonable in terms of applications: indeed, as we have pointed out before, the multiplicative noise in (1.3) could model the angiogenesis phenomenon, i.e. the formation of new blood vessels to convey more nutrient towards the tumoral cells. Hence, it is very natural to assume that this is neglectable when $\sigma = 1$ (i.e. when the nutrient level is saturated) and $\sigma = 0$ (i.e. when in principle no nutrient is available).

In order to keep track of the estimate on the spatial average of φ one usually estimates φ in terms of σ from (1.1)–(1.2), σ in terms of φ from (1.3), and concludes by "closing the estimate" using the boundedness of σ and a Gronwall-type argument. While this procedure perfectly works in the deterministic setting, there are several points deserving attention in the stochastic case. The main problem is that, due to the presence of the noise terms, one is forced to use estimates in expectation through some maximal martingale inequalities, and cannot rely on any pathwise estimate (i.e. with $\omega \in \Omega$ being fixed). Consequently, when trying to close the estimate through the Gronwall lemma, some further regularity is needed on the solutions in terms of existence of exponential moments.

A related issue arising in the stochastic setting concerns the continuous dependence of the solution (φ, σ) on the data of the problem, in particular on the controls (u, w). Indeed, as it is common when dealing with optimal control problems, we need to deduce some continuous dependence properties with respect to the controls in stronger topologies in order to tackle the linearized system. As we have anticipated, this hinges again on the combination of estimates in expectations and the Gronwall lemma, resulting in a need for boundedness of exponential moments of the solutions. The continuous dependence property is achieved with much more difficulty with respect to the deterministic case, as a careful track of the specific moments of the solutions is necessary.

Let us now turn to the major issues arising in the stochastic optimal control problem. In the deterministic setting, the first step consists in showing the Fréchet differentiability of the control-to-state map in order to analyse the linearized system: this is usually achieved (also for degenerate potentials) by proving that $\psi''(\bar{\varphi}) \in L^{\infty}((0,T) \times D)$, where $\bar{\varphi}$ is the state variable corresponding to a fixed control (\bar{u}, \bar{w}) . Such a regularity can be obtained by requiring that the set of admissible controls is bounded in L^{∞} and performing a maximum-principle-argument. Nevertheless, in the stochastic case, there is no hope to prove an L^{∞} -bound on φ , due to the presence of the additive noise in (1.1). We overcome this problem as it was done in [52]. First, we prove that the linearized system admits a unique variational solution. Then, we show that the control-to-state mapping is Gâteaux differentiable only in a weak sense, and that its derivative can be identified as the unique solution to the linearized system: this is still enough to deduce a first-order variational inequality.

The second main difficulty in the stochastic context consists in solving the adjoint problem, which is a system of backward stochastic PDEs. Recall that inverting time in stochastic dynamics is not straightforward and results in the introduction of two additional variables (one for φ and one for σ), which guarantees that solutions are adapted to the fixed filtration $(\mathscr{F}_t)_t$. Moreover, in the present case, the classical variational theory for backward SPDEs cannot be applied directly due to the presence of the coupling terms and to the growth of the potential. To solve the problem, we firstly approximate the backward equations and subsequently exploit the (abstract) duality relation between the adjoint and the linearized dynamics. This permit to get uniform estimates on the adjoint variables through a (more easier) continuous dependence of the linearized system. For an application of this strategy to stochatic reaction-diffusion equations we refer to [26].

Let us finally mention that this work is meant as a first contribution in the direction of analyzing stochastic mathematical models of tumor growth and many relevant questions remain open: it would be interesting to validate the results with numerical simulations and medical data (we refer for example to [33] for a possible stochastic model of MRI data). Moreover, it would be interesting to investigate uniqueness of optimal control problems and to include the duration of the therapy as another control variable in the formulation of the cost functional (cf. [30] and [8] for similar analysis in the deterministic case).

Here is the plan of the paper. In Section 2 we state the main assumptions and the main results of the paper concerning existence and uniqueness of solutions, a refined well-posedness theorem and the results on the optimal control problem: existence of the optimal controls, Gâteaux differentiability of the control-to-state map and first order optimality conditions obtained by solving the adjoint system. The proofs of the first well-posedness result is given in Section 3 and the refined well-posedness is proved in Section 4. The last Section 5 is devoted to the optimal control problem and to the analysis of the adjoint system.

2 General setting and main results

Throughout the paper, $D \subset \mathbb{R}^3$ is a smooth bounded domain with smooth boundary Γ : we shall denote the space-time cylinder $(0,T) \times D$ by Q and use the classical notation Q_t for $(0,t) \times D$ with $t \in [0,T)$.

It is useful to introduce the spaces

$$H := L^2(D), \qquad V := H^1(D), \qquad Z := \left\{ u \in H^2(D) : \partial_{\mathbf{n}} u = 0 \text{ a.e. on } \Gamma \right\}.$$

For a generic element $y \in V^*$ we denote by y_D the average $y_D := \frac{1}{D} \langle y, 1 \rangle_V$ and we recall the Poincaré-Wirtinger inequality: there exists $M_D > 0$, only depending on D, such that

$$\|v - v_D\|_H \le M_D \|\nabla v\|_H \quad \forall v \in V.$$

$$(2.1)$$

Setting

$$V_0^* := \{ y \in V^* : y_D = 0 \}, \qquad V_0 := V \cap V_0^*$$

we can define the operator $\mathcal{N}: V_0^* \to V_0$ as the unique solution of the generalized Neumann problem

$$\int_{D} \nabla \mathcal{N} y \cdot \nabla \phi = \langle y, \phi \rangle_{V} \quad \forall \phi \in V \,, \qquad (\mathcal{N} y)_{D} = 0 \,.$$

This also provides an equivalent norm on V^* given by the following

$$||y||_*^2 := ||\nabla \mathcal{N}(y - y_D)||_H^2 + |y_D|^2, \quad y \in V^*$$

In particular, by the inequality (2.1) this guarantees that there exists $M'_D > 0$, only depending on D, such that

$$\|\mathcal{N}y\|_{V} \le M'_{D} \|\nabla \mathcal{N}y\|_{H} = \|y\|_{*} \quad \forall y \in V_{0}^{*}.$$
(2.2)

We also recall that the following compactness inequality holds: for every $\varepsilon > 0$ there exists $M_{\varepsilon} > 0$ such that

$$\|y\|_{H}^{2} \leq \varepsilon \|\nabla y\|_{H}^{2} + M_{\varepsilon} \|y\|_{V^{*}}^{2} \quad \forall y \in V.$$

$$(2.3)$$

Furthermore, $(\Omega, \mathscr{F}, (\mathscr{F}_t)_{t \in [0,T]}, \mathbb{P})$ is a filtered probability space satisfying the usual conditions, where T > 0 is a fixed final time, and W_1, W_2 are independent cylindrical Wiener processes on separable Hilbert spaces U_1 and U_2 , respectively. We shall assume that the filtration $(\mathscr{F}_t)_t$ is the one generated by (W_1, W_2) . Let also $(e_n)_n$ be an arbitrary, but fixed, complete orthonormal system of U_2 and denote with $\mathscr{L}^2(U_i, H)$ the space of Hilbert-Schmidt operators from U_i to H, i = 1, 2. In the following we frequently use the shorthand $F \cdot W_i := \int_0^T F \, \mathrm{d}W_i$, for suitable stochastically integrable operators F.

Throughout the work, we shall use the symbol $a \leq b$ for any $a, b \geq 0$ to indicate that there exists a constant c > 0, independent of a, b, such that $a \leq cb$. When the implicit constant c depends on some relevant external quantities that we want to keep track of, we shall write it as a subscript.

Fix now $p \ge 6$. The following assumptions will be in order throughout the paper:

- (A1) \mathcal{P} , a, α , b, c, A, B are positive constants, and $h : \mathbb{R} \to [0, 1]$ is a Lipschitz-continuous function with Lipschitz constant $L_h > 0$ such that h(-1) = 0 and h(1) = 1;
- (A2) $\psi : \mathbb{R} \to [0, +\infty)$ is of class C^2 and satisfies

$$\begin{aligned} |\psi''(x)| &\leq C_1(1+|\psi'(x)|),\\ \psi''(x) &\geq -C_2,\\ |\psi'(x) - \psi'(y)| &\leq C_3\left(1+|\psi''(x)| + |\psi''(y)|\right)|x-y|\end{aligned}$$

for some positive constants C_i , i = 1, 2, 3, and for every $x, y \in \mathbb{R}$;

- (A3) $G \in L^4(0,T; \mathscr{L}^2(U_1,V)) \cap L^{\infty}(0,T; \mathscr{L}^2(U_1,V^*))$ is progressively measurable;
- (A4) the sequence $(\hbar_n)_n \subset W^{1,\infty}(\mathbb{R})$ satisfies

$$L_{\mathcal{H}}^{2} := \sum_{n=0}^{\infty} \left\| \hbar_{n}^{\prime} \right\|_{L^{\infty}(\mathbb{R})}^{2} < +\infty, \qquad \hbar_{n}(0) = \hbar_{n}(1) = 0 \quad \forall n \in \mathbb{N}.$$

This implies that it is well-defined and $L_{\mathcal{H}}$ -Lipschitz-continuous the operator

 $\mathcal{H}: H \to \mathscr{L}^2(U_2, H), \qquad \mathcal{H}(x): e_n \mapsto \hbar_n(x) \quad \forall \, n \in \mathbb{N};$

- (A5) u, w are *H*-valued progressively measurable processes such that $u, w \in [0, 1]$ almost everywhere in $\Omega \times (0, T) \times D$;
- (A6) $\varphi_0 \in L^p(\Omega, \mathscr{F}_0, \mathbb{P}; V)$ and $\psi(u_0) \in L^{p/2}(\Omega, \mathscr{F}_0, \mathbb{P}; L^1(D));$
- (A7) $\sigma_0 \in L^2(\Omega, \mathscr{F}_0, \mathbb{P}; H)$ and $\sigma_0 \in [0, 1]$ a.e. in $\Omega \times D$.

Remark 2.1. Let us comment on the assumptions above. First of all, note that assumption (A2) includes both the classical polynomial double-well potential (1.6) as well as any double-well potential with at most first-order exponential growth at infinity. Secondly, a typical example of operator G satisfying (A3) is given by any fixed nonrandom and time-independent $G \in \mathscr{L}^2(U_1, V)$. Moreover, in assumption (A4) we are requiring that the multiplicative noise $\mathcal{H} dW_2$ acts along the directions e_n through the Lipschitzcontinuous function \hbar_n for all n. The requirement that \hbar_n vanishes at 0 and 1 heuristically means that the noise switches off as soon as σ touches the border-values 0 and 1. Finally, the assumptions (A6)–(A7) on the initial data are trivially satisfied when $\varphi_0 \in V$ and $\sigma \in H$ are nonrandom with $\psi(\varphi_0) \in L^1(D)$ and $\sigma \in [0, 1]$ almost everywhere in D.

The first three results concern the well-posedness of the state system. First of all, we prove existence of solutions and continuous dependence on the data in the most general framework (A1)-(A7) in Theorem 2.2–2.3. Then, in Theorem 2.4 we show that under additional assumptions on the data the solution to the system inherits further regularity and a stronger dependence on the controls can be written: this will be necessary in order to tackle the optimal control problem.

Theorem 2.2 (Existence). Under the assumptions (A1)–(A7) there exists a unique triplet (φ, μ, σ) of progressively measurable V-valued processes with

$$\varphi \in L^{p}\left(\Omega; C^{0}([0,T];H) \cap L^{\infty}(0,T;V) \cap L^{2}(0,T;Z)\right), \qquad \varphi - G \cdot W_{1} \in L^{p}(\Omega; H^{1}(0,T;V^{*})), \quad (2.4)$$
$$\psi'(\varphi) \in L^{p/2}(\Omega; L^{2}(0,T;H)), \quad (2.5)$$

$$\mu \in L^{p/2}(\Omega; L^2(0, T; V)), \qquad \nabla \mu \in L^p(\Omega; L^2(0, T; H)), \qquad \mu_D \in L^{p/2}(\Omega; L^\infty(0, T)), \tag{2.6}$$

$$\sigma \in L^{p}\left(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V)\right), \qquad \sigma - \mathcal{H}(\sigma) \cdot W_{2} \in L^{p}(\Omega;H^{1}(0,T;V^{*})), \qquad (2.7)$$

$$\sigma \in [0,1] \quad a.e. \ in \ \Omega \times Q \,, \tag{2.8}$$

$$\varphi(0) = \varphi_0, \quad \sigma(0) = \sigma_0 \qquad a.e. \text{ in } \Omega \times D, \qquad (2.9)$$

and satisfying

$$\left\langle \partial_t (\varphi - G \cdot W_1)(t), \zeta \right\rangle_V + \int_D \nabla \mu(t) \cdot \nabla \zeta = \int_D \left(\mathcal{P}\sigma(t) - a - \alpha u \right) h(\varphi(t))\zeta, \qquad (2.10)$$

$$\int_{D} \mu(t)\zeta = A \int_{D} \nabla\varphi(t) \cdot \nabla\zeta + B \int_{D} \psi'(\varphi(t))\zeta , \qquad (2.11)$$

$$\left\langle \partial_t (\sigma - \mathcal{H}(\sigma) \cdot W_2)(t), \zeta \right\rangle_V + \int_D \left[\nabla \sigma(t) \cdot \nabla \zeta + c\sigma(t)h(\varphi(t))\zeta + b(\sigma(t) - w(t))\zeta \right] = 0$$
(2.12)

for every $\zeta \in V$, for almost every $t \in (0, T)$, \mathbb{P} -almost surely.

Theorem 2.3 (Continuous dependence). Under assumptions (A1)–(A4), let the data $(u_1, w_1, \varphi_0^1, \sigma_0^1)$ and $(u_2, w_2, \varphi_0^2, \sigma_0^2)$ satisfy (A5)–(A8), and let $(\varphi_1, \mu_1, \sigma_1)$ and $(\varphi_2, \mu_2, \sigma_2)$ be two respective solutions satisfying (2.4)–(2.12). Then there exists a sequence $(\tau^N)_{N\in\mathbb{N}}$ of stopping times and a sequence $(M_N)_{N\in\mathbb{N}}$ of positive real numbers, depending on φ_1 and φ_2 , such that $\tau^N \nearrow T$ P-a.s. as $N \to \infty$, and

$$\mathbb{E} \sup_{t \in [0,\tau^{N}]} \left(\| (\varphi_{1} - \varphi_{2})(t) \|_{V^{*}}^{p} + \| (\sigma_{1} - \sigma_{2})(t) \|_{H}^{p} \right) + \mathbb{E} \left(\int_{0}^{\tau^{N}} \| \nabla(\varphi_{1} - \varphi_{2})(s) \|_{H}^{2} \, \mathrm{d}s \right)^{p/2} \\
+ \mathbb{E} \left(\int_{0}^{\tau^{N}} \| (\sigma_{1} - \sigma_{2})(s) \|_{V}^{2} \, \mathrm{d}s \right)^{p/2} \\
\leq M_{N} \left[\| \varphi_{0}^{1} - \varphi_{0}^{2} \|_{L^{p}(\Omega;V^{*})}^{p} + \| \sigma_{0}^{1} - \sigma_{0}^{2} \|_{L^{p}(\Omega;H)}^{p} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \| (u_{1} - u_{2})(s) \|_{V^{*}}^{2} \, \mathrm{d}s \right)^{p/2} \\
+ \mathbb{E} \left(\int_{0}^{\tau^{N}} \| (w_{1} - w_{2})(s) \|_{H}^{2} \, \mathrm{d}s \right)^{p/2} \right] \quad \forall N \in \mathbb{N}.$$
(2.13)

In particular, the problem (2.4)–(2.12) admits a unique solution.

Theorem 2.4 (Refined well-posedness). Under the assumptions (A1)-(A7), suppose also that

$$\exp\left(\beta \left\|\varphi_{0}\right\|_{H}^{2}\right) \in L^{1}(\Omega) \quad \forall \beta \geq 1,$$
(2.14)

$$G \in L^{\infty}(\Omega \times (0,T); \mathscr{L}^2(U_1,H)), \qquad (2.15)$$

$$\psi \in C^3(\mathbb{R}), \quad \exists C_4 > 0: \quad |\psi'''(r)| \le C_4(1+|r|) \quad \forall r \in \mathbb{R}.$$
 (2.16)

Then, the unique solution (φ, μ, σ) to (2.4)–(2.12) also satisfies

$$\varphi \in L^{p/3}(\Omega; L^2(0, T; H^3(D))),$$
(2.17)

$$\exp\left(\beta \left\|\varphi\right\|_{L^{2}(0,T;Z)}^{2}\right) \in L^{1}(\Omega) \quad \forall \beta \geq 1.$$
(2.18)

Furthermore, assume also that

$$\mathcal{H} \equiv 0. \tag{2.19}$$

Let the data $(u_1, w_1, \varphi_0^1, \sigma_0^1)$ and $(u_2, w_2, \varphi_0^2, \sigma_0^2)$ satisfy (A5)–(A7) and (2.14), and let $(\varphi_1, \mu_1, \sigma_1)$ and $(\varphi_2, \mu_2, \sigma_2)$ be the two respective solutions satisfying (2.4)–(2.12) and (2.18): then for every q > p there exists a positive constant $M_{q,p} = M_{q,p}(\varphi_1, \varphi_2)$ such that

$$\begin{aligned} \|\varphi_{1} - \varphi_{2}\|_{L^{p}(\Omega;C^{0}([0,T];V^{*}))} + \|\nabla(\varphi_{1} - \varphi_{2})\|_{L^{p}(\Omega;L^{2}(0,T;H))} + \|\sigma_{1} - \sigma_{2}\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;V))} \\ &\leq M_{q,p} \Big(\|\varphi_{0}^{1} - \varphi_{0}^{2}\|_{L^{q}(\Omega;V^{*})} + \|\sigma_{0}^{1} - \sigma_{0}^{2}\|_{L^{q}(\Omega;H)} \\ &+ \|u_{1} - u_{2}\|_{L^{q}(\Omega;L^{2}(0,T;V^{*}))} + \|w_{1} - w_{2}\|_{L^{q}(\Omega;L^{2}(0,T;H))} \Big) \,. \end{aligned}$$

$$(2.20)$$

If also

$$\exists r > \frac{2pq}{q-p}: \qquad \varphi_0 \in L^r(\Omega; V), \quad G \in L^r(\Omega; L^2(0, T; \mathscr{L}^2(U_1, V))), \qquad (2.21)$$

then there exists a positive constant $M_{p,q,r} = M_{p,q,r}(\varphi_1, \varphi_2)$ such that

$$\begin{aligned}
\varphi_{1} - \varphi_{2} \|_{L^{p}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;Z))} + \|\sigma_{1} - \sigma_{2}\|_{L^{p}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V))} \\
&\leq M_{p,q,r} \Big(\|\varphi_{0}^{1} - \varphi_{0}^{2}\|_{L^{q}(\Omega;H)} + \|\sigma_{0}^{1} - \sigma_{0}^{2}\|_{L^{q}(\Omega;H)} \\
&+ \|u_{1} - u_{2}\|_{L^{q}(\Omega;L^{2}(0,T;H))} + \|w_{1} - w_{2}\|_{L^{q}(\Omega;L^{2}(0,T;H))} \Big).
\end{aligned}$$
(2.22)

Remark 2.5. Note that the refined assumptions (2.14)-(2.16) and (2.21) are not much restrictive: indeed, they are still satisfied by the wide class of data $\varphi_0 \in V$ nonrandom, $G \in \mathscr{L}^2(U_1, V)$ fixed (nonrandom and time-independent) and by the classical double-well potential (1.6) of degree 4. The reason why we need to assume (2.19) is that, in order to obtain the continuous dependence property on the whole time interval [0, T], one needs to rely necessarily on pathwise estimates (i.e. with $\omega \in \Omega$ fixed) of $\sigma_1 - \sigma_2$ in terms of $\varphi_1 - \varphi_2$: this is possible only if the noise in equation (1.3) is of additive type. However, since it is also crucial to prove that $\sigma \in [0, 1]$ (hence also (A4) must be in order), the only reasonable choice is to require that $\mathcal{H} \equiv 0$ and to neglect the noise in the equation for σ .

We are now ready to introduce the general setting for the optimal control problem of the state system (1.1)–(1.5). From now on, we shall assume (A1)–(A7) (2.14)–(2.16), (2.19) and (2.21), so that the strongest continuous dependence property (2.22) holds, and we shall also fix an exponent q > p.

We introduce the set of admissible controls ${\mathcal U}$ as

$$\mathcal{U} := \left\{ (u, w) \in L^2(\Omega; L^2(0, T; H))^2 \text{ progr. measurable: } 0 \le u, w \le 1 \text{ a.e. in } \Omega \times (0, T) \times D \right\},$$

which is clearly a closed convex bounded subset of $L^{\infty}(\Omega \times Q)^2$. It will be useful to embed \mathcal{U} in an open bounded subset $\tilde{\mathcal{U}} \subset L^q(\Omega; L^2(0, T; H))^2 \cap L^{\infty}(\Omega \times Q)^2$.

We define the cost functional J as

$$\begin{split} J: L^2(\Omega; C^0([0,T];H)) \times L^2(\Omega; L^2(0,T;H)) \times L^2(\Omega; L^2(0,T;H)) &\longrightarrow [0,+\infty) \\ J(\varphi,u,w) &:= \frac{\beta_1}{2} \mathbb{E} \int_Q |\varphi - \varphi_Q|^2 + \frac{\beta_2}{2} \mathbb{E} \int_D |\varphi(T) - \varphi_T|^2 + \frac{\beta_3}{2} \mathbb{E} \int_D (\varphi(T) + 1) \\ &+ \frac{\beta_4}{2} \mathbb{E} \int_Q |u|^2 + \frac{\beta_5}{2} \mathbb{E} \int_Q |w|^2 \,, \end{split}$$

where $\beta_i \ge 0, i = 1, ..., 5$, are fixed constants and we assume that

$$\beta_1 \varphi_Q \in L^6(\Omega; L^2(0, T; H)), \qquad \beta_2 \varphi_T \in L^6(\Omega, \mathscr{F}_T; V).$$

Moreover, Theorems 2.2–2.4 ensure that the control-to-state map

$$S: \tilde{\mathcal{U}} \to L^p(\Omega; C^0([0,T]; H) \cap L^\infty(0,T; V) \cap L^2(0,T; Z)), \qquad (u,w) \mapsto \varphi$$

is well-defined and furthermore it is Lipschitz-continuous in the sense specified in (2.22): the proofs of Theorems 2.2–2.4 can be adapted with no difficulty to the more general case where $u, w \in L^{\infty}(\Omega \times Q)$ (not necessarily $0 \le u, w \le 1$). Consequently, it is natural to define the reduced cost functional

$$\tilde{J}:\tilde{\mathcal{U}}\to [0,+\infty)\,,\qquad \tilde{J}(u,w):=J(S(u,w),u,w)\,,\quad (u,w)\in\tilde{\mathcal{U}}$$

The optimal control problem associated to the state system (1.1)-(1.5) consists in minimizing the map \tilde{J} on \mathcal{U} . Note that even if J is convex, the nonlinearity of S does not ensure that \tilde{J} is convex as well, so that the minimization problem in nontrivial and uniqueness of optimal controls is difficult to prove. Furthermore, as already mentioned in the introduction, classical compactness results cannot be directly applied to the stochastic setting and this gives rise to the need of a weaker concept of optimality.

We introduce below the notion of *relaxed* optimal control, which is inspired by [2, Def. 2.4].

Definition 2.6. An optimal control is a pair $(u, w) \in \mathcal{U}$ such that

$$\tilde{J}(u,w) = \inf_{(v,z)\in\mathcal{U}} \tilde{J}(v,z)$$

A relaxed optimal control is a family

$$\left((\Omega',\mathscr{F}',(\mathscr{F}'_t)_{t\in[0,T]},\mathbb{P}'),W_1',\varphi_0',\sigma_0',u',w',\varphi',\mu',\sigma',\varphi_Q',\varphi_T'\right),$$

where $(\Omega', \mathscr{F}', (\mathscr{F}'_t)_{t \in [0,T]}, \mathbb{P}')$ is a complete filtered probability space, W'_1 is a $(\mathscr{F}_t)_t$ -cylindrical Wiener process, (φ'_0, σ'_0) is a \mathscr{F}'_0 -measurable $(V \times H)$ -valued random variable with the same law of (φ_0, σ_0) , (u', w') is a $(\mathscr{F}'_t)_t$ -progressively measurable H^2 -valued process such that $0 \leq u', w' \leq 1$ almost everywhere in $\Omega' \times Q$, $(\varphi', \mu', \sigma')$ is the unique solution to the state system (2.4)–(2.12) on the probability space $(\Omega' \mathscr{F}', \mathbb{P}')$ with respect to the data $(\varphi'_0, \sigma'_0, u', w', W'_1)$ (and the choice $\mathcal{H} \equiv 0$), $\beta_1 \varphi'_Q \in L^2(\Omega'; L^2(0, T; H))$ has the same law of $\beta_1 \varphi_Q$, $\beta_2 \varphi'_T \in L^6(\Omega', \mathscr{F}'_T; V)$ has the same law of $\beta_2 \varphi_T$, and

$$\begin{split} \tilde{J}'(u',w') &:= \frac{\beta_1}{2} \operatorname{\mathbb{E}}' \int_Q |\varphi' - \varphi'_Q|^2 + \frac{\beta_2}{2} \operatorname{\mathbb{E}}' \int_D |\varphi'(T) - \varphi'_T|^2 + \frac{\beta_3}{2} \operatorname{\mathbb{E}}' \int_D (\varphi'(T) + 1) \\ &+ \frac{\beta_4}{2} \operatorname{\mathbb{E}}' \int_Q |u'|^2 + \frac{\beta_5}{2} \operatorname{\mathbb{E}}' \int_Q |w'|^2 \leq \tilde{J}(v,z) \quad \forall (v,z) \in \mathcal{U} \,. \end{split}$$

Let us briefly comment on the notion of relaxed optimality. First notice that the idea of relaxation mimics the concept of weak (or martingale) solutions to stochastic equations, thus involving in the definition the probability space itself along with the Wiener process and the parameters of the model. This naturally allows for the application of the Skorohod representation theorem. Moreover, let us observe that by setting $\beta_4 = \beta_5 = 0$ in the definition of the cost functional, the problem reduces to the existence of nearest points of a fixed target in uniformly convex Banach spaces (see [2] for further details). Finally, given a relaxed optimal control it seems natural to wonder whether it admits or not a *strong* formulation. A possible way to tackle the problem could be the combination of a Gyöngy-Krylov argument with a uniqueness result for optimal controls.

The first result that we prove in this context is very natural and ensures that a relaxed optimal control for our problem always exists.

Theorem 2.7. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Then there exists a relaxed optimal control for problem (1.1)–(1.5).

Let us now concentrate on necessary conditions for optimality. To this end, we first introduce the following linearized system

$$\partial_t x_k - \Delta y_k = h(\varphi)(\mathcal{P}z_k - \alpha k_u) + h'(\varphi)x_k(\mathcal{P}\sigma - a - \alpha u)$$
$$y_k = A\Delta x_k + B\psi''(\varphi)x_k$$
$$\partial_t z_k - \Delta z_k + cz_k h(\varphi) + c\sigma h'(\varphi)x_k + b(z_k - k_w) = 0,$$

complemented with homogeneous Neumann boundary conditions for x_k, y_k and z_k with initial conditions given by $x_k(0) = z_k(0) = 0$. Existence and uniqueness of variational solutions to the above system is the content of the following theorem.

Theorem 2.8. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Let $(u, w) \in \mathcal{U}$, $k := (k_u, k_w) \in \mathcal{U}$ and set $\varphi := S(u, w)$. Then there exists a unique triple (x_k, y_k, z_k) with

$$x_k \in L^p(\Omega; H^1(0, T; Z^*) \cap L^2(0, T; Z)), \qquad y_k \in L^p(\Omega; L^2(0, T; H)),$$
(2.23)

$$z_k \in L^p(\Omega; H^1(0, T; V^*) \cap L^2(0, T; V)), \qquad (2.24)$$

such that $x_k(0) = z_k(0) = 0$ and

$$\langle \partial_t x_k, \zeta \rangle_V - \int_D y_k \Delta \zeta = \int_D \left[h(\varphi) (\mathcal{P} z_k - \alpha k_u) + h'(\varphi) x_k (\mathcal{P} \sigma - a - \alpha u) \right] \zeta, \qquad (2.25)$$

$$\int_{D} y_k \zeta = A \int_{D} \nabla x_k \cdot \nabla \zeta + B \int_{D} \psi''(\varphi) x_k \zeta , \qquad (2.26)$$

$$\langle \partial_t z_k, \zeta \rangle_V + \int_D \nabla z_k \cdot \nabla \zeta + \int_D \left[c z_k h(\varphi) + c \sigma h'(\varphi) x_k + b(z_k - k_w) \right] \zeta = 0$$
(2.27)

for every $\zeta \in Z$, for almost every $t \in (0,T)$, \mathbb{P} -almost surely.

Furthermore, the control-to-state map S is differentiable in a suitable sense and we characterize its derivative as the unique solution to the linearized system.

Proposition 2.9. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Then the control-to-state map $S: \tilde{\mathcal{U}} \to L^p(\Omega; C^0([0,T];H) \cap L^2(0,T;Z))$ is Gâteaux-differentiable in the following sense: for every $(u,w) \in \tilde{\mathcal{U}}$ and $(k_u,k_w) \in L^q(\Omega; L^2(0,T;H))^2$ we have

$$\frac{S((u,w) + \varepsilon(k_u, k_w)) - S(u,w)}{\varepsilon} \to x_k \quad in \ L^{\ell}(\Omega; L^2(0,T;V)) \quad \forall \ell \in [1,p),$$

$$\frac{S((u,w) + \varepsilon(k_u, k_w)) - S(u,w)}{\varepsilon} \to x_k \quad in \ L^p(\Omega; H^1(0,T;Z^*) \cap L^2(0,T;Z))$$

as $\varepsilon \to 0$, where x_k is the unique first solution component to the linearized system (2.23)–(2.27).

The next result is a first version of necessary conditions for optimality.

Proposition 2.10. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Let $(\bar{u}, \bar{w}) \in \mathcal{U}$ be an optimal control and let $\bar{\varphi} := S(\bar{u}, \bar{w})$. Then, for every $(u, w) \in \mathcal{U}$, setting $k := (k_u, k_w) := (u - \bar{u}, w - \bar{w})$, we have

$$\beta_1 \mathbb{E} \int_Q (\bar{\varphi} - \varphi_Q) x_k + \beta_2 \mathbb{E} \int_D (\bar{\varphi}(T) - \varphi_T) x_k(T) + \frac{\beta_3}{2} \mathbb{E} \int_D x_k(T) + \beta_4 \mathbb{E} \int_Q \bar{u} k_u + \beta_5 \mathbb{E} \int_Q \bar{w} k_w \ge 0,$$

where x_k is the unique first solution component to the linearized system (2.23)–(2.27).

In the last results of the paper we show how to remove the dependence on x_k in the first-order necessary conditions for optimality. To this end, we analyse the corresponding adjoint problem, which is a system of backward SPDEs of the form

$$-\mathrm{d}\pi - A\Delta\tilde{\pi}\,\mathrm{d}t + B\psi''(\varphi)\tilde{\pi}\,\mathrm{d}t = h'(\varphi)(\mathcal{P}\sigma - a - \alpha u)\pi\,\mathrm{d}t - ch'(\varphi)\sigma\rho\,\mathrm{d}t + \beta_1(\varphi - \varphi_Q)\,\mathrm{d}t - \xi\,\mathrm{d}W_1\,,$$
$$\tilde{\pi} = -\Delta\pi\,,$$
$$-\mathrm{d}\rho - \Delta\rho\,\mathrm{d}t + ch(\varphi)\rho\,\mathrm{d}t + b\rho\,\mathrm{d}t = \mathcal{P}h(\varphi)\pi\,\mathrm{d}t - \theta\,\mathrm{d}W_2\,,$$

complemented with homogeneous Neumann boundary conditions for π , $\tilde{\pi}$ and ρ , with final conditions

$$\pi(T) = \beta_2(\varphi(T) - \varphi_T) + \frac{\beta_3}{2}, \qquad \rho(T) = 0.$$

The last two results deal with the existence and uniqueness of solutions to the adjoint system, and with the simplified version of the first-order necessary conditions for optimality, respectively.

Theorem 2.11. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Let $(u, w) \in \tilde{\mathcal{U}}$ be fixed and set $\varphi := S(u, w)$. Then there exists a unique quintuplet $(\pi, \tilde{\pi}, \xi, \rho, \theta)$ with

$$\pi \in L^2(\Omega; C^0([0, T]; V) \cap L^2(0, T; Z \cap H^3(D))), \qquad (2.28)$$

$$\tilde{\pi} \in L^2(\Omega; C^0([0, T]; V^*) \cap L^2(0, T; V)), \qquad (2.29)$$

$$\rho \in L^2(\Omega; C^0([0, T]; H) \cap L^2(0, T; V)), \qquad (2.30)$$

$$\xi \in L^2(\Omega; L^2(0, T; \mathscr{L}^2(U_1, V))), \qquad \theta \in L^2(\Omega; L^2(0, T; \mathscr{L}^2(U_2, H))),$$
(2.31)

such that, for every $\zeta \in V$,

$$\int_{D} \pi(t)\zeta + A \int_{t}^{T} \int_{D} \nabla \tilde{\pi}(s) \cdot \nabla \zeta \, \mathrm{d}s + B \int_{t}^{T} \int_{D} \psi''(\varphi(s))\tilde{\pi}(s)\zeta \, \mathrm{d}s = -\int_{D} \left(\int_{t}^{T} \xi(s) \, \mathrm{d}W_{1}(s)\right)\zeta + \int_{t}^{T} \int_{D} \left[h'(\varphi(s))(\mathcal{P}\sigma(s) - a - \alpha u(s))\pi(s) - ch'(\varphi(s))\sigma(s)\rho(s) + \beta_{1}(\varphi - \varphi_{Q})(s)\right]\zeta \, \mathrm{d}s,$$

$$(2.32)$$

$$\langle \tilde{\pi}(t), \zeta \rangle_V = \int_D \nabla \pi(t) \cdot \nabla \zeta ,$$
 (2.33)

$$\int_{D} \rho(t)\zeta + \int_{t}^{T} \int_{D} \nabla \rho(s) \cdot \nabla \zeta \, \mathrm{d}s + \int_{t}^{T} \int_{D} \left[ch(\varphi(s))\rho(s) + b\rho(s) \right] \zeta \, \mathrm{d}s$$
$$= \mathcal{P} \int_{t}^{T} \int_{D} h(\varphi(s))\pi(s)\zeta \, \mathrm{d}s - \int_{D} \left(\int_{t}^{T} \theta(s) \, \mathrm{d}W_{2}(s) \right) \zeta$$
(2.34)

for every $t \in [0, T]$, \mathbb{P} -almost surely.

The final version of the stochastic maximum principle reads as follows

Theorem 2.12. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Let $(\bar{u}, \bar{w}) \in \mathcal{U}$ be an optimal control and let $\bar{\varphi} := S(\bar{u}, \bar{w})$. Then

$$\mathbb{E}\int_{Q}(\beta_{4}\bar{u}-\alpha h(\bar{\varphi})\pi)(u-\bar{u})+\mathbb{E}\int_{Q}(\beta_{5}\bar{w}+b\rho)(w-\bar{w})\geq 0\qquad\forall(u,w)\in\mathcal{U}\,,$$

where (π, ρ) are the first-component solutions to the adjoint system (2.32)-(2.34). In particular, if $\beta_4 > 0$ and $\beta_5 > 0$ then (\bar{u}, \bar{w}) is the orthogonal projection of $(\frac{\alpha}{\beta_4}h(\bar{\varphi})\pi, -\frac{b}{\beta_5}\rho)$ on \mathcal{U} .

Remark 2.13. Note that the first-order condition for optimality can be equivalently rewritten pointwise in $\Omega \times [0, T]$ by using standard localization techniques: see e.g. [60].

3 Well-posedness of the state system

This section contains the proof of the Theorem 2.2. First of all, we consider an approximated problem where the nonlinearity ψ' is smoothed out through a Yosida-type regularization. Secondly, we prove uniform estimates on the approximated solutions in suitable spaces, and through the theorems of Prokhorov and Skorokhod we are then able to show existence of (probabilistic) weak solutions. Finally, we prove a pathwise uniqueness result for the original system, yielding existence and uniqueness also of strong solutions thanks to a well-known criterion.

3.1 The approximated problem

Thanks to assumption (A2), the function $\gamma(r) := \psi'(r) + C_2 r$, $r \in \mathbb{R}$, is nondecreasing, hence can be identified with a maximal monotone graph in $\mathbb{R} \times \mathbb{R}$. For any $\lambda > 0$, let γ_{λ} be its Yosida approximation and set

$$\psi'_{\lambda}(r) := \gamma_{\lambda}(r) - C_2 r, \quad r \in \mathbb{R}$$

We consider the approximated problem

$$d\varphi_{\lambda} - \Delta \mu_{\lambda} dt = (\mathcal{P}\sigma_{\lambda} - a - \alpha u)h(\varphi_{\lambda}) dt + G dW_{1} \quad \text{in } (0, T) \times D, \qquad (3.1)$$

 $\mu_{\lambda} = -A\Delta\varphi_{\lambda} + B\psi_{\lambda}'(\varphi_{\lambda}) \qquad \text{in } (0,T) \times D, \qquad (3.2)$

$$d\sigma_{\lambda} - \Delta \sigma_{\lambda} dt + c\sigma_{\lambda} h(\varphi_{\lambda}) dt + b(\sigma_{\lambda} - w) dt = \mathcal{H}(\sigma_{\lambda}) dW_2 \quad \text{in } (0, T) \times D, \qquad (3.3)$$

$$\partial_{\mathbf{n}}\varphi_{\lambda} = \partial_{\mathbf{n}}\mu_{\lambda} = \partial_{\mathbf{n}}\sigma_{\lambda} = 0 \qquad \text{in } (0,T) \times \partial D , \qquad (3.4)$$

$$\varphi_{\lambda}(0) = \varphi_0, \quad \sigma_{\lambda}(0) = \sigma_0 \quad \text{in } D.$$
 (3.5)

We prove existence and uniqueness of a solution $(\varphi_{\lambda}, \mu_{\lambda}, \sigma_{\lambda})$ through a fixed point argument: we fix a suitable ϕ in the third equation in place of φ_{λ} , and we solve (3.3) obtaining thus a solution component σ_{λ}^{ϕ} depending on ϕ . We substitute σ_{λ}^{ϕ} into (3.1) and we solve (3.1)–(3.2), getting the other two solution components $(\varphi_{\lambda}^{\phi}, \mu_{\lambda}^{\phi})$ in terms on ϕ . Finally, we show that the map $\phi \mapsto \varphi_{\lambda}^{\phi}$ is well-defined in a suitable space and is a contraction.

Let us fix a progressively measurable *H*-valued process ϕ with

$$\phi \in L^2\left(\Omega; L^2(0, T; H)\right)$$

Since h takes values in [0, 1] it is clear now that the initial-value problem

$$\begin{cases} \mathrm{d}\sigma_{\lambda}^{\phi} - \Delta\sigma_{\lambda}^{\phi}\,\mathrm{d}t + c\sigma_{\lambda}^{\phi}h(\phi)\,\mathrm{d}t + b(\sigma_{\lambda}^{\phi} - w)\,\mathrm{d}t = \mathcal{H}(\sigma_{\lambda})\,\mathrm{d}W_{2} & \text{in } (0,T) \times D \,, \\ \partial_{\mathbf{n}}\sigma_{\lambda}^{\phi} = 0 & \text{in } (0,T) \times \partial D \\ \sigma_{\lambda}^{\phi}(0) = \sigma_{0} & \text{in } D \,, \end{cases}$$

admits a unique solution (see [30, 38])

$$\sigma_{\lambda}^{\phi} \in L^2\left(\Omega; C^0([0,T];H) \cap L^2(0,T;V)\right) \,.$$

Let us show now that $\sigma_{\lambda} \in [0, 1]$ almost everywhere in $\Omega \times Q$. Introduce first a smooth approximation $s_{\varepsilon} : \mathbb{R} \to \mathbb{R}$ of the function $r \mapsto (r-1)_+, r \in \mathbb{R}$: for example take

$$s_{\varepsilon}(r) := \begin{cases} 0 & \text{if } r \leq 1 \,, \\ \frac{1}{2\varepsilon}(r-1)^2 & \text{if } r \in [1, 1+\varepsilon] \,, \\ r-1-\frac{\varepsilon}{2} & \text{if } r > 1+\varepsilon \,, \end{cases}$$

so that $\hat{s}_{\varepsilon} := \int_{1}^{\cdot} s_{\varepsilon}(r) dr \in C^{2}(\mathbb{R})$. Associate to \hat{s}_{ε} an integral operator of the form $\mathcal{S}_{\varepsilon} := \int_{D} \hat{s}_{\varepsilon}(\sigma(x)) dx$, where $\sigma \in H$, and exploit the (sub-quadratic) growth of \hat{s}_{ε} to show that for any $\sigma, v \in H$

$$\mathcal{S}'_{\varepsilon}(\sigma)v = \int_{D} s_{\varepsilon}(\sigma)v, \qquad [\mathcal{S}''_{\varepsilon}(\sigma)v] = s'_{\varepsilon}(\sigma(x))v(x).$$
 (3.6)

This permit the application of Itô's formula in the form of [48, Thm. 4.1]

$$\mathbb{E} \int_{D} \hat{s}_{\varepsilon}(\sigma_{\lambda}^{\phi}(t)) + \mathbb{E} \int_{Q_{t}} s_{\varepsilon}'(\sigma_{\lambda}^{\phi}) |\nabla \sigma_{\lambda}^{\phi}|^{2} + c \mathbb{E} \int_{Q_{t}} h(\phi) \sigma_{\lambda}^{\phi} s_{\varepsilon}(\sigma_{\lambda}^{\phi}) + b \mathbb{E} \int_{Q_{t}} (\sigma_{\lambda}^{\phi} - w) s_{\varepsilon}(\sigma_{\lambda}^{\phi}) \\ = \mathbb{E} \int_{D} \hat{s}_{\varepsilon}(\sigma_{0}) + \frac{1}{2} \mathbb{E} \sum_{n=0}^{\infty} \int_{Q_{t}} s_{\varepsilon}'(\sigma_{\lambda}^{\phi}) |\hbar_{n}(\sigma_{\lambda}^{\phi})|^{2}.$$

Since s_{ε} is increasing and identically 0 in $(-\infty, 0]$, the second and third term on the left-hand side are nonnegative. Moreover, using the fact that $s'_{\varepsilon} \leq 1$ and that \hbar_n is Lipschitz-continuous with $\hbar_n(1) = 0$, recalling that $\sigma_0 \in [0, 1]$ a.e. in $\Omega \times D$ we get

$$\mathbb{E}\int_{D} \hat{s}_{\varepsilon}(\sigma_{\lambda}^{\phi}(t)) + b \mathbb{E}\int_{Q_{t}} (\sigma_{\lambda}^{\phi} - w) s_{\varepsilon}(\sigma_{\lambda}^{\phi}) \lesssim \frac{1}{2} \mathbb{E}\int_{\{\sigma_{\lambda}^{\phi} > 1\}} |\sigma_{\lambda}^{\phi} - 1|^{2} = \frac{1}{2} \mathbb{E}\int_{Q_{t}} |(\sigma_{\lambda}^{\phi} - 1)_{+}|^{2}.$$

Letting $\varepsilon \searrow 0$ we deduce that

$$\frac{1}{2} \mathbb{E} \int_D |(\sigma_\lambda^\phi - 1)_+(t)|^2 + b \mathbb{E} \int_{Q_t} (\sigma_\lambda^\phi - w)(\sigma_\lambda^\phi - 1)_+ \le \frac{1}{2} \mathbb{E} \int_{Q_t} |(\sigma_\lambda^\phi - 1)_+|^2,$$

where, since $w \in [0, 1]$ a.e. in $\Omega \times Q$,

$$(\sigma_{\lambda}^{\phi} - w)(\sigma_{\lambda}^{\phi} - 1)_{+} = (\sigma_{\lambda}^{\phi} - 1)(\sigma_{\lambda}^{\phi} - 1)_{+} + (1 - w)(\sigma_{\lambda}^{\phi} - 1)_{+} \ge 0 \quad \text{a.e. in } \Omega \times Q \,.$$

The Gronwall lemma yields then $(\sigma_{\lambda}^{\phi} - 1)_{+} = 0$, hence $\sigma_{\lambda}^{\phi} \leq 1$, a.e in $\Omega \times Q$. Arguing similarly with $-(\sigma_{\lambda}^{\phi})_{-}$, one also deduces that $\sigma_{\lambda}^{\phi} \geq 0$ almost everywhere, so that

$$\sigma_{\lambda}^{\phi} \in [0, 1]$$
 a.e. in $\Omega \times Q$

Let us substitute now σ^{ϕ}_{λ} in (3.1) and consider the Cahn-Hilliard system

$$\begin{cases} \mathrm{d}\varphi_{\lambda}^{\phi} - \Delta \mu_{\lambda}^{\phi} \, \mathrm{d}t = (\mathcal{P}\sigma_{\lambda}^{\phi} - a - \alpha u)h(\varphi_{\lambda}^{\phi}) \, \mathrm{d}t + G \, \mathrm{d}W_{1} & \text{ in } (0, T) \times D \,, \\ \mu_{\lambda}^{\phi} = -A\Delta \varphi_{\lambda}^{\phi} + B\psi_{\lambda}'(\varphi_{\lambda}^{\phi}) & \text{ in } (0, T) \times D \,, \\ \partial_{\mathbf{n}}\varphi_{\lambda}^{\phi} = \partial_{\mathbf{n}}\mu_{\lambda}^{\phi} = 0 & \text{ in } (0, T) \times \partial D \,, \\ \varphi_{\lambda}^{\phi}(0) = \varphi_{0} & \text{ in } D \,. \end{cases}$$

Arguing as in [50,51], by the Lipschitz-continuity of ψ'_{λ} and the fact that h, u and σ^{ϕ}_{λ} are bounded almost everywhere, there is a unique solution $(\varphi^{\phi}_{\lambda}, \mu^{\phi}_{\lambda})$ with

$$\begin{split} \varphi^{\phi}_{\lambda} \in L^2\left(\Omega; C^0([0,T];H) \cap L^{\infty}(0,T;V) \cap L^2(0,T;Z)\right) \,, \\ \mu^{\phi}_{\lambda} \in L^2(\Omega; L^2(0,T;V)) \,. \end{split}$$

For every $\lambda > 0$, it is well-defined then the map

$$\Phi_{\lambda}: L^2\left(\Omega; L^2(0,T;H)\right) \to L^2\left(\Omega; C^0([0,T];H) \cap L^{\infty}(0,T;V) \cap L^2(0,T;Z)\right), \qquad \phi \mapsto \varphi_{\lambda}^{\phi}$$

Let us show that Φ_{λ} is a contraction: let $\phi_1, \phi_2 \in L^2(\Omega; L^2(0, T; H))$ progressively measurable. First of all, from the third equation we have that

$$d(\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}) - \Delta(\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}) dt + c \left(\sigma_{\lambda}^{\phi_1} h(\phi_1) - \sigma_{\lambda}^{\phi_2} h(\phi_2)\right) dt + b(\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}) dt$$
$$= \left(\mathcal{H}(\sigma_{\lambda}^{\phi_1}) - \mathcal{H}(\sigma_{\lambda}^{\phi_2})\right) dW_2,$$

with $(\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2})(0) = 0$. Hence, Itô's formula for the square of the *H*-norm yields

$$\frac{1}{2} \left\| (\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}})(t) \right\|_{H}^{2} + \int_{Q_{t}} |\nabla(\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}})|^{2} + \int_{Q_{t}} (ch(\phi_{1}) + b) |\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}}|^{2} \\
= c \int_{Q_{t}} (h(\phi_{2}) - h(\phi_{1})) \sigma_{\lambda}^{\phi_{2}} (\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}}) + \frac{1}{2} \int_{0}^{t} \left\| (\mathcal{H}(\sigma_{\lambda}^{\phi_{1}}) - \mathcal{H}(\sigma_{\lambda}^{\phi_{2}}))(s) \right\|_{\mathscr{L}^{2}(U_{2},H)}^{2} \, \mathrm{d}s \qquad (3.7) \\
+ \int_{0}^{t} \left((\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}})(s), (\mathcal{H}(\sigma_{\lambda}^{\phi_{1}}) - \mathcal{H}(\sigma_{\lambda}^{\phi_{2}}))(s) \right)_{H} \, \mathrm{d}W_{2}(s) \, .$$

Now, employing the Lipschitz-continuity of h and the fact that $\sigma_{\lambda}^{\phi_2} \in [0,1]$ a.e. in $\Omega \times Q$ we have

$$c\int_{Q_t} \left(h(\phi_2) - h(\phi_1)\right) \sigma_{\lambda}^{\phi_2}(\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}) \le \frac{1}{2}\int_{Q_t} |\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}|^2 + \frac{c^2 L_h^2}{2}\int_Q |\phi_1 - \phi_2|^2.$$

Moreover, the Lipschitz continuity of \mathcal{H} yields

$$\int_0^t \left\| (\mathcal{H}(\sigma_{\lambda}^{\phi_1}) - \mathcal{H}(\sigma_{\lambda}^{\phi_2}))(s) \right\|_{\mathscr{L}^2(U_2, H)}^2 \, \mathrm{d}s \le L_{\mathcal{H}}^2 \int_{Q_t} |\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}|^2$$

while the Burkholder-Davis-Gundy (cf. [27, Lemma 4.1]) and Young inequalities (see [39, Lemma 4.3]) imply that, for some positive constants M and M',

$$\begin{split} \mathbb{E} \sup_{r \in [0,t]} \left| \int_{0}^{r} \left((\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}})(s), (\mathcal{H}(\sigma_{\lambda}^{\phi_{1}}) - \mathcal{H}(\sigma_{\lambda}^{\phi_{2}}))(s) \right)_{H} \mathrm{d}W_{2}(s) \right| \\ & \leq M \mathbb{E} \left(\int_{0}^{t} \left\| (\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}})(s) \right\|_{H}^{2} \left\| (\mathcal{H}(\sigma_{\lambda}^{\phi_{1}}) - \mathcal{H}(\sigma_{\lambda}^{\phi_{2}}))(s) \right\|_{\mathscr{L}^{2}(U_{2},H)}^{2} \mathrm{d}s \right)^{1/2} \\ & \leq M L_{\mathcal{H}} \mathbb{E} \left(\left\| \sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}} \right\|_{C^{0}([0,t];H)} \left\| \sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}} \right\|_{L^{2}(0,t;H)} \right) \\ & \leq \frac{1}{4} \mathbb{E} \left\| \sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}} \right\|_{C^{0}([0,t];H)}^{2} + M' \mathbb{E} \int_{Q_{t}} |\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}}|^{2} \,. \end{split}$$

Hence, taking supremum in time and expectations, rearranging the terms and using the Gronwall lemma, we deduce that

$$\left\|\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}}\right\|_{L^{2}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V))} \leq M \left\|\phi_{1} - \phi_{2}\right\|_{L^{2}(\Omega; L^{2}(0,T;H))},$$
(3.8)

where M > 0 only depends on c, L_h , and L_H . Secondly, from the first two equations we have

$$\partial_t (\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2}) - \Delta(\mu_{\lambda}^{\phi_1} - \mu_{\lambda}^{\phi_2}) = (\mathcal{P}\sigma_{\lambda}^{\phi_1} - a - \alpha u)h(\varphi_{\lambda}^{\phi_1}) - (\mathcal{P}\sigma_{\lambda}^{\phi_2} - a - \alpha u)h(\varphi_{\lambda}^{\phi_2}),$$
$$\mu_{\lambda}^{\phi_1} - \mu_{\lambda}^{\phi_2} = -A\Delta(\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2}) + B\psi_{\lambda}'(\varphi_{\lambda}^{\phi_1}) - B\psi_{\lambda}'(\varphi_{\lambda}^{\phi_2})$$

with initial condition $(\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2})(0) = 0$. Hence testing the first by $\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2}$ and integrating by parts yield

$$\begin{split} \left\| (\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}})(t) \right\|_{H}^{2} + A \int_{Q_{t}} |\Delta(\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}})|^{2} \\ &= \int_{Q_{t}} \left((\mathcal{P}\sigma_{\lambda}^{\phi_{1}} - a - \alpha u)h(\varphi_{\lambda}^{\phi_{1}}) - (\mathcal{P}\sigma_{\lambda}^{\phi_{2}} - a - \alpha u)h(\varphi_{\lambda}^{\phi_{2}}) \right) (\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}}) \\ &+ B \int_{Q_{t}} \left(\psi_{\lambda}'(\varphi_{\lambda}^{\phi_{1}}) - \psi_{\lambda}'(\varphi_{\lambda}^{\phi_{2}}) \right) \Delta(\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}}) \,. \end{split}$$

Hence, using the Young inequality on the right-hand side yields

$$\begin{split} \left\| (\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}})(t) \right\|_{H}^{2} + A \int_{Q_{t}} |\Delta(\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}})|^{2} \\ &\leq \int_{Q_{t}} |\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}}|^{2} + \int_{Q_{t}} |\mathcal{P}\sigma_{\lambda}^{\phi_{1}} - a - \alpha u|^{2} |h(\varphi_{\lambda}^{\phi_{1}}) - h(\varphi_{\lambda}^{\phi_{2}})|^{2} + \mathcal{P}^{2} \int_{Q_{t}} |\sigma_{\lambda}^{\phi_{1}} - \sigma_{\lambda}^{\phi_{2}}|^{2} |h(\varphi_{\lambda}^{\phi_{2}})|^{2} \\ &+ \frac{A}{2} \int_{Q_{t}} |\Delta(\varphi_{\lambda}^{\phi_{1}} - \varphi_{\lambda}^{\phi_{2}})|^{2} + \frac{B^{2}}{2A} \int_{Q_{t}} |\psi_{\lambda}'(\varphi_{\lambda}^{\phi_{1}}) - \psi_{\lambda}'(\varphi_{\lambda}^{\phi_{2}})|^{2} \,, \end{split}$$

from the Lipschitz-continuity of h and ψ'_{λ} , and the fact that σ^{ϕ}_{λ} and h are bounded in [0,1] we get

$$\begin{split} \left\| (\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2})(t) \right\|_{H}^{2} + \frac{A}{2} \int_{Q_t} |\Delta(\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2})|^2 \\ & \leq \left(1 + (\mathcal{P}L_h)^2 + \frac{B^2}{2A} \left| \frac{1}{\lambda} \vee C_2 \right|^2 \right) \int_{Q_t} |\varphi_{\lambda}^{\phi_1} - \varphi_{\lambda}^{\phi_2}|^2 + \mathcal{P}^2 \int_{Q_t} |\sigma_{\lambda}^{\phi_1} - \sigma_{\lambda}^{\phi_2}|^2 \, . \end{split}$$

The Gronwall lemma implies that, possibly updating the value of the implicit constant M,

$$\left\|\varphi_{\lambda}^{\phi_{1}}-\varphi_{\lambda}^{\phi_{2}}\right\|_{L^{2}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;Z))} \leq M_{\mathcal{P},L_{h},A,B,C_{2},\lambda}\left\|\sigma_{\lambda}^{\phi_{1}}-\sigma_{\lambda}^{\phi_{2}}\right\|_{L^{2}(\Omega;L^{2}(0,T;H))}.$$
(3.9)

Combining now (3.8) and (3.9) we obtain that

$$\left\|\varphi_{\lambda}^{\phi_{1}}-\varphi_{\lambda}^{\phi_{2}}\right\|_{L^{2}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;Z))} \leq M_{c,\mathcal{P},L_{h},L_{\mathcal{H}},A,B,C_{2},\lambda}\sqrt{T} \left\|\phi_{1}-\phi_{2}\right\|_{L^{2}(\Omega;L^{2}(0,T;H))}$$

We deduce that there is $T_0 > 0$ sufficiently small such that Φ_{λ} is a contraction on $L^2(\Omega; L^2(0, T_0; H))$, hence admits a fixed point in $L^2(\Omega; L^2(0, T_0; H))$. A standard patching technique yields now the existence and uniqueness of a fixed point φ_{λ} in the whole space $L^2(\Omega; L^2(0, T; H))$. Setting now $\mu_{\lambda} := \mu_{\lambda}^{\varphi_{\lambda}}$ and $\sigma_{\lambda} := \sigma_{\lambda}^{\varphi_{\lambda}}$, the definition of Φ_{λ} itself ensures also that

$$\varphi_{\lambda} \in L^{2}\left(\Omega; C^{0}([0,T];H) \cap L^{\infty}(0,T;V) \cap L^{2}(0,T;Z)\right), \qquad \mu_{\lambda} \in L^{2}(\Omega; L^{2}(0,T;V)), \\ \sigma_{\lambda} \in L^{2}\left(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V)\right), \qquad \sigma_{\lambda} \in [0,1] \quad \text{a.e. in } \Omega \times Q,$$

and that $(\varphi_{\lambda}, \mu_{\lambda}, \sigma_{\lambda})$ is the unique solution to the approximated problem (3.1)–(3.5).

3.2 Uniform estimates

Here we perform uniform a-priori estimates with respect to the parameter λ .

First estimate. We write Itô's formula for $\frac{1}{2} \|\sigma_{\lambda}\|_{H}^{2}$, getting

$$\frac{1}{2} \|\sigma_{\lambda}(t)\|_{H}^{2} + \int_{Q_{t}} |\nabla\sigma_{\lambda}|^{2} + \int_{Q_{t}} (ch(\varphi_{\lambda}) + b) |\sigma_{\lambda}|^{2} = \frac{1}{2} \|\sigma_{0}\|_{H}^{2} + b \int_{Q_{t}} w\sigma_{\lambda} \\
+ \frac{1}{2} \int_{0}^{t} \|\mathcal{H}(\sigma_{\lambda}(s))\|_{\mathscr{L}^{2}(U_{2},H)}^{2} \, \mathrm{d}s + \int_{0}^{t} (\sigma_{\lambda}(s), \mathcal{H}(\sigma_{\lambda}(s)))_{H} \, \mathrm{d}W_{2}(s) \,.$$
(3.10)

Since $w \leq 1$ almost everywhere we have, by the Young inequality,

$$b\int_{Q_t} w\sigma_\lambda \lesssim_b 1 + \int_{Q_t} |\sigma_\lambda|^2 \, .$$

Moreover, the Lipschitz continuity of \mathcal{H} and the fact that $\hbar_n(0) = 0$ for all $n \in \mathbb{N}$ yields

$$\int_0^t \left\| \mathcal{H}(\sigma_\lambda(s)) \right\|_{\mathscr{L}^2(U_2,H)}^2 \, \mathrm{d}s \le L_{\mathcal{H}}^2 \int_{Q_t} |\sigma_\lambda|^2$$

Finally, by the Burkholder-Davis-Gundy and Young inequalities (see e.g. [40, Lem. 4.1]), for every $\varepsilon > 0$ and a certain $M_{\varepsilon} > 0$ we get

$$\begin{aligned} \left\| \sup_{r \in [0,t]} \int_0^r \left(\sigma_{\lambda}(s), \mathcal{H}(\sigma_{\lambda}(s)) \right)_H \, \mathrm{d}W_2(s) \right\|_{L^{p/2}(\Omega)} &\lesssim \left\| \left(\int_0^t \left\| \sigma_{\lambda}(s) \right\|_H^2 \left\| \mathcal{H}(\sigma_{\lambda}(s)) \right\|_{\mathscr{L}^2(U_2,H)}^2 \, \mathrm{d}s \right)^{1/2} \right\|_{L^{p/2}(\Omega)} \\ &\leq \varepsilon \left\| \sigma_{\lambda} \right\|_{L^p(\Omega;C^0([0,t];H))}^2 + M_{\varepsilon} \left\| \mathcal{H}(\sigma_{\lambda}) \right\|_{L^p(\Omega;L^2(0,t;\mathscr{L}^2(0,t;\mathscr{L}^2(0,T;H)))}^2 \\ &\leq \varepsilon \left\| \sigma_{\lambda} \right\|_{L^p(\Omega;C^0([0,t];H))}^2 + M_{\varepsilon} L^2_{\mathcal{H}} \left\| \sigma_{\lambda} \right\|_{L^p(\Omega;L^2(0,T;H))}^2 . \end{aligned}$$

Taking supremum in time and $L^{p/2}(\Omega)$ -norm in (3.10), recalling that $\sigma_{\lambda} \in [0, 1]$, there exists M > 0, independent of λ and depending only on T and $\|\sigma_0\|_{L^p(\Omega;H)}^2$, such that

$$\|\sigma_{\lambda}\|_{L^{p}(\Omega; C^{0}([0,T];H)\cap L^{2}(0,T;V))\cap L^{\infty}(\Omega\times Q)} \leq M, \qquad (3.11)$$

hence, by the Lipschitz continuity of ${\mathcal H}$ and by comparison in the equation itself,

$$\left\|\mathcal{H}(\sigma_{\lambda})\right\|_{L^{p}(\Omega; C^{0}([0,T]; \mathscr{L}^{2}(U_{2},H)))} + \left\|\sigma_{\lambda} - \mathcal{H}(\sigma_{\lambda}) \cdot W_{2}\right\|_{L^{p}(\Omega; H^{1}(0,T; V^{*}))} \leq M.$$
(3.12)

Second estimate. We write Itô's formula for the square of the *H*-norm of φ_{λ} : taking into account the boundary conditions for φ_{λ} and μ_{λ} , integrating by parts we have

$$\frac{1}{2} \|\varphi_{\lambda}(t)\|_{H}^{2} + A \int_{Q_{t}} |\Delta\varphi_{\lambda}|^{2} - B \int_{Q_{t}} \psi_{\lambda}'(\varphi_{\lambda}) \Delta\varphi_{\lambda} = \frac{1}{2} \|\varphi_{0}\|_{H}^{2} + \int_{Q_{t}} (\mathcal{P}\sigma_{\lambda} - a - \alpha u)h(\varphi_{\lambda})\varphi_{\lambda} + \frac{1}{2} \int_{0}^{t} \|G(s)\|_{\mathscr{L}^{2}(U_{1},H)}^{2} \, \mathrm{d}s + \int_{0}^{t} (\varphi_{\lambda}(s), G(s))_{H} \, \mathrm{d}W_{1}(s) \, .$$

Recalling the definition of ψ'_{λ} and the fact that σ_{λ} , h and u are bounded in [0, 1], we deduce

$$\begin{aligned} \left|\varphi_{\lambda}(t)\right|_{H}^{2} + 2A \int_{Q_{t}} \left|\Delta\varphi_{\lambda}\right|^{2} + 2B \int_{Q_{t}} \gamma_{\lambda}''(\varphi_{\lambda}) \left|\nabla\varphi_{\lambda}\right|^{2} &\leq \left\|\varphi_{0}\right\|_{H}^{2} + 2\mathcal{P} \int_{Q_{t}} \left|\varphi_{\lambda}\right| \\ &- C_{2} \int_{Q_{t}} \varphi_{\lambda} \Delta\varphi_{\lambda} + \left\|G\right\|_{L^{2}(0,T;\mathscr{L}^{2}(U_{1},H))}^{2} + \int_{0}^{t} \left(\varphi_{\lambda}(s), G(s)\right)_{H} \, \mathrm{d}W_{1}(s) \,. \end{aligned}$$

The Young inequality yields then

$$\begin{aligned} \|\varphi_{\lambda}(t)\|_{H}^{2} + \int_{Q_{t}} |\Delta\varphi_{\lambda}|^{2} + \int_{Q_{t}} \gamma_{\lambda}''(\varphi_{\lambda})|\nabla\varphi_{\lambda}|^{2} \\ \lesssim_{A,B,\mathcal{P},a,\alpha,C_{2}} 1 + \|\varphi_{0}\|_{H}^{2} + \int_{Q_{t}} |\varphi_{\lambda}|^{2} + \|G\|_{L^{2}(0,T;\mathscr{L}^{2}(U,H))}^{2} + \int_{0}^{t} (\varphi_{\lambda}(s),G(s))_{H} \, \mathrm{d}W_{1}(s) \,, \end{aligned}$$

where the implicit constant is independent of λ . It is a standard matter to check that the Burkholder-Davis-Gundy and Young inequalities on the right-hand side ensure that, for every $\delta > 0$,

$$\left\|\sup_{r\in[0,t]}\int_0^r \left(\varphi_{\lambda}(s), G(s)\right)_H \,\mathrm{d}W(s)\right\|_{L^{p/2}(\Omega)}^{p/2} \lesssim \delta \operatorname{\mathbb{E}}\sup_{r\in[0,t]} \left\|\varphi_{\lambda}(r)\right\|_H^p + \frac{1}{4\delta} \left\|G\right\|_{L^p(\Omega; L^2(0,T; \mathscr{L}^2(U,H)))}^p$$

Hence, taking supremum in time and $L^{p/2}(\Omega)$ -norm, choosing δ sufficiently small and rearranging the terms, we infer that there exists a positive constant M, independent of λ , such that

$$\|\varphi_{\lambda}\|_{L^{p}(\Omega; C^{0}([0,T];H)\cap L^{2}(0,T;Z))} \leq M.$$
(3.13)

Third estimate. The idea is to write Itô's formula for the free-energy functional

$$\mathcal{E}_{\lambda}(\varphi_{\lambda}) := \frac{A}{2} \int_{D} |\nabla \varphi_{\lambda}|^{2} + B \int_{D} \psi_{\lambda}(\varphi_{\lambda})$$

Note that the regularity of φ_{λ} and ψ_{λ} are not enough in order to do so, as \mathcal{E}_{λ} may not be twice Fréchetdifferentiable in V, and φ_{λ} is not necessarily continuous in V: consequently, a rigorous approach would require a further approximation on the problem. However, since this is not restrictive in our direction, we shall proceed formally in order to avoid heavy notations, and refer to [50,51] for a rigorous approach instead. Noting that, by (3.2), for every $x, y, y_1, y_2 \in V$

$$D\mathcal{E}_{\lambda}(x)[y] = \langle -A\Delta(x) + B\psi'_{\lambda}(x), y \rangle_{V} = \langle \mu_{\lambda}, y \rangle,$$
$$D_{2}\mathcal{E}_{\lambda}(x)[y_{1}, y_{2}] = \int_{D} \nabla y_{1} \cdot \nabla y_{2} + \int_{D} \psi''(x)y_{1}y_{2},$$

we have that

$$\begin{split} \frac{A}{2} \int_{D} |\nabla \varphi_{\lambda}(t)|^{2} + B \int_{D} \psi_{\lambda}(\varphi_{\lambda}(t)) + \int_{Q_{t}} |\nabla \mu_{\lambda}|^{2} \\ &= \frac{A}{2} \int_{D} |\nabla \varphi_{0}|^{2} + B \int_{D} \psi_{\lambda}(\varphi_{0}) + \int_{Q_{t}} (\mathcal{P}\sigma_{\lambda} - a - \alpha u)h(\varphi_{\lambda})\mu_{\lambda} \\ &+ \frac{1}{2} \int_{0}^{t} \operatorname{Tr} \left(G^{*}(s)D^{2}\mathcal{E}_{\lambda}(\varphi_{\lambda}(s))G(s) \right) \, \mathrm{d}s + \int_{0}^{t} (\mu_{\lambda}(s), G(s))_{H} \, \mathrm{d}W_{1}(s) \, . \end{split}$$

On the right-hand side, the fact that σ_{λ} , h and u are bounded in [0, 1] together with the Poincaré and Young inequalities imply that, for every $\delta > 0$,

$$\begin{split} \int_{Q_t} (\mathcal{P}\sigma_{\lambda} - a - \alpha u) h(\varphi_{\lambda}) \mu_{\lambda} &\leq (\mathcal{P} + a + \alpha) \int_{Q_t} |\mu_{\lambda} - (\mu_{\lambda})_D| + (\mathcal{P} + a + \alpha) \int_0^t |(\mu_{\lambda}(s))_D| \,\mathrm{d}s \\ &\lesssim_{\mathcal{P}, a, \alpha} \frac{1}{\delta} + \delta \int_{Q_t} |\nabla \mu_{\lambda}|^2 + \int_0^t |(\mu_{\lambda}(s))_D| \,\mathrm{d}s \,, \end{split}$$

where the implicit constant is independent of λ . Now, formally writing

$$(\mu_{\lambda}, G)_{H} = (\mu_{\lambda} - (\mu_{\lambda})_{D}, G)_{H} + |D|(\mu_{\lambda})_{D}G_{D} \lesssim \|\nabla\mu_{\lambda}\|_{H} \|G\|_{V^{*}} + (\mu_{\lambda})_{D} \|G\|_{V^{*}},$$

the quadratic variation of the stochastic integral on the right-hand side can be bounded employing the Young inequality by

$$\left(\int_0^t (\|\nabla \mu_\lambda(s)\|_H^2 + |(\mu_\lambda(s))_D|^2) \|G(s)\|_{\mathscr{L}^2(U_1,V^*)}^2 \, \mathrm{d}s \right)^{1/2} \\ \lesssim \delta \int_{Q_t} |\nabla \mu_\lambda|^2 + \frac{1}{\delta} \|G\|_{L^\infty(0,T;\mathscr{L}^2(U_1,V^*))}^2 + \|G\|_{L^\infty((0,T);\mathscr{L}^2(U_1,V^*))} \|(\mu_\lambda)_D\|_{L^2(0,t)}$$

Hence, using the Burkholder-Davis-Gundy inequality we end up with

$$\begin{split} & \left\| \sup_{r \in [0,t]} \int_{0}^{r} \left(\mu_{\lambda}(s), G(s) \right)_{H} \, \mathrm{d}W_{1}(s) \right\|_{L^{p/2}(\Omega)}^{p/2} \\ & \lesssim \delta^{p/2} \mathbb{E} \left(\int_{Q_{t}} |\nabla \mu_{\lambda}|^{2} \right)^{p/2} + \frac{1}{\delta^{p/2}} \left\| G \right\|_{L^{p}(\Omega; L^{\infty}(0,T; \mathscr{L}^{2}(U_{1}, V^{*})))}^{p} \\ & + \left\| G \right\|_{L^{\infty}(\Omega \times (0,T); \mathscr{L}^{2}(U_{1}, V^{*}))}^{p/2} \mathbb{E} \left\| (\mu_{\lambda})_{D} \right\|_{L^{2}(0,t)}^{p/2} \\ & \lesssim \delta^{p/2} \mathbb{E} \left(\int_{Q_{t}} |\nabla \mu_{\lambda}|^{2} \right)^{p/2} + \frac{1}{\delta^{p/2}} + \sqrt{t} \mathbb{E} \sup_{r \in [0,t]} |(\mu_{\lambda}(r))_{D}|^{p/2} \,, \end{split}$$

where the implicit constant depends on $||G||_{L^{\infty}(\Omega \times (0,T);\mathscr{L}^2(U_1,V^*))}$ but is independent of λ . Finally, given a complete orthonormal system $(e_j)_j$ of U_1 , since $\psi'' \leq 1 + |\psi'|$ and $V \hookrightarrow L^6(D)$, by the Hölder inequality we have

$$\int_{0}^{t} \operatorname{Tr} \left(G^{*}(s) D^{2} \mathcal{E}_{\lambda}(\varphi_{\lambda}(s)) G(s) \right) \, \mathrm{d}s = \int_{0}^{t} \sum_{j=0}^{\infty} \left(\left\| \nabla G(s) e_{j} \right\|_{H}^{2} + \int_{D} \psi_{\lambda}''(\varphi_{\lambda}(s)) ||G(s) e_{j}|^{2} \right) \, \mathrm{d}s$$
$$\leq \left\| G \right\|_{L^{2}(0,T;\mathscr{L}^{2}(U_{1},V))}^{2} + \int_{0}^{t} \left\| \psi_{\lambda}''(\varphi_{\lambda}(s)) \right\|_{L^{3/2}(D)} \left\| G(s) \right\|_{\mathscr{L}^{2}(U_{1},V)}^{2} \, \mathrm{d}s$$
$$\lesssim 1 + \left\| G \right\|_{L^{2}(0,T;\mathscr{L}^{2}(U_{1},V))}^{2} + \left\| \psi_{\lambda}'(\varphi_{\lambda}) \right\|_{L^{2}(0,t;H)} \left\| G \right\|_{L^{4}(0,T;\mathscr{L}^{2}(U_{1},V))}^{2} \, .$$

Hence, exploiting (3.2), we have that, for every $\varepsilon > 0$ and a certain $C_{\varepsilon} > 0$,

$$\int_{0}^{t} \operatorname{Tr}\left(G^{*}(s)D^{2}\mathcal{E}_{\lambda}(\varphi_{\lambda}(s))G(s)\right) \, \mathrm{d}s \lesssim 1 + \|\Delta\varphi_{\lambda}\|_{L^{2}(0,T;H)}^{2} + \varepsilon \int_{Q_{t}} |\nabla\mu_{\lambda}|^{2} + C_{\varepsilon} \|G\|_{L^{4}(0,T;\mathscr{L}^{2}(U_{1},V))}^{4} + \|G\|_{L^{\infty}(\Omega;L^{4}(0,T;\mathscr{L}^{2}(U_{1},V)))}^{2} \sup_{r\in[0,t]} \sqrt{t} |(\mu_{\lambda}(r))_{D}|.$$

Consequently, taking supremum in time and $L^{p/2}(\Omega)$ norm in Itô's formula, choosing δ, ε sufficiently small and rearranging the terms yield, for every $T_0 \in (0, T]$,

$$\begin{aligned} & \left\| \sup_{r \in [0,T_0]} \int_D |\nabla \varphi_{\lambda}(r)|^2 \right\|_{L^{p/2}(\Omega)} + \left\| \sup_{r \in [0,T_0]} \int_D \psi_{\lambda}(\varphi_{\lambda}(r)) \right\|_{L^{p/2}(\Omega)} + \left\| \int_{Q_{T_0}} |\nabla \mu_{\lambda}|^2 \right\|_{L^{p/2}(\Omega)} \\ & \lesssim 1 + \left\| \varphi_0 \right\|_{L^p(\Omega;V)}^2 + \left\| \psi(\varphi_0) \right\|_{L^{p/2}(\Omega;L^1(D))} + \left\| \Delta \varphi_{\lambda} \right\|_{L^p(\Omega;L^2(0,T;H))}^2 + \sqrt{T_0} \|(\mu_{\lambda})_D\|_{L^{p/2}(\Omega;L^{\infty}(0,T_0))} \end{aligned}$$

where the implicit constant only depends on $||G||_{L^{\infty}(\Omega \times (0,T); \mathscr{L}^{2}(U_{1},H))}$, $||G||_{L^{\infty}(\Omega; L^{4}(0,T; \mathscr{L}^{2}(U_{1},V)))}$ and on the data $A, B, \mathcal{P}, a, \alpha, C_{3}$. Now, thanks to (3.13), the fact that $|\psi'| \leq 1 + \psi$ on the left-hand side yields, by comparison in (3.2),

$$\left\| \sup_{r \in [0,T_0]} \int_D |\nabla \varphi_{\lambda}(r)|^2 \right\|_{L^{p/2}(\Omega)} + \left\| \sup_{r \in [0,T_0]} |(\mu_{\lambda}(r))_D| \right\|_{L^{p/2}(\Omega)} + \left\| \int_{Q_{T_0}} |\nabla \mu_{\lambda}|^2 \right\|_{L^{p/2}(\Omega)} \\ \lesssim 1 + \left\| \varphi_0 \right\|_{L^p(\Omega;V)}^2 + \left\| \psi(\varphi_0) \right\|_{L^{p/2}(\Omega;L^1(D))} + \sqrt{T_0} \left\| (\mu_{\lambda})_D \right\|_{L^{p/2}(\Omega;L^{\infty}(0,T_0))},$$

so that, choosing first T_0 small enough and then using a classical patching argument, we infer that

$$\|\varphi_{\lambda}\|_{L^{p}(\Omega;L^{\infty}(0,T;V))} + \|\nabla\mu_{\lambda}\|_{L^{p}(\Omega;L^{2}(0,T;H))} + \|(\mu_{\lambda})_{D}\|_{L^{p/2}(\Omega;L^{\infty}(0,T))} \le M.$$
(3.14)

By comparison in (3.2) we deduce then also that

$$\|\mu_{\lambda}\|_{L^{p/2}(\Omega;L^{2}(0,T;V))} + \|\psi_{\lambda}'(\varphi_{\lambda})\|_{L^{p/2}(\Omega;L^{2}(0,T;H))} \le M.$$
(3.15)

3.3 Continuous dependence on the data

Here we prove the continuous dependence contained in Theorem 2.3. Given $(\varphi_1, \mu_1, \sigma_1)$ and $(\varphi_2, \mu_2, \sigma_2)$ two solutions to (2.4)–(2.12) with sources (u_1, w_1) , (u_2, w_2) , and initial data $(\varphi_0^1, \sigma_0^1)$, $(\varphi_0^2, \sigma_0^2)$, respectively, we denote by $\varphi := \varphi_1 - \varphi_2$, $\mu := \mu_1 - \mu_2$, $\sigma := \sigma_1 - \sigma_2$, $u := u_1 - u_2$, and $w := w_1 - w_2$.

The equation for the differences in the Cahn-Hilliard system reads

$$\partial_t \varphi - \Delta \mu = (\mathcal{P}\sigma - \alpha u) h(\varphi_1) + (\mathcal{P}\sigma_2 - a - \alpha u_2) (h(\varphi_1) - h(\varphi_2))$$
(3.16)

$$\mu = -A\Delta\varphi + B(\psi'(\varphi_1) - \psi'(\varphi_2)), \qquad (3.17)$$

with initial data given by $\varphi_0 := \varphi_0^1 - \varphi_0^2$ and $\sigma_0 := \sigma_0^1 - \sigma_0^2$. To get the required estimate we use the same strategy as in [42]. Firstly, we integrate (3.16) over D

$$\partial_t \varphi_D = \int_D \left(\mathcal{P}\sigma + \alpha u \right) h(\varphi_1) + \int_D \left(\mathcal{P}\sigma_2 - a - \alpha u_2 \right) \left(h(\varphi_1) - h(\varphi_1) \right)$$
(3.18)

then we test by φ_D to get

$$\|\varphi_D\|_{C^0([0,t])}^2 \lesssim |(\varphi_0)_D|^2 + \|\sigma\|_{L^2(0,t;H)}^2 + \|u\|_{L^2(0,t;V^*)}^2 + \|\varphi\|_{L^2(0,t;H)}^2,$$

thanks to the boundedness of h, σ_2, u_2 and the Gronwall lemma. Taking the difference between (3.16) and (3.18) and testing by $\mathcal{N}(\varphi - \varphi_D)$ we have

$$\langle \partial_t (\varphi - \varphi_D), \mathcal{N}(\varphi - \varphi_D) \rangle_V = \frac{1}{2} \frac{\mathrm{d}}{\mathrm{d}t} \| \varphi - \varphi_D \|_*^2, \qquad \langle -\Delta \mu, \mathcal{N}(\varphi - \varphi_D) \rangle_V = (\mu, \varphi - \varphi_D)_H$$

so that, integrating in time,

$$\begin{split} &\frac{1}{2} \| (\varphi - \varphi_D)(t) \|_*^2 + \int_{Q_t} \mu(\varphi - \varphi_D) \\ &\lesssim \| \varphi_0 - (\varphi_0)_D \|_{V^*}^2 + \| \sigma \|_{L^2(0,t;H)}^2 + \| u \|_{L^2(0,t;V^*)}^2 + \| \varphi - \varphi_D \|_{L^2(0,t;V^*)}^2 + \| \varphi \|_{L^2(0,t;H)}^2 , \end{split}$$

where we used the boundedness of $h(\varphi_1), \sigma_1$, the Lipschitz continuity of h and the characterization (2.2). We test now equation (3.17) by $\varphi - \varphi_D$ and we employ (A2) to obtain

$$A \int_{Q_{t}} |\nabla \varphi|^{2} = \int_{Q_{t}} \mu(\varphi - \varphi_{D}) - B \int_{Q_{t}} (\psi'(\varphi_{1}) - \psi'(\varphi_{2})) (\varphi - \varphi_{D})$$

$$\leq \int_{Q_{t}} \mu(\varphi - \varphi_{D}) + B \int_{Q_{t}} \varphi_{D} (\psi'(\varphi_{1}) - \psi'(\varphi_{2})) + C_{2} \|\varphi\|^{2}_{L^{2}(0,t;H)}$$

$$\lesssim \int_{Q_{t}} \mu(\varphi - \varphi_{D}) + \int_{Q_{t}} \varphi_{D} |\varphi| (1 + |\psi''(\varphi_{1})| + |\psi''(\varphi_{2})|) + \|\varphi\|^{2}_{L^{2}(0,t;H)}.$$
(3.19)

From the previous estimates we have

$$\begin{split} \|(\varphi - \varphi_D)(t)\|_{V^*}^2 + |\varphi_D(t)|^2 + \int_{Q_t} |\nabla\varphi(s)|^2 \\ \lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \int_0^t \left(\varphi_D(s) \int_D |\varphi(s)| \left(1 + |\psi''(\varphi_1(s))| + |\psi''(\varphi_2(s))|\right)\right) ds \\ + \int_0^t \left(\|(\varphi(s) - \varphi_D(s))\|_{V^*}^2 + \|\sigma(s)\|_H^2 + \|u(s)\|_{V^*}^2 + \|\varphi(s) - \varphi_D(s)\|_H^2 + |\varphi_D(s)|^2\right) ds \\ \lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \int_0^t |\varphi_D(s)|^2 \int_D \left(1 + |\psi''(\varphi_1(s))|^2 + |\psi''(\varphi_2(s))|^2\right) ds \\ + \int_0^t \left(\|(\varphi(s) - \varphi_D(s))\|_{V^*}^2 + \|\sigma(s)\|_H^2 + \|u(s)\|_{V^*}^2 + \varepsilon \|\nabla\varphi(s)\|_H^2 + |\varphi_D(s)|^2\right) ds, \end{split}$$
(3.20)

where we used Young inequality and the compactness inequality (2.3) with $\varepsilon < 1$. We introduce now a sequence of stopping times $(\tau^N)_{N\in\mathbb{N}}$ as

$$\tau^{N} := \inf \left\{ t \ge 0 : \|\psi''(\varphi_{1})\|_{L^{2}(0,t;H)}^{2} + \|\psi''(\varphi_{2})\|_{L^{2}(0,t;H)}^{2} \ge N \right\} \wedge T.$$

Notice that $\tau^N = \tau^N(\varphi_1, \varphi_2)$ depends on φ_1 and φ_2 but we will simply write τ^N not to weight too much the notation. For every $N \in \mathbb{N}$, the application of Gronwall's lemma (recall that ε can be chosen arbitrarily small) in the time horizon $[0, \tau^N]$ gives, for \mathbb{P} -almost every $\omega \in \Omega$,

$$\begin{aligned} \|(\varphi - \varphi_D)(t)\|_{V^*}^2 + |\varphi_D(t)|^2 + \int_0^t \|\nabla\varphi\|_H^2 \,\mathrm{d}s \\ \lesssim \left(1 + e^{T+N}(T+N)\right) \left(\|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \|\sigma\|_{L^2(0,\tau^N;H)}^2 + \|u\|_{L^2(0,\tau^N;V^*)}^2 \right) \end{aligned}$$

for every $t \in [0, \tau^N(\omega)]$. Taking the power p/2, the supremum in time and expectation on both sides we deduce that

$$\mathbb{E} \sup_{t \in [0,\tau^{N}]} \|\varphi(t) - \varphi_{D}(t)\|_{V^{*}}^{p} + \mathbb{E} \sup_{t \in [0,\tau^{N}]} |\varphi_{D}(t)|^{p} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|\nabla\varphi(s)\|_{H}^{2} \,\mathrm{d}s \right)^{p/2} \\
\lesssim_{N} \mathbb{E} \|\varphi_{0} - (\varphi_{0})_{D}\|_{V^{*}}^{p} + \mathbb{E} |(\varphi_{0})_{D}|^{p} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|u(s)\|_{V^{*}}^{2} \,\mathrm{d}s \right)^{p/2} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|\sigma(s)\|_{H}^{2} \,\mathrm{d}s \right)^{p/2}.$$
(3.21)

For what concerns σ , it holds that

$$d\sigma - \Delta\sigma dt + c\sigma h(\varphi_1)dt + c\sigma_2 \left(h(\varphi_1) - h(\varphi_2)\right)dt + b(\sigma - w)dt = \left(\mathcal{H}(\sigma_1) - \mathcal{H}(\sigma_2)\right)dW_2$$

and the same strategy as in (3.7) (taking into account also the term w) gives

$$\mathbb{E} \sup_{t \in [0, \tau^{N}]} \|\sigma(t)\|_{H}^{p} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|\sigma(s)\|_{V}^{2} \, \mathrm{d}s \right)^{p/2} \\
\lesssim_{c, L_{h}, L_{\mathcal{H}}} \|\sigma_{0}\|_{L^{p}(\Omega; H)} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|w(s)\|_{H}^{2} \, \mathrm{d}s \right)^{p/2} + \mathbb{E} \left(\int_{0}^{\tau^{N}} \|\varphi(s)\|_{H}^{2} \, \mathrm{d}s \right)^{p/2}.$$
(3.22)

Note again that, for every $t \in [0, T]$

$$\begin{aligned} \|\varphi\|_{L^{p}(\Omega;L^{2}(0,t;H))}^{p} &\lesssim \|\varphi-\varphi_{D}\|_{L^{p}(\Omega;L^{2}(0,t;H))}^{2} + \|\varphi_{D}\|_{L^{p}(\Omega;L^{2}(0,t))}^{2} \\ &\lesssim \|\varphi-\varphi_{D}\|_{L^{p}(\Omega;L^{2}(0,t;V^{*}))}^{2} + \varepsilon \|\nabla\varphi\|_{L^{p}(\Omega;L^{2}(0,t;H))}^{2} + \|\varphi_{D}\|_{L^{p}(\Omega;L^{2}(0,t))}^{2} \\ \end{aligned}$$

Combining (3.22) and (3.21), choosing ε sufficiently small, and applying the stochastic Gronwall's Lemma (see [41, Lem. 29.1]), we end up with (2.13).

For what concerns uniqueness, if we take the same initial conditions and set u, w = 0 in (2.13), a consistency result holds: for every $N \in \mathbb{N}$

$$\|\varphi - \varphi_D\|_{C^0([0,t];V^*))} + \|\varphi_D\| + \|\nabla\varphi\|_{L^2(0,t;H)}^2 + \|\sigma\|_{C^0([0,t];H)\cap L^2(0,T;V)} = 0$$
(3.23)

for every $t \in [0, \tau^N(\omega)]$ and for \mathbb{P} -almost every $\omega \in \Omega$. Hence, and the validity of (3.23) can be extended to the stochastic interval $[0, \tau]$ where

$$\tau := \lim_{N \to +\infty} \tau^N = \sup_{n \in \mathbb{N}} \tau^N.$$

It remains to show that $\tau = T$ \mathbb{P} -almost surely: by contradiction, if $\mathbb{P}\{\tau < T\} > 0$, then by definition of τ we would have $\mathbb{P}\{\|\psi'(\varphi_1)\|_{L^2(\tau,T;H)} + \|\psi'(\varphi_2)\|_{L^2(\tau,T;H)} = +\infty\} > 0$, which clearly contradicts that $\psi'(\varphi_1), \psi'(\varphi_2) \in L^1(\Omega; L^2(0,T;H)).$

3.4 Stochastic compactness and passage to the limit

In this section we prove Theorem 2.2. From the uniform estimates (3.11)-(3.15) obtained above, there exists a constant M, independent of λ , such that

$$\begin{aligned} \|\sigma_{\lambda}\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;V))\cap L^{\infty}(\Omega\times Q)} &\leq M, \\ \|\varphi_{\lambda}\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{\infty}(0,T;V)\cap L^{2}(0,T;Z))} &\leq M, \\ \|\mu_{\lambda}\|_{L^{p/2}(\Omega;L^{2}(0,T;V))} + \|\nabla\mu_{\lambda}\|_{L^{p}(\Omega;L^{2}(0,T;H))} + \|(\mu_{\lambda})_{D}\|_{L^{p/2}(\Omega;L^{\infty}(0,T))} &\leq M, \\ \|\psi_{\lambda}'(\varphi_{\lambda})\|_{L^{p/2}(\Omega;L^{2}(0,T;H))} &\leq M. \end{aligned}$$

$$(3.24)$$

Moreover, given $s \in (0, 1/2)$, from Hypothesis (A3) and [23, Lem. 2.1]

$$G \cdot W_1 \in L^p(\Omega; W^{s,2}(0,T;V)) \cap L^{\kappa}(\Omega; W^{s,\kappa}(0,T);V^*)) \quad \forall \kappa \ge 2,$$

and by comparison in (2.10) it holds that

$$\|\varphi_{\lambda}\|_{L^{p}(\Omega; W^{s,\kappa}(0,T;V^{*}))} \leq M$$

This is crucial when dealing with the method of compactness: fixing $\kappa > \frac{1}{s}$, by [54, Cor. 4-5] the following inclusions are compact

$$\begin{split} L^{\infty}(0,T;V) \cap W^{s,\kappa}(0,T;V^*) &\hookrightarrow C^0([0,T];H) \,, \\ L^2(0,T;Z) \cap W^{s,2}(0,T;V^*) &\hookrightarrow L^2(0,T;V) \,. \end{split}$$

We are now in position to show that the family of laws $(\pi_{\lambda})_{\lambda} := (\mathscr{L}(\varphi_{\lambda}))_{\lambda}$ is a tight family of probability measures on $C([0,T];H) \cap L^2(0,T;V)$. Denoting $\mathcal{X} := L^{\infty}(0,T;V) \cap W^{s,\kappa}(0,T;V^*) \cap L^2(0,T;Z)$, we know that $\mathcal{X} \hookrightarrow C([0,T];H) \cap L^2(0,T;V)$ compactly and

$$\|\varphi_{\lambda}\|_{L^{p}(\Omega;\mathcal{X})} \leq M.$$

Hence, denoting by B_n the closed ball of radius n in \mathcal{X} , we have that B_n is compact in $C([0,T];H) \cap L^2(0,T;V)$ and the Markov inequality implies that

$$\sup_{\lambda>0} \pi_{\lambda}(B_n^c) = \sup_{\lambda>0} \mathbb{P}\{\|\varphi_{\lambda}\|_{\mathcal{X}}^p > n^p\} \le \frac{1}{n^p} \mathbb{E} \|\varphi_{\lambda}\|_{\mathcal{X}}^p \le \frac{M^p}{n^p} \to 0$$

as $n \to \infty$. We deduce that for every $\varepsilon > 0$ there is $n \in \mathbb{N}$ sufficiently large such that the compact set B_n of $C([0,T];H) \cap L^2(0,T;V)$ satisfies $\pi_{\lambda}(B_n) > 1 - \varepsilon$ for every $\lambda > 0$, and this proves the tightness of the laws $(\pi_{\lambda})_{\lambda}$ on $C([0,T];H) \cap L^2(0,T;V)$.

In order to reconstruct a process from the limit law, we need to exhibit a random variable φ : $(\Omega, \mathscr{F}, \mathbb{P}) \to C([0, T]; H) \cap L^2(0, T; V)$ such that for every $\delta > 0$

$$\mathbb{P}\{\|\varphi_{\lambda} - \varphi\|_{C([0,T];H) \cap L^2(0,T;V)} > \delta\} \longrightarrow 0 \quad \text{as } \lambda \to 0.$$

Thanks to a method of Gyöngy-Krylov [32, Lem. 1.1], this is equivalent to the following: for any subsequences $(\varphi_k)_k := (\varphi_{\lambda_k})_k$ and $(\varphi_j)_j := (\varphi_{\lambda_j})_j$ of the original sequence $(\varphi_{\lambda})_{\lambda}$, there exists a joint subsequence $(\varphi_{k_i}, \varphi_{j_i})_i$ converging in law to a probability measure ν on $(C([0, T]; H) \cap L^2(0, T; V))^2$ with the property

$$\nu\left(\left\{(g^1, g^2) \in (C([0, T]; H) \cap L^2(0, T; V))^2 : g^1 = g^2\right\}\right) = 1.$$
(3.25)

Let us now prove the validity of (3.25). Using the tightness of the laws $(\pi_{\lambda})_{\lambda}$ and Skorokhod theorem (see [36, Thm. 2.7]) we can find a new probability space $(\tilde{\Omega}, \tilde{\mathscr{F}}, \tilde{\mathbb{P}})$ along with a sequence of random variables $(\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i}) : \tilde{\Omega} \to (C([0, T]; H) \cap L^2(0, T; V))^2$ such that

$$(\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i}) \longrightarrow (\varphi^1, \varphi^2)$$
 in $(C([0, T]; H) \cap L^2(0, T; V))^2$, $\tilde{\mathbb{P}}$ -a.s. (3.26)

and $\mathscr{L}(\varphi_{k_i}, \varphi_{j_i}) = \mathscr{L}(\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i})$, for every $i \in \mathbb{R}$. Moreover, there exists a sequence of measurable maps $\Upsilon_i : (\tilde{\Omega}, \tilde{\mathscr{F}}) \to (\Omega, \mathscr{F})$ such that for every $i \in \mathbb{R}$ it holds that $\mathbb{P} = \tilde{\mathbb{P}} \circ (\Upsilon_i)^{-1}$, $(\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i}) := (\varphi_{k_i}, \varphi_{j_i}) \circ \Upsilon_i$ and (3.26) takes place (c.f. [57, Thm. 1.10.4, Add. 1.10.5]).

If we define $(\tilde{\sigma}_{k_i}, \tilde{\sigma}_{j_i}) := (\sigma_{k_i}, \sigma_{j_i}) \circ \Upsilon_i$ and $(\tilde{\mu}_{k_i}, \tilde{\mu}_{j_i}) := (\mu_{k_i}, \mu_{j_i}) \circ \Upsilon_i$, the uniform bounds in (3.24) still holds in the form

$$\begin{aligned} \|(\tilde{\sigma}_{k_{i}},\tilde{\sigma}_{j_{i}})\|_{(L^{p}(\tilde{\Omega};C^{0}([0,T];H)\cap L^{2}(0,T;V))\cap L^{\infty}(\Omega\times Q))^{2}} &\leq M, \\ \|(\tilde{\varphi}_{k_{i}},\tilde{\varphi}_{j_{i}})\|_{(L^{p}(\tilde{\Omega};C^{0}([0,T];H)\cap L^{\infty}(0,T;V)\cap L^{2}(0,T;Z)))^{2}} &\leq M, \\ \|(\tilde{\mu}_{k_{i}},\tilde{\mu}_{j_{i}})\|_{(L^{p/2}(\tilde{\Omega};L^{2}(0,T;V)))^{2}} + \|\nabla(\tilde{\mu}_{k_{i}},\tilde{\mu}_{j_{i}})\|_{(L^{p}(\tilde{\Omega};L^{2}(0,T;H)))^{2}} &\leq M, \\ \|(\psi_{\lambda}'(\tilde{\varphi}_{k_{i}}),\psi_{\lambda}'(\tilde{\varphi}_{j_{i}}))\|_{(L^{p/2}(\tilde{\Omega};L^{2}(0,T;H)))^{2}} &\leq M. \end{aligned}$$

This guarantees that, up to passing to subsequences, recalling the strong convergence (3.26) and employing the strong-weak closure of maximal monotone operators as well as the weak lower semicontinuity of the norms,

$$\begin{aligned} & (\tilde{\sigma}_{k_i}, \tilde{\sigma}_{j_i}) \rightharpoonup (\sigma^1, \sigma^2) & \text{ in } (L^p(\Omega; L^2(0, T; V)))^2, \\ & (\tilde{\sigma}_{k_i}, \tilde{\sigma}_{j_i}) \stackrel{*}{\rightharpoonup} (\sigma^1, \sigma^2) & \text{ in } (L^p(\tilde{\Omega}; L^1(0, T; H))^* \cap L^{\infty}(\tilde{\Omega} \times Q))^2, \\ & (\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i}) \rightarrow (\varphi^1, \varphi^2) & \text{ in } (L^\ell(\tilde{\Omega}; C^0([0, T]; H) \cap L^2(0, T; V)))^2 \quad \forall \ell \in [1, p), \\ & (\tilde{\varphi}_{k_i}, \tilde{\varphi}_{j_i}) \rightharpoonup (\varphi^1, \varphi^2) & \text{ in } (L^p(\tilde{\Omega}; L^2(0, T; Z)))^2, \\ & (\tilde{\mu}_{k_i}, \tilde{\mu}_{j_i}) \rightharpoonup (\mu^1, \mu^2) & \text{ in } (L^{p/2}(\tilde{\Omega}; L^2(0, T; V)))^2, \\ & (\psi'_{\lambda}(\tilde{\varphi}_{k_i}), \psi'_{\lambda}(\tilde{\varphi}_{j_i})) \rightharpoonup (\psi'(\varphi^1), \psi'(\varphi^2)) & \text{ in } (L^{p/2}(\tilde{\Omega}; L^2(0, T; H)))^2, \end{aligned}$$

where the limit objects belong to the same spaces of the respective sequences. Furthermore, for every $t \in [0, T]$ and $\tilde{\mathbb{P}}$ -a.s. the following holds in V^* ,

$$\partial_t (\tilde{\varphi}_{k_i} - \tilde{\varphi}_{j_i}) - \Delta(\tilde{\mu}_{k_i} - \tilde{\mu}_{j_i}) = \mathcal{P}(\tilde{\sigma}_{k_i} - \tilde{\sigma}_{j_i})h(\tilde{\varphi}_{k_i}) + (\mathcal{P}\tilde{\sigma}_{j_i} - a - \alpha u)(h(\tilde{\varphi}_{k_i}) - h(\tilde{\varphi}_{j_i}))$$
$$\tilde{\mu}_{k_i} - \tilde{\mu}_{j_i} = -\Delta(\tilde{\varphi}_{k_i} - \tilde{\varphi}_{j_i}) + \psi'(\tilde{\varphi}_{k_i}) - \psi'(\tilde{\varphi}_{j_i}),$$

and passing to the limit as $i \to +\infty$ we get

$$\partial_t(\varphi^1 - \varphi^2) - \Delta(\mu^1 - \mu^2) = \mathcal{P}(\sigma^1 - \sigma^2)h(\varphi^1) + (\mathcal{P}\sigma^1 - a - \alpha u)\left(h(\varphi^1) - h(\varphi^2)\right)$$
$$\mu^1 - \mu^2 = -\Delta(\varphi^1 - \varphi^2) + \psi'(\varphi^1) - \psi'(\varphi^2).$$

From the uniqueness of the solution proved in Section 3.3 (see (3.23), with $t \in [0,T]$) we deduce that $\varphi^1(t) = \varphi^2(t)$ for every $t \in [0,T]$, $\tilde{\mathbb{P}}$ -a.s. and this readily implies that

$$\nu\left(\left\{(g^1,g^2)\in (C^0([0,T];H)\cap L^2(0,T;V))^2:g^1=g^2\right\}\right)=\tilde{\mathbb{P}}(\varphi^1=\varphi^2)=1\,,$$

and (3.25) holds. From Gyöngy-Krylov's criterium, it holds that

$$\varphi_{\lambda} \to \varphi$$
 in $C^0([0,T];H) \cap L^2(0,T;V)$, in \mathbb{P} -measure on Ω .

Moreover, using the uniform bounds (3.24), the lower semicontinuity of the norms and the strong-weak closure of maximal monotone graphs we also have that

$$\begin{split} \sigma_{\lambda} &\stackrel{*}{\rightharpoonup} \sigma & \text{ in } L^{p}(\Omega; L^{2}(0,T;V)) \cap L^{\infty}(\Omega \times Q) \,, \\ \varphi_{\lambda} &\to \varphi & \text{ in } L^{\ell}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V)) \quad \forall \ell \in [1,p) \,, \\ \varphi_{\lambda} &\rightharpoonup \varphi & \text{ in } L^{p}(\Omega; L^{2}(0,T;Z)) \,, \\ \mu_{\lambda} &\rightharpoonup \mu & \text{ in } L^{p/2}(\Omega; L^{2}(0,T;V)) \,, \\ \psi_{\lambda}'(\varphi_{\lambda}) &\rightharpoonup \psi'(\varphi) & \text{ in } L^{p/2}(\Omega; L^{2}(0,T;H)) \,, \end{split}$$

where

$$\begin{split} \varphi \in L^p(\Omega; C^0([0,T];H) \cap L^\infty(0,T;V) \cap L^2(0,T;Z)) \,, \\ \sigma \in L^p(\Omega; C^0([0,T];H) \cap L^2(0,T;V)) \cap L^\infty(\Omega \times Q) \,, \\ \mu \in L^{p/2}(\Omega; L^2(0,T;V)) \,, \qquad \nabla \mu \in L^p(\Omega; L^2(0,T;H)), \qquad (\mu)_D \in L^{p/2}(\Omega; L^\infty(0,T)) \,, \\ \psi'(\varphi) \in L^{p/2}(\Omega; L^2(0,T;H)) \,. \end{split}$$

Let us show a further convergence for $(\sigma_{\lambda})_{\lambda}$. To do it, we exploit the continuous dependence on data given in (3.22). For every subsequences $\lambda_i, \lambda_j, i, j \in \mathbb{N}$ and every $\ell \in [1, p)$, it holds

$$0 \le \left\| \sigma_{\lambda_i} - \sigma_{\lambda_j} \right\|_{L^{\ell}(\Omega; C^0([0,T];H) \cap L^2(0,T;V))} \lesssim \left\| \varphi_{\lambda_i} - \varphi_{\lambda_j} \right\|_{L^{\ell}(\Omega; L^2(0,T;H))} \longrightarrow 0, \quad \text{as } i \to +\infty.$$

This implies that $\sigma_{\lambda} \to \sigma$ in $L^{\ell}(\Omega; C^{0}([0, T]; H) \cap L^{2}(0, T; V))$, which allows to pass to the limit in the stochastic integrals, getting by the Burkholder-Davis-Gundy inequality that

$$\mathcal{H}(\sigma_{\lambda}) \cdot W_2 \to \mathcal{H}(\sigma) \cdot W_2 \qquad \text{in } L^{\ell}(\Omega; C^0([0, T]; H)) \quad \forall \ell \in [1, p) \,.$$

Putting this information together and letting $\lambda \searrow 0$ in the approximated problem, it is now straightforward to see that (φ, μ, σ) is a solution to (2.4)–(2.12).

Finally, the uniqueness result proved in Proposition 2.13 (see estimate (3.23), with $t \in [0, T]$) implies that (φ, μ, σ) is the unique solution to (1.1)-(1.5). This concludes the proof of Theorem 2.2.

4 Refined well-posedness

This section is devoted to the proof of the refined well-posedness result contained in Theorem 2.4.

4.1 Refined existence

In order to prove the additional regularity (2.17)–(2.18), we show that under assumptions (2.14)–(2.16) some additional estimates hold on the approximated solutions $(\varphi_{\lambda}, \mu_{\lambda}, \sigma_{\lambda})$ obtained in the previous section.

Fourth estimate. Thanks to (2.16) we have that

$$\left\|\psi_{\lambda}'(\varphi_{\lambda})\right\|_{H}^{2} \lesssim 1 + \left\|\varphi_{\lambda}\right\|_{L^{6}(D)}^{6}$$

hence also, by the embedding $V \hookrightarrow L^6(D)$,

$$\left\|\psi_{\lambda}'(\varphi_{\lambda})\right\|_{L^{2}(0,T;H)} \lesssim 1 + \left\|\varphi_{\lambda}\right\|_{L^{\infty}(0,T;V)}^{3}.$$

Similarly, since $\nabla \psi'_{\lambda}(\varphi_{\lambda}) = \psi''_{\lambda}(\varphi_{\lambda}) \nabla \varphi_{\lambda}$, we have again by (2.16) that

$$\left\|\nabla\psi_{\lambda}'(\varphi_{\lambda})\right\|_{H}^{2} \lesssim \left\|\nabla\varphi_{\lambda}\right\|_{H}^{2} + \int_{D} |\varphi_{\lambda}|^{4} |\nabla\varphi_{\lambda}|^{2} \le \left\|\nabla\varphi_{\lambda}\right\|_{H}^{2} + \left\|\varphi_{\lambda}\right\|_{L^{6}(D)}^{4} \left\|\nabla\varphi_{\lambda}\right\|_{L^{6}(D)}^{2},$$

from which, thanks again to the embedding $V \hookrightarrow L^6(D)$,

$$\|\nabla \psi_{\lambda}'(\varphi_{\lambda})\|_{L^{2}(0,T;H)} \lesssim \|\varphi\|_{L^{2}(0,T;Z)} + \|\varphi_{\lambda}\|_{L^{\infty}(0,T;V)}^{2} \|\varphi_{\lambda}\|_{L^{2}(0,T;Z)}.$$

We infer that

$$\|\psi_{\lambda}'(\varphi_{\lambda})\|_{L^{2}(0,T;V)} \lesssim 1 + \|\varphi_{\lambda}\|_{L^{\infty}(0,T;V)}^{3} + \|\varphi_{\lambda}\|_{L^{\infty}(0,T;V)}^{2} \|\varphi_{\lambda}\|_{L^{2}(0,T;Z)},$$

hence from (3.13)-(3.14) we deduce that

$$\|\psi_{\lambda}'(\varphi_{\lambda})\|_{L^{p/3}(\Omega;L^2(0,T;V))} \le M$$

By comparison in (3.2) and estimate (3.15) we infer that

$$\left\|\Delta\varphi_{\lambda}\right\|_{L^{p/3}(\Omega;L^{2}(0,T;V))} \leq M,$$

which yields (2.17) by elliptic regularity.

Fifth estimate. Arguing as in the Second estimate, we easily get that

$$\begin{aligned} \|\varphi_{\lambda}(t)\|_{H}^{2} + \int_{Q_{t}} |\Delta\varphi_{\lambda}|^{2} \lesssim_{A,B,\mathcal{P},a,\alpha,C_{2}} 1 + \|\varphi_{0}\|_{H}^{2} + \|G\|_{L^{2}(0,T;\mathscr{L}^{2}(U_{1},H))}^{2} \\ + \int_{Q_{t}} |\varphi_{\lambda}|^{2} + \int_{0}^{t} (\varphi_{\lambda}(s),G(s))_{H} \, \mathrm{d}W_{1}(s) \end{aligned}$$

for every $t \in [0, T]$. Now, we add and subtract the term

$$\frac{r}{2} \int_0^t |(\varphi_\lambda(s), G(s))_H|^2 \,\mathrm{d}s$$

on the right-hand side, with r > 0 to be fixed later, getting

$$\begin{split} \|\varphi_{\lambda}(t)\|_{H}^{2} &+ \int_{Q_{t}} |\Delta\varphi_{\lambda}|^{2} \lesssim 1 + \|\varphi_{0}\|_{H}^{2} + \|G\|_{L^{2}(0,T;\mathscr{L}^{2}(U_{1},H))}^{2} + \int_{Q_{t}} |\varphi_{\lambda}|^{2} \\ &+ \left(\int_{0}^{t} (\varphi_{\lambda}(s), G(s))_{H} \, \mathrm{d}W_{1}(s) - \frac{r}{2} \int_{0}^{t} |(\varphi_{\lambda}(s), G(s))_{H}|^{2} \, \mathrm{d}s\right) + \frac{r}{2} \int_{0}^{t} |(\varphi_{\lambda}(s), G(s))_{H}|^{2} \, \mathrm{d}s \,, \end{split}$$

where

$$\int_0^t |(\varphi_\lambda(s), G(s))_H|^2 \,\mathrm{d}s \le \|G\|_{L^\infty(\Omega \times (0,T);\mathscr{L}^2(U_1,H))}^2 \int_{Q_t} |\varphi_\lambda|^2 \,.$$

Hence, we infer that

$$\begin{aligned} \|\varphi_{\lambda}(t)\|_{H}^{2} + \int_{Q_{t}} |\Delta\varphi_{\lambda}|^{2} \lesssim_{\|G\|_{L^{\infty}(\Omega\times(0,T;\mathscr{L}^{2}(U_{1},H))}^{2}} 1 + \|\varphi_{0}\|_{H}^{2} + \int_{Q_{t}} |\varphi_{\lambda}|^{2} \\ + \left(\int_{0}^{t} (\varphi_{\lambda}(s), G(s))_{H} \, \mathrm{d}W_{1}(s) - \frac{r}{2} \int_{0}^{t} |(\varphi_{\lambda}(s), G(s))_{H}|^{2} \, \mathrm{d}s\right) \,,\end{aligned}$$

and consequently, for every $\beta \geq 1$,

$$\begin{split} &\exp\left(\beta \left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right) + \exp\left(\beta \int_{Q_{t}} \left|\Delta\varphi_{\lambda}\right|^{2}\right) \\ &\lesssim \exp\left(\beta \left\|\varphi_{0}\right\|_{H}^{2}\right) \cdot \exp\left(\beta \int_{Q_{t}} \left|\varphi_{\lambda}\right|^{2}\right) \cdot \exp\left(\beta \int_{0}^{t} \left(\varphi_{\lambda}(s), G(s)\right)_{H} \, \mathrm{d}W_{1}(s) - \frac{\beta r}{2} \int_{0}^{t} \left|\left(\varphi_{\lambda}(s), G(s)\right)_{H}\right|^{2} \mathrm{d}s\right) \\ &\lesssim \exp\left(3\beta \left\|\varphi_{0}\right\|_{H}^{2}\right) + \exp\left(3\beta \int_{Q_{t}} \left|\varphi_{\lambda}\right|^{2}\right) \\ &+ \exp\left(3\beta \int_{0}^{t} \left(\varphi_{\lambda}(s), G(s)\right)_{H} \, \mathrm{d}W_{1}(s) - \frac{3\beta r}{2} \int_{0}^{t} \left|\left(\varphi_{\lambda}(s), G(s)\right)_{H}\right|^{2} \mathrm{d}s\right) \end{split}$$

for every $t \in [0, T]$, P-almost surely. Now, choosing $r := 3\beta$, it is well-know that the process

$$t \mapsto \exp\left(3\beta \int_0^t \left(\varphi_\lambda(s), G(s)\right)_H \, \mathrm{d}W_1(s) - \frac{9\beta^2}{2} \int_0^t \left|\left(\varphi_\lambda(s), G(s)\right)_H\right|^2 \mathrm{d}s\right)$$

is a real positive local martingale, hence also a real supermartingale, so that

$$\mathbb{E}\exp\left(3\beta\int_0^t \left(\varphi_\lambda(s), G(s)\right)_H \,\mathrm{d}W_1(s) - \frac{9\beta^2}{2}\int_0^t |(\varphi_\lambda(s), G(s))_H|^2 \,\mathrm{d}s\right) \le 1 \qquad \forall t \in [0, T] \,.$$

Consequently, taking expectations and supremum in time yields, for all $T_0 \in (0, T]$,

$$\begin{split} &\sup_{t\in[0,T_0]} \mathbb{E}\exp\left(\beta \left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right) + \mathbb{E}\exp\left(\beta \left\|\Delta\varphi_{\lambda}\right\|_{L^{2}(0,T_0;H)}^{2}\right) \lesssim 1 + \mathbb{E}\exp\left(3\beta \left\|\varphi_{0}\right\|_{H}^{2}\right) \\ &+ \mathbb{E}\exp\left(3\beta \int_{Q_{T_0}} |\varphi_{\lambda}|^{2}\right). \end{split}$$

Noting that, by the Jensen inequality and Fubini's theorem,

$$\begin{split} \mathbb{E}\exp\left(\frac{\beta}{T_{0}}\left\|\varphi_{\lambda}\right\|_{L^{2}(0,T_{0};H)}^{2}\right) &\leq \mathbb{E}\frac{1}{T_{0}}\int_{0}^{T_{0}}\exp\left(\beta\left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right)\,\mathrm{d}t = \frac{1}{T_{0}}\int_{0}^{T_{0}}\mathbb{E}\exp\left(\beta\left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right)\,\mathrm{d}t \\ &\leq \sup_{t\in[0,T_{0}]}\mathbb{E}\exp\left(\beta\left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right)\,,\end{split}$$

rearranging the terms yields

$$\mathbb{E} \exp\left(\frac{\beta}{T_0} \left\|\varphi_{\lambda}\right\|_{L^2(0,T_0;H)}^2\right) + \mathbb{E} \exp\left(\beta \left\|\Delta\varphi_{\lambda}\right\|_{L^2(0,T_0;H)}^2\right)$$
$$\lesssim 1 + \mathbb{E} \exp\left(3\beta \left\|\varphi_{0}\right\|_{H}^2\right) + \mathbb{E} \exp\left(3\beta \left\|\varphi_{\lambda}\right\|_{L^2(0,T_0;H)}^2\right).$$

Hence, choosing $T_0 \in (0,T]$ such that $\frac{\beta}{T_0} > 3\beta$, for example $T_0 = \frac{1}{6} \wedge T$, and using the Young inequality yields, for every $\varepsilon > 0$ and for $C_{\varepsilon} > 0$,

$$\begin{split} \mathbb{E} \exp\left(6\beta \|\varphi_{\lambda}\|_{L^{2}(0,T_{0};H)}^{2}\right) + \mathbb{E} \exp\left(\beta \|\Delta\varphi_{\lambda}\|_{L^{2}(0,T_{0};H)}^{2}\right) \\ \lesssim 1 + \mathbb{E} \exp\left(3\beta \|\varphi_{0}\|_{H}^{2}\right) + \mathbb{E} \exp\left(3\beta \|\varphi_{\lambda}\|_{L^{2}(0,T_{0};H)}^{2}\right) \\ \lesssim 1 + \mathbb{E} \exp\left(3\beta \|\varphi_{0}\|_{H}^{2}\right) + \varepsilon \mathbb{E} \exp\left(6\beta \|\varphi_{\lambda}\|_{L^{2}(0,T_{0};H)}^{2}\right) + C_{\varepsilon} \,. \end{split}$$

Choosing $\varepsilon > 0$ sufficiently small, rearranging the terms and using a standard patching argument implies that there exists $M_{\beta} > 0$, independent of λ , such that

$$\sup_{t\in[0,T]} \mathbb{E}\left(\beta \left\|\varphi_{\lambda}(t)\right\|_{H}^{2}\right) + \mathbb{E}\exp\left(\beta \left\|\varphi_{\lambda}\right\|_{L^{2}(0,T;Z)}^{2}\right) \leq M_{\beta}.$$
(4.1)

This concludes the proof of (2.18).

4.2 Refined continuous dependence

We prove the refined continuous dependence property (2.20) by using the additional regularity (2.18). We use the same notation of Section 3.3. From (3.19)–(3.20) and the mean-value theorem we get that

$$\begin{aligned} \|\varphi(t) - \varphi_D(t)\|_{V^*}^2 + |\varphi_D(t)|^2 + \int_0^t \|\nabla\varphi(s)\|_H^2 \, \mathrm{d}s \\ &\lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \int_0^t |\varphi_D(s)| \int_D (\psi'(\varphi_1(s)) - \psi'(\varphi_2(s))) \, \mathrm{d}s \\ &+ \int_0^t \left(\|\varphi(s) - \varphi_D(s)\|_{V^*}^2 + |\varphi_D(s)|^2 \right) \, \mathrm{d}s + \|\sigma\|_{L^2(0,t;H)}^2 + \|u\|_{L^2(0,t;V^*)}^2 \end{aligned}$$

$$&\lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \int_0^t |\varphi_D(s)| \int_D \varphi(s) \int_0^1 \psi''(\iota\varphi_1(s) + (1-\iota)\varphi_2(s)) \, \mathrm{d}\iota \, \mathrm{d}s \\ &+ \int_0^t \left(\|(\varphi(s) - \varphi_D(s))\|_{V^*}^2 + |\varphi_D(s)|^2 \right) \, \mathrm{d}s + \|\sigma\|_{L^2(0,t;H)}^2 + \|u\|_{L^2(0,t;V^*)}^2 . \end{aligned}$$

$$(4.2)$$

Note that, by the growth assumption (2.16), for every $\iota \in [0, 1]$ we have that

$$|\psi''(\iota\varphi_1 + (1-\iota)\varphi_2)| \lesssim 1 + |\varphi_1|^2 + |\varphi_2|^2$$

and

$$\begin{split} |\nabla\psi''(\iota\varphi_1 + (1-\iota)\varphi_2)| &= |\psi'''(\iota\varphi_1 + (1-\iota)\varphi_2)(\iota\nabla\varphi_1 + (1-\iota)\nabla\varphi_2)| \\ &\lesssim (1+|\varphi_1| + |\varphi_2|)(|\nabla\varphi_1| + |\nabla\varphi_2|) \\ &\lesssim 1+|\varphi_1|^2 + |\varphi_2|^2 + |\nabla\varphi_1|^2 + |\nabla\varphi_2|^2 \,, \end{split}$$

from which we deduce, thanks to the continuous embedding $V \hookrightarrow L^6(D)$, that $\|\psi''(\iota\varphi_1 + (1-\iota)\varphi_2)\|_V^2 \lesssim 1 + \|\varphi_1\|_{L^4(D)}^4 + \|\varphi_2\|_{L^4(D)}^4 + \|\nabla\varphi_1\|_{L^4(D)}^4 + \|\nabla\varphi_2\|_{L^4(D)}^4 \lesssim 1 + \|\varphi_1\|_Z^4 + \|\varphi_2\|_Z^4 .$ Hence, we infer that, for every $s \in [0,T]$,

$$\left\| \int_0^1 \psi''(\iota\varphi_1(s) + (1-\iota)\varphi_2(s)) \,\mathrm{d}\iota \right\|_V \lesssim 1 + \|\varphi_1(s)\|_Z^2 + \|\varphi_2(s)\|_Z^2 \,.$$

Hence, substituting in (4.2) we have

$$\begin{aligned} \|\varphi(t) - \varphi_D(t)\|_{V^*}^2 + |\varphi_D(t)|^2 + \int_0^t \|\nabla\varphi(s)\|_H^2 \,\mathrm{d}s \\ &\lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \int_0^t \|\varphi(s)\|_{V^*}^2 \left(1 + \|\varphi_1(s)\|_Z^2 + \|\varphi_2(s)\|_Z^2\right) \,\mathrm{d}s \\ &+ \int_0^t \left(\|(\varphi(s) - \varphi_D(s))\|_{V^*}^2 + |\varphi_D(s)|^2\right) \,\mathrm{d}s + \|\sigma\|_{L^2(0,t;H)}^2 + \|u\|_{L^2(0,t;V^*)}^2. \end{aligned}$$

Now, since (2.19) is in order, we have that

$$\partial_t \sigma - \Delta \sigma + c \sigma h(\varphi_1) + c \sigma_2 \left(h(\varphi_1) - h(\varphi_2) \right) + b(\sigma - w) = 0, \qquad \sigma(0) = \sigma_0$$

so that it is not difficult to prove that, for every $\varepsilon>0$ and a certain $C_{\varepsilon}>0$

$$\begin{aligned} \|\sigma\|_{C^{0}([0,t];H)\cap L^{2}(0,t;V))}^{2} \lesssim_{c,L_{h}} \|\sigma_{0}\|_{H}^{2} + \|\varphi\|_{L^{2}(0,t;H)}^{2} + \|w\|_{L^{2}(0,t;H)}^{2} \\ & \leq \|\sigma_{0}\|_{H}^{2} + \varepsilon \|\nabla\varphi\|_{L^{2}(0,t;H)}^{2} + C_{\varepsilon} \|\varphi\|_{L^{2}(0,t;V^{*})}^{2} + \|w\|_{L^{2}(0,t;H)}^{2} \qquad (4.3) \end{aligned}$$

Substituting in the estimate above and choosing ε small enough we get that

$$\begin{split} \|\varphi(t) - \varphi_D(t)\|_{V^*}^2 + |\varphi_D(t)|^2 + \int_0^t \|\nabla\varphi(s)\|_H^2 \,\mathrm{d}s \\ \lesssim \|\varphi_0 - (\varphi_0)_D\|_{V^*}^2 + |(\varphi_0)_D|^2 + \|\sigma_0\|_H^2 + \|u\|_{L^2(0,T;V^*)}^2 + \|w\|_{L^2(0,T;H)}^2 \\ + \int_0^t \|\varphi(s)\|_{V^*}^2 \left(1 + \|\varphi_1(s)\|_Z^2 + \|\varphi_2(s)\|_Z^2\right) \,\mathrm{d}s + \int_0^t \|(\varphi(s) - \varphi_D(s))\|_{V^*}^2 \,\mathrm{d}s \end{split}$$

for every $t \in [0, T]$, P-almost surely. The Gronwall lemma implies then that

$$\begin{aligned} \|\varphi\|_{C^{0}([0,T];V^{*})\cap L^{2}(0,T;V)}^{2} \lesssim \left(\|\varphi_{0}\|_{V^{*}}^{2} + \|\sigma_{0}\|_{H}^{2} + \|u\|_{L^{2}(0,T;V^{*})}^{2} + \|w\|_{L^{2}(0,T;H)}^{2}\right) \\ \times \exp\left(\|\varphi_{1}\|_{L^{2}(0,T;Z)}^{2} + \|\varphi_{2}\|_{L^{2}(0,T;Z)}^{2}\right) \end{aligned}$$
(4.4)

so that, taking p/2-power, expectations, applying the Hölder inequality on the right-hand side with exponents $\frac{q}{p} > 1$ and $\frac{q/p}{(q/p)-1}$, and recalling condition (2.18) holds with the choice $\beta = \frac{q/p}{(q/p)-1}$, we infer that

$$\|\varphi\|_{L^{p}(\Omega;C^{0}([0,T];V^{*})\cap L^{2}(0,T;V))}^{p} \lesssim \|\varphi_{0}\|_{L^{q}(\Omega;V^{*})}^{p} + \|\sigma_{0}\|_{L^{q}(\Omega;H)}^{p} + \|u\|_{L^{q}(\Omega;L^{2}(0,T;V^{*}))}^{p} + \|w\|_{L^{q}(\Omega;L^{2}(0,T;H))}^{p} ,$$

so that (2.20) follows from (4.3).

Now, testing (3.16) by φ , (3.17) by $-\Delta\varphi$ and taking the difference, yields

$$\frac{1}{2} \|\varphi(t)\|_{H}^{2} + \int_{Q_{t}} |\Delta\varphi|^{2} - \int_{Q_{t}} \left(\psi'(\varphi_{1}) - \psi'(\varphi_{2})\right) \Delta\varphi$$
$$= \frac{1}{2} \|\varphi_{0}\|_{H}^{2} + \int_{Q_{t}} \left(\mathcal{P}\sigma - \alpha u\right) h(\varphi_{1})\varphi + \int_{Q_{t}} \left(\mathcal{P}\sigma_{2} - a - \alpha u_{2}\right) \left(h(\varphi_{1}) - h(\varphi_{2})\right)\varphi$$

for every $t \in [0, T]$. Using the boundedness of σ_2 , u_2 , h and the Lipschitz-continuity of h we get

$$\|\varphi\|_{C^{0}([0,T];H)\cap L^{2}(0,T;Z)}^{2} \lesssim \|\varphi_{0}\|_{H}^{2} + \|\varphi\|_{L^{2}(0,T;H)}^{2} + \|\sigma\|_{L^{2}(0,T;H)}^{2} + \|u\|_{L^{2}(0,T;H)}^{2} + \int_{Q} \left(1 + |\varphi_{1}|^{4} + |\varphi_{2}|^{4}\right) |\varphi|^{2} .$$

Noting now that, by the Hölder inequality and the embedding $V \hookrightarrow L^6(D)$,

$$\begin{split} \int_{Q} \left(1 + |\varphi_{1}|^{4} + |\varphi_{2}|^{4} \right) |\varphi|^{2} &\lesssim \int_{0}^{T} \|\varphi(s)\|_{L^{6}(D)}^{2} \left(1 + \|\varphi_{1}(s)\|_{L^{6}(D)}^{4} + \|\varphi_{2}(s)\|_{L^{6}(D)}^{4} \right) \,\mathrm{d}s \\ &\lesssim \|\varphi\|_{L^{2}(0,T;V)}^{2} \left(1 + \|\varphi_{1}\|_{L^{\infty}(0,T;V)}^{4} + \|\varphi_{2}\|_{L^{\infty}(0,T;V)}^{4} \right), \end{split}$$

taking power p/2 and using (4.4) again to the power p/2 we infer that

$$\begin{split} \|\varphi\|_{C^{0}([0,T];H)\cap L^{2}(0,T;Z)}^{p} \lesssim \left(\|\varphi_{0}\|_{H}^{p} + \|\sigma_{0}\|_{H}^{p} + \|u\|_{L^{2}(0,T;H)}^{p} + \|w\|_{L^{2}(0,T;H)}^{p}\right) \\ \times \left(1 + \|\varphi_{1}\|_{L^{\infty}(0,T;V)}^{2p} + \|\varphi_{2}\|_{L^{\infty}(0,T;V)}^{2p}\right) \\ \times \exp\left(\frac{p}{2}\|\varphi_{1}\|_{L^{2}(0,T;Z)}^{2} + \frac{p}{2}\|\varphi_{2}\|_{L^{2}(0,T;Z)}^{2}\right). \end{split}$$

Now, note that by (2.21) it is easy to check that

$$1 - \frac{1}{\beta_0} := \frac{p}{q} + \frac{2p}{r} < 1 \,,$$

hence we can take expectations and use the Hölder inequality on the right-hand side with exponents q/p, r/(2p) and β_0 , respectively, getting

$$\begin{split} \|\varphi\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;Z))}^{p} \lesssim \left(\|\varphi_{0}\|_{L^{q}(\Omega;V^{*})}^{p} + \|\sigma_{0}\|_{L^{q}(\Omega;H)}^{p} + \|u\|_{L^{q}(\Omega;L^{2}(0,T;H))}^{p} + \|w\|_{L^{q}(\Omega;L^{2}(0,T;H))}^{p}\right) \\ \times \left(1 + \|\varphi_{1}\|_{L^{r}(\Omega;L^{\infty}(0,T;V))}^{2p} + \|\varphi_{2}\|_{L^{r}(\Omega;L^{\infty}(0,T;V))}^{2p}\right) \\ \times \left\|\exp\left(\frac{\beta_{0}p}{2}\|\varphi_{1}\|_{L^{2}(0,T;Z)}^{2} + \frac{\beta_{0}p}{2}\|\varphi_{2}\|_{L^{2}(0,T;Z)}^{2}\right)\right\|_{L^{1}(\Omega)}^{1/\beta_{0}}. \end{split}$$

Noting that $\varphi_1, \varphi_2 \in L^r(\Omega; L^{\infty}(0, T; V))$ by (2.21), (2.4), and Theorem 2.2, recalling also (2.18) the last two factors are finite, and we can conclude the proof of Theorem 2.4.

5 Optimal control problem

This section is devoted to the analysis of the optimal control problem associated to the state system (1.1)-(1.5) and the cost functional J.

5.1 Existence of an optimal control

We prove here Theorem 2.7, showing that a relaxed optimal control always exists.

Let $(u_n, w_n)_n \subset \mathcal{U}$ be a minimizing sequence for \tilde{J} in \mathcal{U} , i.e. such that

$$\lim_{n \to \infty} \tilde{J}(u_n, w_n) = \inf_{(v,z) \in \mathcal{U}} \tilde{J}(v, z) \,.$$

By definition of \mathcal{U} the sequence $(u_n, w_n)_n$ is uniformly bounded in $L^{\infty}(\Omega \times Q)$. In particular, if we denote by $L^2_w(Q)$ the space $L^2(Q)$ equipped with its weak topology, it is immediate to see that the sequence of laws of (u_n, w_n) is tight on $L^2_w(Q)^2$. Furthermore, for every $n \in \mathbb{N}$, let $(\varphi_n, \mu_n, \sigma_n)_n$ be the corresponding solution to (2.4)–(2.12) with respect to the data $(\varphi_0, \sigma_0, u_n, w_n)$: recalling the proof of Theorem 2.2 (see section 3), we know that there exists a positive constant M, depending only on the initial data (φ_0, σ_0) , but not on n, such that

$$\begin{split} \|\varphi_n\|_{L^p(\Omega; W^{s,\kappa}(0,T;V^*)\cap C^0([0,T];H)\cap L^{\infty}(0,T;V)\cap L^2(0,T;Z))} &\leq M \,, \\ \|\mu_n\|_{L^{p/2}(\Omega; L^2(0,T;H))} + \|\nabla\mu_n\|_{L^p(\Omega; L^2(0,T;H))} + \|(\mu_n)_D\|_{L^{p/2}(\Omega; L^{\infty}(0,T))} &\leq M \,, \\ \|\psi'(\varphi_n)\|_{L^{p/2}(\Omega; L^2(0,T;H))} &\leq M \,, \\ \|\sigma_n\|_{L^p(\Omega; H^1(0,T;V^*)\cap L^2(0,T;V))\cap L^{\infty}(\Omega \times Q)} &\leq M \,, \end{split}$$

where $s \in (0, 1/2)$ and $\kappa > 1/s$ are fixed. In particular, the sequence of laws of $(\varphi_n)_n$ and $(\sigma_n)_n$ are tight on the spaces $L^2(0,T;V) \cap C^0([0,T];H)$ and $L^2(0,T;H)$, respectively (see again section 3).

Hence, the sequence $(W_1, G \cdot W_1, \varphi_0, \sigma_0, u_n, w_n, \varphi_n, \mu_n, \sigma_n, \varphi_Q, \varphi_T)_n$ is tight on the product space

$$C^{0}([0,T];U_{1}) \times C^{0}([0,T];H) \times V \times H \times L^{2}_{w}(Q)^{2} \times C^{0}([0,T];H) \times L^{2}_{w}(0,T;V) \times L^{2}(0,T;H) \times L^{2}(0,T;H) \times V.$$

Recalling that $L^2_w(Q)$ and $L^2_w(0,T;V)$ are a sub-Polish spaces, by Jakubowski-Skorokhod theorem (see e. g. [4, Thm. 2.7.1]) there is a probability space $(\Omega', \mathscr{F}', \mathbb{P}')$ and a sequence of measurable mappings $\phi_n:(\Omega',\mathscr{F}')\to(\Omega,\mathscr{F})$ such that $\mathbb{P}:=\mathbb{P}'\circ\phi_n^{-1}$ for every $n\in\mathbb{N}$ and

$$\begin{split} (W'_{1,n},I'_n) &:= (W_1,G\cdot W_1) \circ \phi_n \to (W'_1,I') & \text{ in } C^0([0,T];U_1) \times C^0([0,T];H) \\ (\varphi_{0,n},\sigma_{0,n})' &:= (\varphi_0,\sigma_0) \circ \phi_n \to (\varphi'_0,\sigma'_0) & \text{ in } V \times H , \\ (u'_n,w'_n) &:= (u_n,w_n) \circ \phi_n \to (u',w') & \text{ in } L^2(0,T;H)^2 , \\ (\varphi'_n,\sigma'_n) &:= (\varphi_n,\sigma_n) \circ \phi_n \to (\varphi',\sigma') & \text{ in } C^0([0,T];H) \times L^2(0,T;H) , \\ \mu'_n &:= \mu_n \circ \phi_n \to \mu' & \text{ in } L^2(0,T;V) , \\ (\varphi'_{O,n},\varphi'_{T,n}) &:= (\varphi_Q,\varphi_T) \circ \phi_n \to (\varphi'_Q,\varphi'_T) & \text{ in } L^2(0,T;H) \times V , \end{split}$$

 \mathbb{P}' -almost surely on Ω' . Since $(W_1, G \cdot W_1, \varphi_0, \sigma_0, \varphi_Q, \varphi_T)_n$ is constant, it follows immediately that the law of $(W'_1, I', \varphi'_0, \sigma'_0, \varphi'_Q, \varphi'_T)$ coincides with the law of $(W_1, \varphi_0, \sigma_0, \varphi_Q, \varphi_T)$. Moreover, by weak lower semicontinuity we also have that $0 \leq u', w' \leq 1$ almost everywhere in $\Omega' \times Q$. Finally, using a classical procedure based on martingale representation theorems (for a detailed argument the reader can refer to [58, § 4]), it is possible to show that $W'_{1,n}$ is a $(\mathscr{F}'_n, t)_t$ -cylindrical Wiener process in U_1 and W'_1 is a $(\mathscr{F}'_t)_t$ -cylindrical Wiener process in U_1 , where

$$\mathscr{F}_{n,t}' := \sigma(W_{1,n}'(s))_{s \in [0,t]} \,, \qquad \mathscr{F}_t' := \sigma(W_1'(s), I'(s), \varphi(s), \sigma(s))_{s \in [0,t]} \,,$$

and that $I' = G \cdot W'$. Since the uniform estimates on $(\varphi_n, \mu_n, \sigma_n, \psi'(\varphi_n))_n$ are also satisfied by $(\varphi'_n, \mu'_n, \sigma'_n, \psi'(\varphi'_n))_n$, passing to the weak limit as $n \to \infty$ in the variational formulation of the problem on $(\Omega', \mathscr{F}', \mathbb{P}')$, by the strong-weak closure of maximal monotone operators it follows that $(\varphi', \mu', \sigma')$ is the unique solution to (2.4)–(2.12) on Ω' with respect to $(\varphi'_0, \sigma'_0, u', w')$. Consequently, by weak lower semicontinuity, the properties of $(\phi_n)_n$, and the definition of minimizing sequence, we have

$$\begin{split} \tilde{J}'(u',w') &\leq \liminf_{n \to \infty} \frac{\beta_1}{2} \mathbb{E}' \int_Q |\varphi'_n - \varphi'_{Q,n}|^2 + \frac{\beta_2}{2} \mathbb{E}' \int_D |\varphi'_n(T) - \varphi'_{T,n}|^2 + \frac{\beta_3}{2} \mathbb{E}' \int_D (\varphi'_n(T) + 1) \\ &+ \frac{\beta_4}{2} \mathbb{E}' \int_Q |u'_n|^2 + \frac{\beta_5}{2} \mathbb{E}' \int_Q |w'_n|^2 \\ &= \liminf_{n \to \infty} J(\varphi_n, u_n, w_n) = \liminf_{n \to \infty} \tilde{J}(u_n, w_n) = \inf_{(v,z) \in \mathcal{U}} \tilde{J}(v, z) \,, \end{split}$$

so that (u', w') is a relaxed optimal control.

5.2 The linearized system

In this section we prove Theorems 2.8–2.9. First of all, we show that uniqueness of solution holds for the linearized system (2.23)–(2.27). Secondly, we prove Theorem 2.9 and (hence) existence of solutions for the linearized system.

Uniqueness. Let us show that the linearized system (2.23)–(2.27) admits a unique solution. Let (x_k^i, y_k^i, z_k^i) solve (2.23)–(2.27) for i = 1, 2. Then we have

$$\begin{split} \partial_t (x_k^1 - x_k^2) &- \Delta (y_k^1 - y_k^2) = h(\varphi) \mathcal{P}(z_k^1 - z_k^2) + h'(\varphi)(x_k^1 - x_k^2) (\mathcal{P}\sigma - a - \alpha u) & \text{ in } (0, T) \times D \,, \\ y_k^1 - y_k^2 &= -A\Delta (x_k^1 - x_k^2) + B\psi''(\varphi)(x_k^1 - x_k^2) & \text{ in } (0, T) \times D \,, \\ \partial_t (z_k^1 - z_k^2) - \Delta (z_k^1 - z_k^2) + c(z_k^1 - z_k^2)h(\varphi) + c\sigma h'(\varphi)(x_k^1 - x_k^2) + b(z_k^1 - z_k^2) = 0 & \text{ in } (0, T) \times D \,, \\ \partial_n (x_k^1 - x_k^2) &= \partial_n (z_k^1 - z_k^2) = 0 & \text{ in } (0, T) \times \partial D \,, \\ (x_k^1 - x_k^2)(0) &= (z_k^1 - z_k^2)(0) = 0 & \text{ in } D \,. \end{split}$$

Testing the first equation by $\frac{1}{|D|}$, using the boundedness of h, h' and σ , we deduce that there exists M > 0 such that

$$\left\| (x_k^1 - x_k^2)_D \right\|_{C^0([0,t])}^2 \le M \left(\left\| z_k^1 - z_k^2 \right\|_{L^1(Q_t)}^2 + \left\| x_k^1 - x_k^2 \right\|_{L^1(Q_t)}^2 \right) \qquad \forall t \in [0,T] \,, \quad \mathbb{P}\text{-a.s.}$$

Testing the first equation by $\mathcal{N}((x_k^1 - x_k^2) - (x_k^1 - x_k^2)_D)$, the second one by $x_k^1 - x_k^2$ and taking the

difference yields

$$\begin{split} &\frac{1}{2} \left\| (x_k^1 - x_k^2 - (x_k^1 - x_k^2)_D)(t) \right\|_{V^*}^2 + A \int_{Q_t} |\nabla(x_k^1 - x_k^2)|^2 + B \int_{Q_t} \psi''(\varphi) |(x_k^1 - x_k^2)|^2 \\ &= B \int_{Q_t} \psi''(\varphi) (x_k^1 - x_k^2) (x_k^1 - x_k^2)_D \\ &+ \int_{Q_t} \left[h(\varphi) \mathcal{P}(z_k^1 - z_k^2) + h'(\varphi) (x_k^1 - x_k^2) (\mathcal{P}\sigma - a - \alpha u) \right] \mathcal{N}(x_k^1 - x_k^2 - (x_k^1 - x_k^2)_D) \,. \end{split}$$

Summing the two inequalities and recalling that $\psi'' \geq -C_2$ we infer then that

$$\begin{split} \left\| (x_k^1 - x_k^2)(t) \right\|_{V^*}^2 + \int_{Q_t} |\nabla(x_k^1 - x_k^2)|^2 \lesssim_{M, C_2, \mathcal{P}} \int_{Q_t} |x_k^1 - x_k^2|^2 + \int_{Q_t} |z_k^1 - z_k^2|^2 \\ + \int_0^t (x_k^1 - x_k^2)_D(s) \left\| (x_k^1 - x_k^2)(s) \right\|_{V^*} \left\| \psi''(\varphi(s)) \right\|_V \, \mathrm{d}s + \int_0^t \left\| (x_k^1 - x_k^2)(s) \right\|_{V^*}^2 \, \mathrm{d}s \, , \end{split}$$

where a direct computation based on (2.16), the embedding $V \hookrightarrow L^6(D)$ and the Young inequality yields (as already performed in Section 4.2)

$$\left\|\psi''(\varphi)\right\|_V \lesssim 1 + \left\|\varphi\right\|_Z^2$$

Testing the third equation by $z_k^1 - z_k^2$ it follows easily by the Gronwall lemma and the boundedness of h, h' and σ that

$$\left\|z_{k}^{1}-z_{k}^{2}\right\|_{C^{0}([0,t];H)\cap L^{2}(0,t;V)} \leq M\left\|x_{k}^{1}-x_{k}^{2}\right\|_{L^{2}(0,t;H)}$$

Hence, substituting in the previous inequality and using a compactness inequality in the form

$$\left\|x_{k}^{1}-x_{k}^{2}\right\|_{L^{2}(0,t;H)}^{2} \leq \varepsilon \left\|\nabla(x_{k}^{1}-x_{k}^{2})\right\|_{L^{2}(0,t;H)}^{2} + C_{\varepsilon} \left\|x_{k}^{1}-x_{k}^{2}\right\|_{L^{2}(0,t;V^{*})}^{2},$$

choosing ε sufficiently small and rearranging the terms we have

$$\left\| (x_k^1 - x_k^2)(t) \right\|_{V^*}^2 + \int_{Q_t} |\nabla(x_k^1 - x_k^2)|^2 \lesssim \int_0^t \left(1 + \|\varphi(s)\|_Z^2 \right) \left\| (x_k^1 - x_k^2)(s) \right\|_{V^*}^2 \, \mathrm{d}s \,,$$

yielding $x_k^1(t) = x_k^2(t)$ for every $t \in [0, T]$ thanks to the Gronwall lemma and recalling $\varphi \in L^2(0, T; Z)$. It follows as a consequence the uniqueness $y_k^1 = y_k^2$ and $z_k^1 = z_k^2$.

Gâteaux-differentiability and existence. Let $(u, w), (k_u, k_w) \in \tilde{\mathcal{U}}$ and let us set $\varphi := \mathcal{S}(u, w)$. Let now $\varepsilon \in (-\varepsilon_0, \varepsilon_0)$ where $\varepsilon_0 > 0$ is chosen sufficiently small so that $(u, w) + \varepsilon(k_u, k_w) \in \tilde{\mathcal{U}}$ for all $\varepsilon \in (-\varepsilon_0, \varepsilon_0)$ (note that this is possibly since $\tilde{\mathcal{U}}$ is an open subset of $L^2(\Omega \times Q)$). Defining $\varphi_{\varepsilon} := \mathcal{S}((u, w) + \varepsilon(k_u, k_w))$, we have then that

$$\partial_t \left(\frac{\varphi_{\varepsilon} - \varphi}{\varepsilon}\right) - \Delta \left(\frac{\mu_{\varepsilon} - \mu}{\varepsilon}\right) = \frac{h(\varphi_{\varepsilon}) - h(\varphi)}{\varepsilon} (\mathcal{P}\sigma_{\varepsilon} - a - u - \varepsilon k_u) + h(\varphi) \left(\mathcal{P}\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} - k_u\right),$$
$$\frac{\mu_{\varepsilon} - \mu}{\varepsilon} = -A\Delta \left(\frac{\varphi_{\varepsilon} - \varphi}{\varepsilon}\right) + B\frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon},$$
$$\partial_t \left(\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon}\right) - \Delta \left(\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon}\right) + c\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon}h(\varphi_{\varepsilon}) + c\frac{h(\varphi_{\varepsilon}) - h(\varphi)}{\varepsilon}\sigma + b\left(\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} - k_w\right) = 0,$$

where $\frac{\varphi_{\varepsilon}-\varphi}{\varepsilon}(0) = \frac{\sigma_{\varepsilon}-\sigma}{\varepsilon}(0) = 0$. Now, from the continuous dependence property (2.22), we have

$$\begin{aligned} \|\varphi_{\varepsilon} - \varphi\|_{L^{p}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;Z))} + \|\sigma_{\varepsilon} - \sigma\|_{L^{p}(\Omega; C^{0}([0,T];H) \cap L^{2}(0,T;V))} \\ &\leq M\varepsilon \Big(\|k_{u}\|_{L^{q}(\Omega; L^{2}(0,T;H))} + \|k_{w}\|_{L^{q}(\Omega; L^{2}(0,T;H))}\Big), \end{aligned}$$

so that we deduce the uniform estimate (updating M)

$$\left\|\frac{\varphi_{\varepsilon}-\varphi}{\varepsilon}\right\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;Z))}+\left\|\frac{\sigma_{\varepsilon}-\sigma}{\varepsilon}\right\|_{L^{p}(\Omega;C^{0}([0,T];H)\cap L^{2}(0,T;V))}\leq M.$$
(5.5)

Moreover, thanks to the growth condition (2.16), the Hölder inequality and the embedding $V \hookrightarrow L^6(D)$, we also have

$$\begin{split} &\int_{Q} \left| \frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon} \right|^{2} = \int_{Q} \int_{0}^{1} |\psi''(\varphi + \tau(\varphi_{\varepsilon} - \varphi))|^{2} \left| \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \right|^{2} d\tau \lesssim \int_{Q} (1 + |\varphi|^{4} + |\varphi_{\varepsilon}|^{4}) \left| \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \right|^{2} \\ &\leq \int_{0}^{T} \left(1 + \|\varphi(s)\|_{L^{6}(D)}^{4} + \|\varphi_{\varepsilon}(s)\|_{L^{6}(D)}^{4} \right) \left\| \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} (s) \right\|_{L^{6}(D)}^{2} ds \\ &\lesssim \left(1 + \|\varphi\|_{L^{\infty}(0,T;V)}^{4} + \|\varphi_{\varepsilon}\|_{L^{\infty}(0,T;V)}^{4} \right) \left\| \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \right\|_{L^{2}(0,T;V)}^{2}, \end{split}$$

yielding

$$\left\|\frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon}\right\|_{L^{2}(0,T;H)} \lesssim \left(1 + \|\varphi\|_{L^{\infty}(0,T;V)}^{2} + \|\varphi_{\varepsilon}\|_{L^{\infty}(0,T;V)}^{2}\right) \left\|\frac{\varphi_{\varepsilon} - \varphi}{\varepsilon}\right\|_{L^{2}(0,T;V)}$$

where by (4.4) we have that

$$\left\|\frac{\varphi_{\varepsilon}-\varphi}{\varepsilon}\right\|_{L^{2}(0,T;V)} \lesssim \left(\|k_{u}\|_{L^{2}(0,T;V^{*})} + \|k_{w}\|_{L^{2}(0,T;H)}\right) \exp\left(\|\varphi\|_{L^{2}(0,T;Z)}^{2} + \|\varphi_{\varepsilon}\|_{L^{2}(0,T;Z)}^{2}\right)$$

Now, thanks to (2.21) we have that $\|\varphi\|_{L^{\infty}(0,T;V)}^{2} + \|\varphi_{\varepsilon}\|_{L^{\infty}(0,T;V)}^{2}$ is uniformly bounded (w.r.t. ε) in $L^{r/2}(\Omega)$, where $\frac{r}{2} > \frac{pq}{q-p} > p$. Moreover, $\|k_{u}\|_{L^{2}(0,T;V^{*})} + \|k_{w}\|_{L^{2}(0,T;H)} \in L^{q}(\Omega)$ by definition of $\tilde{\mathcal{U}}$, where q > p by assumption, and by (2.14)–(2.16) we also have

$$\left\| \exp\left(\left\| \varphi \right\|_{L^{2}(0,T;Z)}^{2} + \left\| \varphi_{\varepsilon} \right\|_{L^{2}(0,T;Z)}^{2} \right) \right\|_{L^{\beta}(\Omega)} \le M_{\beta} \qquad \forall \beta > 1$$

In particular, noting that

$$\frac{1}{q} + \frac{2}{r} < \frac{1}{q} + \frac{q-p}{pq} = \frac{1}{p} < 1$$

choosing $\frac{1}{\beta} := \frac{1}{q} + \frac{2}{r}$ in the estimates above and using the Hölder inequality yields

$$\left\|\frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon}\right\|_{L^{p}(\Omega; L^{2}(0,T;H))} \le M,$$
(5.6)

hence also, by comparison in the equations,

$$\left\|\partial_t \left(\frac{\varphi_{\varepsilon} - \varphi}{\varepsilon}\right)\right\|_{L^p(\Omega; L^2(0,T;Z^*))} + \left\|\frac{\mu_{\varepsilon} - \mu}{\varepsilon}\right\|_{L^p(\Omega; L^2(0,T;H))} + \left\|\partial_t \left(\frac{\sigma_{\varepsilon} - \sigma}{\varepsilon}\right)\right\|_{L^p(\Omega; L^2(0,T;V^*))} \le M.$$
(5.7)

By the uniform estimates (5.5)-(5.7) we deduce that there are

$$x_k \in L^p(\Omega; H^1(0, T; Z^*) \cap L^2(0, T; Z)),$$

$$y_k \in L^p(\Omega; L^2(0, T; H)), \qquad z_k \in L^p(\Omega; H^1(0, T; V^*) \cap L^2(0, T; V))$$

such that, as $\varepsilon \to 0$,

$$\begin{split} & \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \rightharpoonup x_k \qquad \text{in } L^p(\Omega; H^1(0,T;Z^*) \cap L^2(0,T;Z)) \,, \\ & \frac{\mu_{\varepsilon} - \mu}{\varepsilon} \rightharpoonup y_k \qquad \text{in } L^p(\Omega; L^2(0,T;H)) \,, \\ & \frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} \rightharpoonup z_k \qquad \text{in } L^p(\Omega; H^1(0,T;V^*) \cap L^2(0,T;V)) \end{split}$$

and

$$\begin{split} \varphi_{\varepsilon} &\to \varphi & \quad \text{in } L^p(\Omega; C^0([0,T];H) \cap L^2(0,T;Z)) \,, \\ \sigma_{\varepsilon} &\to \sigma & \quad \text{in } L^p(\Omega; C^0([0,T];H) \cap L^2(0,T;V)) \,. \end{split}$$

Since we have the compact inclusions

$$H^{1}(0,T;Z^{*}) \cap L^{2}(0,T;Z) \xrightarrow{c} L^{2}(0,T,V), \qquad H^{1}(0,T;V^{*}) \cap L^{2}(0,T;V) \xrightarrow{c} L^{2}(0,T;H),$$

by Skorokhod theorem there exists a probability space $(\Omega', \mathscr{F}', \mathbb{P}')$ and measurable mappings

$$\phi_{\varepsilon}: (\Omega', \mathscr{F}') \to (\Omega, \mathscr{F})$$

with $\mathbb{P} = \mathbb{P}' \circ \phi_{\varepsilon}^{-1}$ such that

$$\begin{split} \varphi \circ \phi_{\varepsilon} \to \varphi' & \text{in } L^2(0,T;V) \,, \\ \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \circ \phi_{\varepsilon} \to x'_k & \text{in } L^2(0,T;V) \,, \qquad \frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} \circ \phi_{\varepsilon} \to z'_k & \text{in } L^2(0,T;H) \end{split}$$

 $\mathbb{P}'\text{-almost}$ surely, and

$$\begin{split} & \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \circ \phi_{\varepsilon} \rightharpoonup x'_{k} & \text{ in } L^{p}(\Omega'; H^{1}(0, T; Z^{*}) \cap L^{2}(0, T; Z)) \,, \\ & \frac{\mu_{\varepsilon} - \mu}{\varepsilon} \circ \phi_{\varepsilon} \rightharpoonup y'_{k} & \text{ in } L^{p}(\Omega'; L^{2}(0, T; H)) \,, \\ & \frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} \circ \phi_{\varepsilon} \rightharpoonup z'_{k} & \text{ in } L^{p}(\Omega'; H^{1}(0, T; V^{*}) \cap L^{2}(0, T; V)) \end{split}$$

for some φ' , x'_k , y'_k , and z'_k . Moreover, up to extracting a subsequence, the continuity of ψ'' guarantees that

$$\frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon} \circ \phi_{\varepsilon} = \int_{0}^{1} \psi''(\varphi \circ \phi_{\varepsilon} + \tau(\varphi_{\varepsilon} - \varphi) \circ \phi_{\varepsilon}) \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \circ \phi_{\varepsilon} \, \mathrm{d}\tau \to \psi''(\varphi') x'_{k}$$

a.e. in $\Omega \times Q$. Recalling then that the left-hand side is uniformly bounded in $L^p(\Omega'; L^2(0, T; H))$ by (5.6), we deduce also the convergence of the whole sequence

$$\frac{\psi'(\varphi_{\varepsilon}) - \psi'(\varphi)}{\varepsilon} \circ \phi_{\varepsilon} \rightharpoonup \psi''(\varphi') x'_k \quad \text{in } L^p(\Omega'; L^2(0, T; H))$$

Similarly, by the Lipschitz-continuity of h it is immediate to show that

$$\frac{h(\varphi_{\varepsilon}) - h(\varphi)}{\varepsilon} \circ \phi_{\varepsilon} \to h'(\varphi') x'_k \quad \text{in } L^{\ell}(\Omega'; L^2(0, T; H)) \quad \forall \ell \in [1, p)$$

Passing then to the weak limit in the variational formulation of the equations on Ω' we deduce that (x'_k, y'_k, z'_k) solves the linearized system (2.23)–(2.27) on $(\Omega', \mathscr{F}', \mathbb{P}')$ with respect to φ' . Since e have already proved uniqueness for such system, the well-known results by Gyöngy and Krylov [32, Lem 1.1.] ensures that the strong convergences hold in the original probability space $(\Omega, \mathscr{F}, \mathbb{P})$, i.e. that

$$\frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \to x_k \quad \text{in } L^2(0,T;V), \qquad \frac{\sigma_{\varepsilon} - \sigma}{\varepsilon} \to z_k \quad \text{in } L^2(0,T;H)$$

P-almost surely in Ω . Hence, repeating the same argument on $(\Omega, \mathscr{F}, \mathbb{P})$, we have that (x_k, y_k, z_k) is the unique solution to the linearized system in the sense of (2.23)–(2.27). This completes the proof of existence of Theorem 2.8.

Finally, as a consequence of estimates (5.5) and (5.7), we also have that

$$\|x_k\|_{L^p(\Omega; H^1(0,T;Z^*)\cap L^2(0,T;Z))} \le M \|(k_u,k_w)\|_{L^q(\Omega; L^2(0,T;H))^2},$$

so that the map $k \mapsto x_k$ is linear and continuous from $L^q(\Omega; L^2(0,T;H))^2$ to $L^p(\Omega; H^1(0,T;Z^*) \cap L^2(0,T;Z))$, and the Gâteaux-differentiability of Theorem 2.9 is also proved.

5.3 The adjoint system

We prove here existence (and uniqueness) of solutions for the adjoint system. We firstly introduce a suitable approximation of the system so that classical variational theory for Backward SPDEs can be

applied. Then we derive uniform estimates on the solution by using a duality argument and pass to the limit exploiting the linear character of the equations. As regards uniqueness we use again a duality relation.

The approximated problem. For every $n \in \mathbb{N}$, let

$$\psi_n'': \mathbb{R} \to \mathbb{R}, \qquad \psi_n''(r) := \begin{cases} n & \text{if } \psi''(r) > n, \\ \psi''(r) & \text{if } |\psi''(r)| \le n, \quad r \in \mathbb{R}. \\ -n & \text{if } \psi''(r) < -n, \end{cases}$$

We consider the approximated problem

$$-\mathrm{d}\pi_n - A\Delta\tilde{\pi}_n \,\mathrm{d}t + B\psi_n''(\varphi)\tilde{\pi}_n \,\mathrm{d}t = h'(\varphi)(\mathcal{P}\sigma - a - \alpha u)\pi_n \,\mathrm{d}t - ch'(\varphi)\sigma\rho_n \,\mathrm{d}t + \beta_1(\varphi - \varphi_Q) \,dt - \xi_n \,\mathrm{d}W_1 \,,$$
$$\tilde{\pi}_n = -\Delta\pi_n \,,$$
$$-d\rho_n - \Delta\rho_n \,\mathrm{d}t + ch(\varphi)\rho_n \,\mathrm{d}t + b\rho_n \,\mathrm{d}t = \mathcal{P}h(\varphi)\pi_n \,\mathrm{d}t - \theta_n \,\mathrm{d}W_2 \,,$$
$$\pi_n(T) = \beta_2(\varphi(T) - \varphi_T) + \frac{\beta_3}{2} \,, \quad \rho(T) = 0 \,.$$

From the boundedness of $\psi_n''(\varphi)$, h', σ and u and the linear character of the system, we can infer existence and uniqueness of a variational solution due to the classical theory for backward SPDEs (see e.g. [17, § 3]). By rewriting the system as a unique equation in the corresponding product spaces (or by using a fixed point technique) it can be easily shown that that

$$\pi_n \in L^2 \left(\Omega; C^0([0,T];H) \cap L^2(0,T;Z)\right), \tilde{\pi}_n \in L^2 \left(\Omega; C^0([0,T];Z^*) \cap L^2(0,T;H)\right), \rho_n \in L^2 \left(\Omega; C^0([0,T];H) \cap L^2(0,T;V)\right), \xi_n \in L^2(\Omega; L^2(0,T;\mathscr{L}^2(U_1,H))), \qquad \theta_n \in L^2(\Omega; L^2(0,T;\mathscr{L}^2(U_2,H))).$$

Furthermore, by assumption on φ_Q and φ_T , we have $\beta_1(\varphi - \varphi_Q) \in L^6(\Omega; L^2(0, T; H))$ and $\beta_2(\varphi(T) - \varphi_T) \in L^6(\Omega, \mathscr{F}_T; V)$ (recall that $p \ge 6$). Hence, by computing Itô formula for $\|\nabla \pi_n\|_H^2$ and subsequently derive L^6_{Ω} -estimates, it can be shown that the variational solution $(\pi_n, \tilde{\pi}_n, \xi_n, \rho_n, \theta_n)$ given by (2.32)–(2.34) (where ψ'' is replaced by ψ''_n), is actually more regular:

$$\begin{aligned} \pi_n &\in L^6 \left(\Omega; C^0([0,T];V) \cap L^2(0,T;Z \cap H^3(D)) \right) ,\\ \tilde{\pi}_n &\in L^6 \left(\Omega; C^0([0,T];V^*) \cap L^2(0,T;V) \right) ,\\ \rho_n &\in L^6 \left(\Omega; C^0([0,T];H) \cap L^2(0,T;V) \right) ,\\ \xi_n &\in L^6(\Omega; L^2(0,T;\mathscr{L}^2(U_1,V))) , \qquad \theta_n \in L^6(\Omega; L^2(0,T;\mathscr{L}^2(U_2,H))). \end{aligned}$$

We omit here the details not to weight to much the readability of the paper, and we refer to [27, Lem. 4.2] = for what concerns L^p_{Ω} estimates on backward SPDEs and to [52] for the improved regularity in space.

In order to compute uniform estimates on the approximated solutions to the adjoint problem, we need some auxiliary results. First of all, we show that the corresponding approximated linearized system is well-posed in a more general setting, where the forcing terms in the equations are represented by an arbitrary term $\gamma := (\gamma_1, \gamma_2)$:

$$\begin{split} \partial_t x_n^\gamma - \Delta y_n^\gamma &= h(\varphi) \mathcal{P} z_n^\gamma + h'(\varphi) x_n^\gamma (\mathcal{P} \sigma - a - \alpha u) + \gamma_1 & & \text{in } (0,T) \times D \,, \\ y_n^\gamma &= -A \Delta x_n^\gamma + B \psi_n''(\varphi) x_n^\gamma & & \text{in } (0,T) \times D \,, \\ \partial_t z_n^\gamma - \Delta z_n^\gamma + c z_n^\gamma h(\varphi) + c \sigma h'(\varphi) x_n^\gamma + b z_n^\gamma &= \gamma_2 & & \text{in } (0,T) \times D \,, \\ \partial_\mathbf{n} x_n^\gamma &= \partial_\mathbf{n} z_n^\gamma &= 0 & & \text{in } (0,T) \times \partial D \,, \\ x_n^\gamma(0) &= z_n^\gamma(0) &= 0 & & \text{in } D \,. \end{split}$$

Then we introduce the linear map $\tau : \gamma \mapsto (x_n^{\gamma}, x_n^{\gamma}(T))$ assigning to the forcing terms the solution and the solution at final time t = T of the first equation. Notice that a more general map can be studied, also involving z_n^{γ} and allowing for stochastic perturbation in the linearized system (this is not necessary in our situation due to the additive character of the noise in (1.1)). By carefully choosing the functional spaces, the map τ turns out to be continuous along with its adjoint: $\tau^* : (f, \zeta) \to (\pi_n^{\gamma}, \rho_n^{\gamma})$. Observe that the dual operator τ^* maps the forcing term f and the final condition ζ to the corresponding solution $(\pi_n^{\gamma}, \rho_n^{\gamma})$ of the backward equation. Hence, the solution (π_n, ρ_n) we are interested in, can be obtained by evaluating τ^* at $f = \beta_1(\varphi - \varphi_Q)$ and $\zeta = \beta_2(\varphi(T) - \varphi_T) + \frac{\beta_2}{2}$.

Let us start by showing the continuity of the map τ and the subsequent duality formula.

Lemma 5.1. Assume (A1)–(A7), (2.14)–(2.16), (2.19) and (2.21). Let $(u, w) \in \tilde{\mathcal{U}}$ and set $\varphi := \mathcal{S}(u, w)$. Then for every $\gamma := (\gamma_1, \gamma_2) \in L^{6/5}(\Omega; L^1(0, T; H))^2$ and for every $n \in \mathbb{N}$ there exists a unique triple $(x_n^{\gamma}, y_n^{\gamma}, z_n^{\gamma})$ with

$$\begin{split} x_n^{\gamma} \in L^1(\Omega; H^1(0,T;Z^*) \cap L^2(0,T;Z)) \cap L^{6/5}(\Omega;C^0([0,T];V^*) \cap L^2(0,T;V)) \,, \\ y_n^{\gamma} \in L^1(\Omega;L^2(0,T;H)) \,, \\ z_n^{\gamma} \in L^{6/5}(\Omega;H^1(0,T;V^*) \cap L^2(0,T;V)) \,, \end{split}$$

such that

$$\begin{split} \langle \partial_t x_k, \zeta \rangle_V &- \int_D y_k \Delta \zeta = \int_D \left[h(\varphi) \mathcal{P} z_k + h'(\varphi) x_k (\mathcal{P} \sigma - a - \alpha u) + \gamma_1 \right] \zeta \,, \\ &\int_D y_k \zeta = A \int_D \nabla x_k \cdot \nabla \zeta + B \int_D \psi''(\varphi) x_k \zeta \,, \\ \langle \partial_t z_k, \zeta \rangle_V &+ \int_D \nabla z_k \cdot \nabla \zeta + \int_D \left[c z_k h(\varphi) + c \sigma h'(\varphi) x_k + b z_k + \gamma_2 \right] \zeta = 0 \end{split}$$

for every $\zeta \in Z$, for almost every $t \in (0,T)$, \mathbb{P} -almost surely. Moreover, there exists a positive constant M > 0, independent of γ and n, such that

$$\|x_n^{\gamma}\|_{L^{6/5}(\Omega;C^0([0,T];V^*)\cap L^2(0,T;V))} \le M\left(\|\gamma_1\|_{L^{6/5}(\Omega;L^1(0,T;V^*))} + \|\gamma_2\|_{L^{6/5}(\Omega;L^1(0,T;H))}\right), \tag{5.8}$$

$$\|x_n^{\gamma}\|_{L^1(\Omega; C^0([0,T];H)\cap L^2(0,T;Z))} \le M\left(\|\gamma_1\|_{L^{6/5}(\Omega; L^1(0,T;H))} + \|\gamma_2\|_{L^{6/5}(\Omega; L^1(0,T;H))}\right).$$
(5.9)

Finally, it holds that

$$\mathbb{E}\int_{Q}\pi_{n}\gamma_{1} + \mathbb{E}\int_{Q}\rho_{n}\gamma_{2} = \beta_{1}\mathbb{E}\int_{Q}(\varphi - \varphi_{Q})x_{n}^{\gamma} + \mathbb{E}\int_{D}\left(\beta_{2}(\varphi(T) - \varphi_{T}) + \frac{\beta_{3}}{2}\right)x_{n}^{\gamma}(T).$$
(5.10)

Proof. The existence and uniqueness of $(x_n^{\gamma}, y_n^{\gamma}, z_n^{\gamma})$ follows from the fact that $\psi_n''(\varphi) \in L^{\infty}(\Omega \times Q)$ and $\gamma \in L^{6/5}(\Omega; L^1(0, T; H))^2$. Let us show the two estimates. We integrate the first equation on D and test it by $(x_n^k)_D$, then we also test the first equation by $\mathcal{N}(x_n^{\gamma} - (x_n^{\gamma})_D)$, the second by $x_n^{\gamma} - (x_n^{\gamma})_D$, the third by z_n^{γ} . Summing up all the contributions we obtain

$$\frac{1}{2} |(x_n^{\gamma})_D(t)|^2 + \frac{1}{2} ||(x_n^{\gamma} - (x_n^{\gamma})_D)(t)||_{V^*}^2 + \frac{1}{2} ||z_n(t)||_H^2 + A \int_{Q_t} |\nabla x_n^{\gamma}|^2 + \int_{Q_t} |\nabla z_n^{\gamma}|^2 \\
\leq -B \int_{Q_t} \psi_n''(\varphi) x_n^{\gamma} (x_n^{\gamma} - (x_n^{\gamma})_D) + \int_{Q_t} (h(\varphi) \mathcal{P} z_n^{\gamma} + h'(\varphi) x_n^{\gamma} (\mathcal{P} \sigma - a - \alpha u) + \gamma_1)_D (x_n^{\gamma})_D \\
+ \int_{Q_t} [h(\varphi) \mathcal{P} z_n^{\gamma} + h'(\varphi) x_n^{\gamma} (\mathcal{P} \sigma - a - \alpha u) + \gamma_1] \mathcal{N} (x_n^{\gamma} - (x_n^{\gamma})_D) + \int_{Q_t} [\gamma_2 - c\sigma h'(\varphi) x_n^{\gamma}] z_n^{\gamma},$$

which yields then by the assumptions on ψ , the Young and Hölder inequalities, and the boundedness of

 h, h', σ and u,

$$\begin{split} |(x_{n}^{\gamma})_{D}(t)|^{2} + \|(x_{n}^{\gamma} - (x_{n}^{\gamma})_{D})(t)\|_{V^{*}}^{2} + \|z_{n}(t)\|_{H}^{2} + \int_{Q_{t}} |\nabla x_{n}^{\gamma}|^{2} + \int_{Q_{t}} |\nabla z_{n}^{\gamma}|^{2} \\ \lesssim_{A,B,C_{2}} \int_{0}^{t} \left(|(x_{n}^{\gamma})_{D}(s)|^{2} + \|x_{n}^{\gamma}(s)\|_{H}^{2} + \|z_{n}^{\gamma}(s)\|_{H}^{2} \right) ds \\ &+ \int_{0}^{t} \|\gamma_{1}(s)\|_{V^{*}} \left(\|(x_{n}^{\gamma} - (x_{n}^{\gamma})_{D})(s)\|_{V^{*}} + |(x_{n}^{\gamma})_{D}(s)|^{2} \right) ds + \int_{0}^{t} \|\gamma_{2}(s)\|_{H} \|z_{n}^{\gamma}(s)\|_{H} ds \\ \leq \delta \int_{Q_{t}} |\nabla x_{n}^{\gamma}|^{2} + C_{\delta} \int_{0}^{t} \left(\|(x_{n}^{\gamma} - (x_{n}^{\gamma})_{D})(s)\|_{V^{*}}^{2} + |(x_{n}^{\gamma})_{D}(s)|^{2} + \|z_{n}^{\gamma}(s)\|_{H}^{2} \right) ds \\ &+ \int_{0}^{t} (\|\gamma_{1}(s)\|_{V^{*}} \|x_{n}^{\gamma}(s)\|_{V^{*}} + \|\gamma_{2}(s)\|_{H} \|z_{n}^{\gamma}(s)\|_{H}) ds \end{split}$$

for every $\delta > 0$. Choosing the $\delta > 0$ sufficiently small, rearranging the terms and applying the Gronwall lemma in the version [5, Lem. A4–A5] we infer that

$$\|x_n^{\gamma}\|_{L^{\infty}(0,T;V^*)\cap L^2(0,T;V)} + \|z_n^{\gamma}\|_{L^{\infty}(0,T;H)\cap L^2(0,T;V)} \le M\left(\|\gamma_1\|_{L^1(0,T;V^*)} + \|\gamma_2\|_{L^1(0,T;H)}\right) \qquad \mathbb{P}\text{-a.s.}$$

for a certain M > 0 independent of n and γ , from which (5.8) follows. Now, we test the first equation by x_n^{γ} , the second by $-\Delta x_n^{\gamma}$ and take the difference, getting

$$\frac{1}{2} \|x_n^{\gamma}(t)\|_H^2 + A \int_{Q_t} |\Delta x_n^{\gamma}|^2 = B \int_{Q_t} \psi_n''(\varphi) x_n^{\gamma} \Delta x_n^{\gamma} + \int_{Q_t} [h(\varphi) \mathcal{P} z_n^{\gamma} + h'(\varphi) x_n^{\gamma} (\mathcal{P}\sigma - a - \alpha u) + \gamma_1] x_n^{\gamma}.$$

The Young and Hölder inequalities together with the growth assumption on ψ , the continuous embedding $V \hookrightarrow L^6(D)$ and the boundedness of h, σ and u yield then

$$\begin{split} \|x_n^{\gamma}(t)\|_H^2 + \int_{Q_t} |\Delta x_n^{\gamma}|^2 &\lesssim \|z_n^{\gamma}\|_{L^2(0,T;H)}^2 + \int_{Q_t} |\psi''(\varphi)|^2 |x_n^{\gamma}|^2 + \int_{Q_t} |x_n^{\gamma}|^2 + \int_0^t \|\gamma_1(s)\|_H \, \|x_n^{\gamma}(s)\|_H \, \mathrm{d}s \\ &\lesssim \|z_n^{\gamma}\|_{L^2(0,T;H)}^2 + \|\varphi\|_{L^{\infty}(0,T;V)}^4 \, \|x_n^{\gamma}\|_{L^2(0,T;V)}^2 + \int_{Q_t} |x_n^{\gamma}|^2 + \int_0^t \|\gamma_1(s)\|_H \, \|x_n^{\gamma}(s)\|_H \, \mathrm{d}s \, . \end{split}$$

The Gronwall lemma and the fact that $\|z_n^{\gamma}\|_{L^2(0,T;H)} \lesssim \|\gamma\|_{L^1(0,T;H)^2}$ imply again that

$$\|x_n^{\gamma}\|_{L^{\infty}(0,T;H)\cap L^2(0,T;Z)} \le M\left(\|\varphi\|_{L^{\infty}(0,T;V)}^2 \|x_n^{\gamma}\|_{L^2(0,T;V)} + \|\gamma\|_{L^1(0,T;H)^2}\right).$$

Now, note that since $q > p \ge 6$ and $\varphi \in L^r(\Omega; L^{\infty}(0,T;V))$ with $r > \frac{2pq}{q-p} > 12$, we have in particular that $\|\varphi\|^2_{L^{\infty}(0,T;V)} \in L^6(\Omega)$: hence, taking expectations in the last inequality and using the Hölder inequality and (5.8), we deduce also (5.9).

Finally, in order to prove the duality relation (5.10) we test the equation for x_n^{γ} by π_n , the equation for z_n^{γ} by ρ_n , and we subtract the equation for π_n tested by x_n^{γ} and the equation for ρ_n tested by z_n^{γ} . Integrating from 0 to T and taking expectations, using the initial conditions for $(x_n^{\gamma}, z_n^{\gamma})$ and the final conditions for (π_n, ρ_n) , the duality relations follows from usual computations involving integration by parts.

We are now ready to show uniform estimates on the approximated solutions to the adjoint problem. The main tool we have at our disposal is the duality relation (5.10).

First estimate. For every $\gamma \in L^{6/5}(\Omega; L^1(0, T; H))^2$, the duality relation (5.10) implies that

$$\mathbb{E} \int_{Q} \pi_{n} \gamma_{1} + \mathbb{E} \int_{Q} \rho_{n} \gamma_{2} \leq \|\beta_{1}(\varphi - \varphi_{Q})\|_{L^{6}(\Omega; L^{2}(0,T;H))} \|x_{n}^{\gamma}\|_{L^{6/5}(\Omega; L^{2}(0,T;H))} + \left\|\beta_{2}(\varphi(T) - \varphi_{T}) + \frac{\beta_{3}}{2}\right\|_{L^{6}(\Omega; V)} \|x_{n}^{\gamma}(T)\|_{L^{6/5}(\Omega; V^{*})}.$$

Since $\varphi \in L^p(\Omega; L^{\infty}(0,T;V) \cap C^0([0,T];H))$ and $p \ge 6$, by Lemma 5.1 we infer that

$$\mathbb{E}\int_{Q}\pi_{n}\gamma_{1} + \mathbb{E}\int_{Q}\rho_{n}\gamma_{2} \leq M\left(\|\gamma_{1}\|_{L^{6/5}(\Omega;L^{1}(0,T;V^{*}))} + \|\gamma_{2}\|_{L^{6/5}(\Omega;L^{1}(0,T;H))}\right)$$

for a positive constant M independent of γ and n. Since H is dense in V^* , taking supremum over $\gamma \in L^{6/5}(\Omega; L^1(0,T;H))^2$ such that $\|\gamma_1\|_{L^{6/5}(\Omega; L^1(0,T;V^*))} \leq 1$ and $\|\gamma_2\|_{L^{6/5}(\Omega; L^1(0,T;H))} \leq 1$, we deduce that for every $\ell \in [1, +\infty)$

$$\|\pi_n\|_{L^6(\Omega; L^\ell(0,T;V))} + \|\rho_n\|_{L^6(\Omega; L^\ell(0,T;H))} \le M_\ell,$$

so that in particular

$$\|\pi_n\|_{L^6(\Omega \times (0,T);V)} + \|\rho_n\|_{L^6(\Omega \times (0,T);H)} \le M.$$
(5.11)

Second estimate. We write Itô's formula for the sum $\frac{1}{2} \|\pi_n\|_V^2 + \frac{1}{2} \|\rho_n\|_H^2$ (it is crucial not to do it seprately). By noting that $\frac{1}{2}D\|\pi_n\|_V^2 = \pi_n + \Delta \pi_n = \pi_n + \tilde{\pi}_n$, we have

$$\begin{split} &\frac{1}{2} \left\| \pi_n(t) \right\|_V^2 + \frac{1}{2} \left\| \rho_n(t) \right\|_H^2 + A \int_t^T \left(\left\| \nabla \tilde{\pi}_n(s) \right\|_H^2 + \left\| \tilde{\pi}_n(s) \right\|_H^2 \right) \, \mathrm{d}s + \int_t^T \left\| \nabla \rho_n(s) \right\|_H^2 \, \mathrm{d}s \\ &+ \int_t^T \!\!\!\!\!\int_D (ch(\varphi(s)) + b) |\rho_n(s)|^2 + \frac{1}{2} \int_t^T \left\| \xi_n(s) \right\|_{\mathscr{L}^2(U_1,V)}^2 \, \mathrm{d}s + \frac{1}{2} \int_t^T \left\| \theta_n(s) \right\|_{\mathscr{L}^2(U_2,H)}^2 \, \mathrm{d}s \\ &= \frac{1}{2} \left\| \beta_2(\varphi(T) - \varphi_T) \right\|_V^2 + \int_t^T \!\!\!\!\int_D \beta_1(\varphi - \varphi_Q)(s) \left(\tilde{\pi}_n(s) + \pi_n(s) \right) \, \mathrm{d}s + \int_t^T \!\!\!\!\int_D \mathcal{P}h(\varphi(s))\pi_n(s)\rho_n(s) \, \mathrm{d}s \\ &+ \int_t^T \!\!\!\!\int_D \left[h'(\varphi(s))(\mathcal{P}\sigma(s) - a - \alpha u(s))\pi_n(s) - ch'(\varphi(s))\sigma(s)\rho_n(s) \right] \left(\tilde{\pi}_n(s) + \pi_n(s) \right) \, \mathrm{d}s \\ &- B \int_t^T \!\!\!\!\int_D \psi_n''(\varphi(s)) |\tilde{\pi}_n(s)|^2 \, \mathrm{d}s - B \int_t^T \!\!\!\!\int_D \psi_n''(\varphi(s))\tilde{\pi}_n(s)\pi_n(s)\pi_n(s) \, \mathrm{d}s \\ &- \int_t^T \left(\xi_n(s), \tilde{\pi}_n(s) + \pi_n(s) \right)_H \, \mathrm{d}W_1(s) - \int_t^T \left(\theta_n(s), \rho_n(s) \right)_H \, \mathrm{d}W_2(s) \, . \end{split}$$

Taking expectations, using the boundedness of h, σ and u together with the Young inequality, the first four terms on the right-hand side are estimated by

$$\begin{aligned} \|\beta_{2}(\varphi(T) - \varphi_{T})\|_{L^{2}(\Omega; V)}^{2} + C_{\varepsilon} \|\beta_{1}(\varphi - \varphi_{Q})\|_{L^{2}(\Omega; L^{2}(0, T; H))}^{2} + \varepsilon \mathbb{E} \int_{t}^{T} \|\tilde{\pi}_{n}(s)\|_{H}^{2} \, \mathrm{d}s \\ + C_{\varepsilon} \|\pi_{n}\|_{L^{2}(\Omega; L^{2}(0, T; H))}^{2} + C_{\varepsilon} \|\rho_{n}\|_{L^{2}(\Omega; L^{2}(0, T; H))}^{2} \end{aligned}$$

for every $\varepsilon > 0$. Secondly, note that by definition of $\tilde{\pi}_n$ and the fact that $\psi'' \ge -C_2$,

$$-B \mathbb{E} \int_{t}^{T} \int_{D} \psi_{n}''(\varphi(s)) |\tilde{\pi}_{n}(s)|^{2} ds \leq BC_{2} \int_{t}^{T} \int_{D} |\tilde{\pi}_{n}|^{2} \leq \varepsilon \mathbb{E} \int_{t}^{T} \|\nabla \tilde{\pi}_{n}(s)\|_{H}^{2} ds + C_{\varepsilon} \mathbb{E} \int_{t}^{T} \|\tilde{\pi}_{n}(s)\|_{V^{*}}^{2} ds$$
$$\leq \varepsilon \mathbb{E} \int_{t}^{T} \|\nabla \tilde{\pi}_{n}(s)\|_{H}^{2} ds + C_{\varepsilon} \mathbb{E} \int_{t}^{T} \|\nabla \pi_{n}(s)\|_{H}^{2} ds.$$

Finally, by the Hölder inequality, the growth assumption on ψ'' and the continuous inclusion $V \hookrightarrow L^6(D)$ we have

$$\begin{split} -B &\mathbb{E} \int_{t}^{T} \int_{D} \psi_{n}''(\varphi(s)) \tilde{\pi}_{n}(s) \pi_{n}(s) \,\mathrm{d}s \lesssim \mathbb{E} \int_{t}^{T} \int_{D} \left(1+|\varphi(s)|^{2} \right) |\tilde{\pi}_{n}(s)| |\pi_{n}(s)| \\ &\leq \varepsilon \,\mathbb{E} \int_{t}^{T} \left\| \tilde{\pi}_{n}(s) \right\|_{H}^{2} \,\mathrm{d}s + C_{\varepsilon} \,\mathbb{E} \int_{t}^{T} \int_{D} \left(1+|\varphi(s)|^{4} \right) |\pi_{n}(s)|^{2} \\ &\leq \varepsilon \,\mathbb{E} \int_{t}^{T} \left\| \tilde{\pi}_{n}(s) \right\|_{H}^{2} \,\mathrm{d}s + C_{\varepsilon} \left\| \varphi \right\|_{L^{6}(\Omega \times (0,T);V)}^{6} + C_{\varepsilon} \left\| \pi_{n} \right\|_{L^{6}(\Omega \times (0,T);V)}^{6} \,. \end{split}$$

Taking into account the positivity of $ch(\varphi) + b$, rearranging the terms and choosing ε sufficiently small, the estimate (5.11) yields

$$\|\pi_n\|_{L^{\infty}(0,T;L^2(\Omega;V))} + \|\tilde{\pi}_n\|_{L^2(\Omega;L^2(0,T;V))} + \|\rho_n\|_{L^{\infty}(0,T;L^2(\Omega;H))\cap L^2(\Omega;L^2(0,T;V))} \le M,$$
(5.12)

$$\|\xi_n\|_{L^2(\Omega;L^2(0,T;\mathscr{L}^2(U_1,V)))} + \|\theta_n\|_{L^2(\Omega;L^2(0,T;\mathscr{L}^2(U_2,H)))} \le M$$
(5.13)

for a positive constant M independent of n.

Going back now to Itô's formula, taking first supremum in time and then expectations, using (5.13) with Burkholder-Davis-Gundy inequality, a classical procedure yields also by elliptic regularity

$$\|\pi_n\|_{L^2(\Omega;L^{\infty}(0,T;V)\cap L^2(0,T;Z\cap H^3(D)))} + \|\tilde{\pi}_n\|_{L^2(\Omega;L^{\infty}(0,T;V^*)\cap L^2(0,T;V))} \le M,$$
(5.14)

$$\|\rho_n\|_{L^2(\Omega; L^{\infty}(0,T;H)\cap L^2(0,T;V))} \le M.$$
(5.15)

Passage to the limit. From (5.14)–(5.15) we obtain that there exist

$$\begin{split} \pi \in L^{\infty}(0,T;L^{2}(\Omega;V)) \cap L^{2}(\Omega;L^{2}(0,T;Z \cap H^{3}(D))) \,, \\ \tilde{\pi} \in L^{\infty}(0,T;L^{2}(\Omega;V^{*})) \cap L^{2}(\Omega;L^{2}(0,T;V)) \,, \\ \rho \in L^{\infty}(0,T;L^{2}(\Omega;H)) \cap L^{2}(\Omega;L^{2}(0,T;V)) \,, \\ \xi \in L^{2}(\Omega;L^{2}(0,T;\mathscr{L}^{2}(U_{1},V))) \,, \qquad \theta \in L^{2}(\Omega;L^{2}(0,T;\mathscr{L}^{2}(U_{2},H))) \,, \end{split}$$

such that

$$\begin{aligned} \pi_n &\stackrel{*}{\rightharpoonup} \pi & \text{ in } L^{\infty}(0,T;L^2(\Omega;V)) \cap L^2(\Omega;L^2(0,T;Z \cap H^3(D))) \\ \tilde{\pi}_n &\stackrel{*}{\rightharpoonup} \tilde{\pi} & \text{ in } L^{\infty}(0,T;L^2(\Omega;V^*)) \cap L^2(\Omega;L^2(0,T;V)) \,, \\ \rho_n &\stackrel{*}{\rightharpoonup} \rho & \text{ in } L^{\infty}(0,T;L^2(\Omega;H)) \cap L^2(\Omega;L^2(0,T;V)) \,, \\ \xi_n &\rightharpoonup \xi & \text{ in } L^2(\Omega;L^2(0,T;\mathscr{L}^2(U_1,V))) \,, \\ \theta_n &\rightharpoonup \theta & \text{ in } L^2(\Omega;L^2(0,T;\mathscr{L}^2(U_2,H))) \,. \end{aligned}$$

Notice that, in order to extract the above (weakly converging) sequences, we modified on purpose the functional setting so to guarantee reflexivity of the functional spaces in consideration. The weak convergences are enough to pass to the limit in each term of the variational formulation of the approximated problem. Let us show in detail only the term involving ψ''_n . First of all, since $\psi''(\varphi) \in L^{p/2}(\Omega; L^{\infty}(0, T; L^3(D)))$ and $p \geq 6$, by the properties of the truncation operator we have

$$\psi_n''(\varphi) \to \psi''(\varphi) \quad \text{in } L^3(\Omega \times (0,T) \times D).$$

Hence, by strong-weak convergence we infer that

$$\psi_n''(\varphi)\tilde{\pi}_n \rightharpoonup \psi''(\varphi)\tilde{\pi} \qquad \text{in } L^{6/5}(\Omega \times (0,T) \times D)$$

By linearity of the approximated adjoint system, letting $n \to \infty$ we obtain exactly conditions (2.32)–(2.34). The further regularity in (2.28)–(2.31) is recovered a posteriori in the limit equation by Itô's formula. This completes the proof of existence for the adjoint system.

Uniqueness. For every $\gamma \in L^{6/5}(\Omega; L^1(0, T; H))^2$, using the estimates (5.8)–(5.9) and arguing as in the proof of Lemma 5.1, it is straightforward to prove the existence of $(x^{\gamma}, y^{\gamma}, z^{\gamma})$ such that

$$\begin{split} x^{\gamma} &\in L^{1}(\Omega; H^{1}(0,T;Z^{*}) \cap L^{2}(0,T;Z)) \cap L^{6/5}(\Omega;C^{0}([0,T];V^{*}) \cap L^{2}(0,T;V)) \,, \\ & y^{\gamma} \in L^{1}(\Omega; L^{2}(0,T;H)) \,, \\ & z^{\gamma} \in L^{6/5}(\Omega; H^{1}(0,T;V^{*}) \cap L^{2}(0,T;V)) \,, \end{split}$$

such that

$$\begin{split} \partial_t x^{\gamma} - \Delta y^{\gamma} &= h(\varphi) \mathcal{P} z^{\gamma} + h'(\varphi) x^{\gamma} (\mathcal{P} \sigma - a - \alpha u) + \gamma_1 & & \text{in } (0, T) \times D \,, \\ y^{\gamma} &= -A \Delta x^{\gamma} + B \psi''(\varphi) x^{\gamma} & & \text{in } (0, T) \times D \,, \\ \partial_t z^{\gamma} - \Delta z^{\gamma} + c z^{\gamma} h(\varphi) + c \sigma h'(\varphi) x^{\gamma} + b z^{\gamma} &= \gamma_2 & & \text{in } (0, T) \times D \,, \\ \partial_{\mathbf{n}} x^{\gamma} &= \partial_{\mathbf{n}} z^{\gamma} = 0 & & & \text{in } (0, T) \times \partial D \,, \\ x^{\gamma}(0) &= z^{\gamma}(0) = 0 & & & \text{in } D \,. \end{split}$$

C. Orrieri, E. Rocca, L. Scarpa

Assume now that $(\pi_i, \tilde{\pi}_i, \rho_i, \xi_i, \theta_i)$ solve (2.28)–(2.34) for i = 1, 2. Testing the equation for x^{γ} by π_i , the equation for z^{γ} by ρ_i , and subtracting the equation for π_i tested by x^{γ} and the equation for ρ_i tested by z^{γ} , integrating from 0 to T and taking expectations, using the initial conditions for (x^{γ}, z^{γ}) and the final conditions for (π_i, ρ_i) , we get the duality relation

$$\mathbb{E}\int_{Q}\pi_{i}\gamma_{1} + \mathbb{E}\int_{Q}\rho_{i}\gamma_{2} = \beta_{1}\mathbb{E}\int_{Q}(\varphi - \varphi_{Q})x^{\gamma} + \mathbb{E}\int_{D}\left(\beta_{2}(\varphi(T) - \varphi_{T}) + \frac{\beta_{3}}{2}\right)x^{\gamma}(T)$$

for every $\gamma \in L^2(\Omega; L^2(0,T;H))^2$ and for i = 1, 2. Hence, subtracting we infer that

$$\mathbb{E}\int_{Q}(\pi_{1}-\pi_{2})\gamma_{1}+\mathbb{E}\int_{Q}(\rho_{1}-\rho_{2})\gamma_{2}=0$$

from which $\pi_1 = \pi_2$ and $\rho_1 = \rho_2$. The uniqueness of the other solution components follows then by comparison in the equations.

5.4 First order conditions for optimality

We prove here the first version of the necessary conditions for optimality contained in Proposition 2.10 and Theorem 2.12.

Let $(\bar{u}, \bar{w}) \in \mathcal{U}$ be an optimal control, $\varphi := \mathcal{S}(\bar{u}, \bar{w})$ the corresponding state and let $(u, w) \in \mathcal{U}$ be arbitrary. Setting $k := (k_u, k_w) := (u - \bar{u}, w - \bar{w})$, by the convexity of \mathcal{U} we have that $(\bar{u} + \varepsilon k_u, \bar{w} + \varepsilon k_w) \in \mathcal{U}$ for every $\varepsilon \in [0, 1]$: hence, by definition of optimal control we have, setting $\varphi_{\varepsilon} := \mathcal{S}(\bar{u} + \varepsilon k_u, \bar{w} + \varepsilon k_w)$,

$$\begin{split} J(\bar{\varphi}, \bar{u}, \bar{w}) &\leq \frac{\beta_1}{2} \mathbb{E} \int_Q |\varphi_{\varepsilon} - \varphi_Q|^2 + \frac{\beta_2}{2} \mathbb{E} \int_D |\varphi_{\varepsilon}(T) - \varphi_T|^2 + \frac{\beta_3}{2} \mathbb{E} \int_D (\varphi_{\varepsilon}(T) + 1) \\ &+ \frac{\beta_4}{2} \mathbb{E} \int_Q |\bar{u} + \varepsilon k_u|^2 + \frac{\beta_5}{2} \mathbb{E} \int_Q |\bar{w} + \varepsilon k_w|^2 \,. \end{split}$$

Solving the square-powers on the right-hand side and plugging in the definition of J yields

$$0 \leq \frac{\beta_1}{2} \mathbb{E} \int_Q \left(|\varphi_{\varepsilon}|^2 - |\bar{\varphi}|^2 - 2(\varphi_{\varepsilon} - \bar{\varphi})\varphi_Q \right) + \frac{\beta_2}{2} \mathbb{E} \int_D \left(|\varphi_{\varepsilon}(T)|^2 - |\bar{\varphi}(T)|^2 - 2(\varphi_{\varepsilon} - \bar{\varphi})(T)\varphi_T \right) \\ + \frac{\beta_3}{2} \mathbb{E} \int_D (\varphi_{\varepsilon} - \bar{\varphi})(T) + \frac{\beta_4}{2} \mathbb{E} \int_Q \left(\varepsilon^2 |k_u|^2 + 2\varepsilon \bar{u}k_u \right) + \frac{\beta_5}{2} \mathbb{E} \int_Q \left(\varepsilon^2 |k_w|^2 + 2\varepsilon \bar{w}k_w \right) \,.$$

Dividing by ε and using the Gâteaux-differentiability of J (all the terms admit Gâteaux differential) we infer that

$$0 \leq \beta_1 \mathbb{E} \int_Q \left(\int_0^1 (\bar{\varphi} + \tau(\varphi_{\varepsilon} - \bar{\varphi})) \, d\tau - \varphi_Q \right) \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} \\ + \beta_2 \mathbb{E} \int_D \left(\int_0^1 (\bar{\varphi} + \tau(\varphi_{\varepsilon} - \bar{\varphi}))(T) \, d\tau - \varphi_T \right) \frac{\varphi_{\varepsilon} - \varphi}{\varepsilon} (T) \\ + \frac{\beta_3}{2} \mathbb{E} \int_D \frac{\varphi_{\varepsilon} - \bar{\varphi}}{\varepsilon} (T) + \beta_4 \mathbb{E} \int_Q \bar{u} k_u + \beta_5 \mathbb{E} \int_Q \bar{w} k_w + \varepsilon \left(\frac{\beta_4}{2} \mathbb{E} \int_Q |k_u|^2 + \frac{\beta_5}{2} \mathbb{E} \int_Q |k_w|^2 \right) \, .$$

By Theorem 2.9 we have that $\varphi_{\varepsilon} - \bar{\varphi} \to 0$ strongly in $L^2(\Omega; C^0([0, T]; H))$, and $\frac{\varphi_{\varepsilon} - \bar{\varphi}}{\varepsilon} \to x_k$ strongly in $L^2(\Omega; L^2(0, T; H))$ and $\frac{\varphi_{\varepsilon} - \bar{\varphi}}{\varepsilon}(T) \to x_k(T)$ weakly in $L^2(\Omega; H)$, so that letting $\varepsilon \searrow 0$ we can conclude that

$$\beta_1 \mathbb{E} \int_Q (\bar{\varphi} - \varphi_Q) x_k + \beta_2 \mathbb{E} \int_D (\bar{\varphi}(T) - \varphi_T) x_k(T) + \frac{\beta_3}{2} \mathbb{E} \int_D x_k(T) + \beta_4 \mathbb{E} \int_Q \bar{u} k_u + \beta_5 \mathbb{E} \int_Q \bar{w} k_w \ge 0$$

and Proposition 2.10 is proved.

Finally, note that choosing $\gamma_1 = -\alpha h(\varphi)k_u$ and $\gamma_2 = bk_w$ we get $x^{\gamma} = x_k$, $y^{\gamma} = y_k$ and $z^{\gamma} = z_k$ by Theorem 2.8, so that as we have already pointed out in the previous sections, the following duality formula holds:

$$-\alpha \mathbb{E} \int_{Q} \pi h(\bar{\varphi}) k_{u} + b \mathbb{E} \int_{Q} \rho k_{w} = \beta_{1} \mathbb{E} \int_{Q} (\bar{\varphi} - \varphi_{Q}) x_{k} + \mathbb{E} \int_{D} \left(\beta_{2}(\bar{\varphi}(T) - \varphi_{T}) + \frac{\beta_{3}}{2} \right) x_{k}(T) \,.$$

By comparison we obtain the desired inequality, and Theorem 2.12 is finally proved.

References

- A. Agosti, P.F. Antonietti, P. Ciarletta, M. Grasselli, and M. Verani. A Cahn-Hilliard-type equation with application to tumor growth dynamics. *Math. Methods Appl. Sci.*, 40:7598–7626, 2017.
- [2] V. Barbu, M. Röckner, and D. Zhang. Optimal bilinear control of nonlinear stochastic Schrödinger equations driven by linear multiplicative noise. Ann. Probab., 46:1957–1999, 2018.
- [3] C. Bauzet, E. Bonetti, G. Bonfanti, F. Lebon, and G. Vallet. A global existence and uniqueness result for a stochastic Allen-Cahn equation with constraint. *Math. Methods Appl. Sci.*, 40(14):5241–5261, 2017.
- [4] D. Breit, E. Feireisl, and M. Hofmanová. Stochastically forced compressible fluid flows, volume 3 of De Gruyter Series in Applied and Numerical Mathematics. De Gruyter, Berlin, 2018.
- [5] H. Brézis. Opérateurs maximaux monotones et semi-groupes de contractions dans les espaces de Hilbert. North-Holland Publishing Co., Amsterdam-London; American Elsevier Publishing Co., Inc., New York, 1973. North-Holland Mathematics Studies, No. 5. Notas de Matemática (50).
- [6] J.W. Cahn and J.E. Hilliard. Free energy of a nonuniform system. I. Interfacial free energy. J. Chem. Phys., 28: 258–267, 1958.
- [7] V. Capasso and D. Morale. Stochastic modelling of tumour-induced angiogenesis. J. Math. Biol., 58:219–233, 2009.
- [8] C. Cavaterra, E. Rocca, and H. Wu. Long-time Dynamics and Optimal Control of a Diffuse Interface Model for Tumor Growth. Appl. Math. Optim., to appear, 2019. DOI: 10.1007/s00245-019-09562-5.
- P. Colli, G. Gilardi, G. Marinoschi, and E. Rocca. Sliding mode control for phase field system related to tumor growth. Appl. Math. Optim., 79(3): 647–670, 2019.
- [10] P. Colli, G. Gilardi, E. Rocca and J. Sprekels. Optimal distributed control of a diffuse interface model of tumor growth. Nonlinearity, 30:2518–2546, 2017.
- [11] P. Colli, H. Gomez, G. Lorenzo, G. Marinoschi, A. Reali, E. Rocca, Mathematical analysis and simulation study of a phase-field model of prostate cancer growth with chemotherapy and antiangiogenic therapy effects, preprint arXiv:1907.11618 (2019).
- [12] F. Cornalba. A nonlocal stochastic Cahn-Hilliard equation. Nonlinear Anal., 140:38–60, 2016.
- [13] V. Cristini, X. Li, J.S. Lowengrub, and S.M. Wise. Nonlinear simulations of solid tumor growth using a mixture model: invasion and branching. J. Math. Biol., 58:723-763, 2009.
- [14] V. Cristini and J. Lowengrub. Multiscale modeling of cancer. An integrated experimental and mathematical modeling approach. Cambridge Univ. Press, 2010.
- [15] G. Da Prato and A. Debussche. Stochastic Cahn-Hilliard equation. Nonlinear Anal., 26(2):241-263, 1996.
- [16] A. Debussche and L. Goudenège. Stochastic Cahn-Hilliard equation with double singular nonlinearities and two reflections. SIAM J. Math. Anal., 43(3):1473–1494, 2011.
- [17] K. Du and Q. Meng. A revisit to W_2^n -theory of super-parabolic backward stochastic partial differential equations in \mathbb{R}^d . Stochastic Process. Appl., 120(10):1996–2015, 2010.
- [18] M. Ebenbeck and P. Knopf. Optimal control theory and advanced optimality conditions for a diffuse interface model of tumor growth. arXiv:1903.00333, 2019.
- [19] M. Ebenbeck and P. Knopf. Optimal medication for tumors modeled by a Cahn-Hilliard-Brinkman equation. Calc. Var. Partial Differential Equations, 58:131, 2019. https://doi.org/10.1007/s00526-019-1579-z
- [20] N. Elezović and A. Mikelić. On the stochastic Cahn-Hilliard equation. Nonlinear Anal., 16(12):1169–1200, 1991.
- [21] E. Feireisl and M. Petcu. A diffuse interface model of a two-phase flow with thermal fluctuations. ArXiv e-prints, Apr. 2018.
- [22] E. Feireisl and M. Petcu. Stability of strong solutions for a model of incompressible two-phase flow under thermal fluctuations. J. Differential Equations, 267(3):1836–1858, 2019.
- [23] F. Flandoli and D. Gatarek. Martingale and stationary solutions for stochastic Navier-Stokes equations. Probab. Theory Related Fields, 102(3):367–391, 1995.
- [24] S. Frigeri, K.F. Lam, E. Rocca, and G. Schimperna. On a multi-species Cahn-Hilliard-Darcy tumor growth model with singular potentials. Comm Math Sci., 16:821–856, 2018.
- [25] M. Fuhrman, Y. Hu, and G. Tessitore. Stochastic maximum principle for optimal control of SPDEs. C. R. Math. Acad. Sci. Paris, 350(13–14):683–688, 2012.
- [26] M. Fuhrman and C. Orrieri. Stochastic maximum principle for optimal control of a class of nonlinear SPDEs with dissipative drift. SIAM J. Control Optim., 54(1):341–371, 2016.
- [27] M. Fuhrman and G. Tessitore. Nonlinear Kolmogorov equations in infinite dimensional spaces: the backward stochastic differential equations approach and applications to optimal control. Ann. Probab., 30(3):1397–1465, 2002.
- [28] H. Garcke, K.F. Lam, R. Nürnberg, and E. Sitka. A multiphase Cahn-Hilliard-Darcy model for tumour growth with necrosis. *Math. Models Methods Appl. Sci.*, 28:525–577, 2018.
- [29] H. Garcke, K.F. Lam, E. Sitka, and V. Styles. A Cahn-Hilliard-Darcy model for tumour growth with chemotaxis and active transport. *Math. Models Methods Appl. Sci.*, 26: 1095–1148, 2016.
- [30] H. Garcke, K. F. Lam, and E. Rocca. Optimal control of treatment time in a diffuse interface model of tumor growth. *Appl. Math. Optim.*, 78: 495–544, 2018.

- [31] L. Goudenège. Stochastic Cahn-Hilliard equation with singular nonlinearity and reflection. Stochastic Process. Appl., 119(10):3516–3548, 2009.
- [32] I. Gyöngy and N. Krylov. Existence of strong solutions for Itô's stochastic equations via approximations. Probab. Theory Related Fields, 105(2):143–158, 1996.
- [33] H. Gudbjartsson and S. Patz. The Rician Distribution of Noisy MRI Data. Magn Reson Med, 34(6): 910–914, 1995.
- [34] A. Hawkins-Daarud, K.G. van der Zee, and J.T. Oden. Numerical simulation of a thermodynamically consistent fourspecies tumor growth model. Int. J. Numer. Meth. Biomed. Engng., 28:3–24, 2011.
- [35] D. Hilhorst, J. Kampmann, T.N. Nguyen, and K.G. van der Zee. Formal asymptotic limit of a diffuse-interface tumorgrowth model. Math. Models Methods Appl. Sci., 25:1011–1043, 2015.
- [36] N. Ikeda and S. Watanabe. Stochastic differential equations and diffusion processes, volume 24 of North-Holland Mathematical Library. North-Holland Publishing Co., Amsterdam; Kodansha, Ltd., Tokyo, second edition, 1989.
- [37] E.A.B.F. Lima, R.C. Almeida, and J.T. Oden. Analysis and numerical solution of stochastic phase-field models of tumor growth. Numerical Methods for Partial Differential Equations, 31(2), 2014.
- [38] W. Liu and M. Röckner. Stochastic partial differential equations: an introduction. Springer, Cham, 2015.
- [39] C. Marinelli and L. Scarpa. A variational approach to dissipative SPDEs with singular drift. Ann. Probab., 46(3):1455– 1497, 2018.
- [40] C. Marinelli and L. Scarpa. Refined existence and regularity results for a class of semilinear dissipative SPDEs. ArXiv e-prints, Nov. 2017.
- [41] M. Métivier. Semimartingales, volume 2 of de Gruyter Studies in Mathematics. Walter de Gruyter & Co., Berlin-New York, 1982. A course on stochastic processes.
- [42] A. Miranville, E. Rocca, and G. Schimperna. On the long time behavior of a tumor growth model. Journal of Differential Equations, 67: 2616–2642, 2019.
- [43] S. Mirrahimi and B. Perthame. Asymptotic analysis of a selection model with space. J. Math. Pures Appl., 104: 1108–1118, 2018.
- [44] Y. Moghadamnia and D. Moslemi. A stochastic mathematical model of avascular tumor growth pattern and its treatment by means of noises. *Caspian J. Intern. Med.*, 8(4): 258–263, 2017.
- [45] A. Niemisto, V. Dunmire, O. Yli-Harja, W. Zhang, and I. Shmulevich. Analysis of angiogenesis using in vitro experiments and stochastic growth models. *Physical Review E*, 72, 2005.
- [46] G. Guatteri, F. Masiero, and C. Orrieri. Stochastic maximum principle for SPDEs with delay. Stochastic Process. Appl., 127(7):2396–2427, 2017.
- [47] C. Orrieri and L. Scarpa. Singular stochastic Allen-Cahn equations with dynamic boundary conditions. J. Differential Equations, 266(8):4624–4667, 2019.
- [48] E. Pardoux. Equations aux derivées partielles stochastiques nonlinéaires monotones. PhD thesis, Université Paris XI, 1975.
- [49] C. Pouchol, J. Clairambault, A. Lorz, and E. Trélat. Asymptotic analysis and optimal control of an integro-differential system modelling healthy and cancer cells exposed to chemotherapy. J. Math. Pures Appl., 116: 268–308, 2018.
- [50] L. Scarpa. On the stochastic Cahn-Hilliard equation with a singular double-well potential. Nonlinear Anal., 171:102– 133, 2018.
- [51] L. Scarpa. The stochastic viscous Cahn-Hilliard equation: well-posedness, regularity and vanishing viscosity limit. ArXiv e-prints, Sept. 2018.
- [52] L. Scarpa. Optimal distributed control of a stochastic Cahn-Hilliard equation. arXiv:1810.09292v1, 2018.
- [53] A. Signori. Vanishing parameter for an optimal control problem modeling tumor growth. arXiv:1903.04930, 2019.
- [54] J. Simon. Compact sets in the space $L^p(0,T;B)$. Ann. Mat. Pura Appl. (4), 146:65–96, 1987.
- [55] J. Sprekels and H. Wu. Optimal distributed control of a Cahn-Hilliard-Darcy system with mass sources. Appl. Math. Optim., to appear, 2019. DOI: 10.1007/s00245-019-09555-4.
- [56] W.Y. Tan, and C.W. Chen. Stochastic modeling of carcinogenesis: Some new insights. Mathematical and Computer Modelling, 28:49–71, 1998.
- [57] A. W. van der Vaart and J. A. Wellner. Weak convergence and empirical processes. Springer Series in Statistics. Springer-Verlag, New York, 1996. With applications to statistics.
- [58] G. Vallet and A. Zimmermann. Well-posedness for a pseudomonotone evolution problem with multiplicative noise. J. Evol. Equ., 19:153–202, 2019.
- [59] S.M. Wise, J.S. Lowengrub, H.B. Frieboes, and V. Cristini. Three-dimensional multispecies nonlinear tumor growth–I: model and numerical method. J. Theoret. Biol., 253:524–543, 2008.
- [60] J. Yong and X. Y. Zhou. Stochastic controls, volume 43 of Applications of Mathematics (New York). Springer-Verlag, New York, 1999.