

1 **ERGODIC BSDES WITH MULTIPLICATIVE AND DEGENERATE NOISE**

2 GIUSEPPINA GUATTERI* GIANMARIO TESSITORE†

3 **Abstract.** In this paper we study an Ergodic Markovian BSDE involving a forward process X that solves
 4 an infinite dimensional forward stochastic evolution equation with multiplicative and possibly degenerate diffusion
 5 coefficient. A concavity assumption on the driver allows us to avoid the typical quantitative conditions relating the
 6 dissipativity of the forward equation and the Lipschitz constant of the driver. Although the degeneracy of the noise
 7 has to be of a suitable type we can give a stochastic representation of a large class of Ergodic HJB equations; moreover
 8 our general results can be applied to get the synthesis of the optimal feedback law in relevant examples of ergodic
 9 control problems for SPDEs.

10 **Keywords:** Ergodic control; infinite dimensional SDEs; BSDEs; Multiplicative Noise

11 **2010 Mathematics Subject Classification:** 60H15, 60H30, 37A50.

12 **1. Introduction.** In this paper we study the following BSDE of ergodic type

$$Y_t^x = Y_T^x + \int_t^T [\widehat{\psi}(X_s^x, Z_s^x, U_s^x) - \lambda] ds - \int_t^T Z_s^x dW_s^1 - \int_t^T U_s^x dW_s^2, \quad 0 \leq t \leq T < \infty,$$

12 where the processes (Y^x, Z^x, U^x) and the constant λ are the unknowns of the above equation while
 13 the diffusion X is the (mild) solution of the infinite dimensional (forward) SDE:

14 (1.1)
$$\begin{cases} dX_s^x = AX_s^x ds + F(X_s^x) ds + QG(X_s^x) dW_s^1 + DdW_s^2, \\ X_t^x = x. \end{cases}$$

15 In the above equation X takes values in an Hilbert space H and W^1, W^2 are independent cylindrical
 16 Wiener processes (see (A.1)-(A.6) in Section 3 and (B.1) in Section 4 for precise description of
 17 the other terms). We just stress that we will assume that $G(x)$ is invertible for all $x \in H$ while Q
 18 and D will be general, possibly degenerate, linear operators.

19 Ergodic BSDEs have been introduced in [19] in relation to optimal stochastic ergodic control
 20 problems and as a tool to study the asymptotic behaviour of parabolic HJB equations and conse-
 21 quently to give a stochastic representation to the limit semilinear elliptic PDEs (see equation (5.1)
 22 below).

23 In [19] the same class of BSDEs have been introduced, already in an infinite dimensional frame-
 24 work, but only in the case in which the noise coefficient was constant ($Q = 0$ in our notation).
 25 Successive works, see [15] and [7] weakened the assumptions and refined the results in the same
 26 *additive noise* case. Then in [24], in a finite dimensional framework, the case of ‘multiplicative noise
 27 ($Q \neq 0$ and G depending on x in our notation) is treated under quantitative conditions relating
 28 the dissipativity constant of the forward equation to the Lipschitz norm of $\widehat{\psi}$ with respect to Z .
 29 Afterwards, in [21], still in finite dimensions, such quantitative assumptions are dropped in the case
 30 of a non degenerate and bounded diffusion coefficient ($Q = I$ and G bounded and invertible in our
 31 notation) by a careful use of smoothing properties of the Kolmogorov semigroup associated to the
 32 non-degenerate underlying diffusion X . Finally in [14] the result is extended to the case of non de-
 33 generate but unbounded (linearly growing) diffusion coefficients ($Q = I$ and G invertible and linearly

*DEPARTMENT OF MATHEMATICS, POLITECNICO DI MILANO, PIAZZA LEONARDO DA VINCI
 32, 20133 MILANO, ITALY, GIUSEPPINA.GUATTERI@POLIMI.IT.

†DEPARTMENT OF MATHEMATICS AND APPLICATIONS, UNIVERSITY OF MILANO-BICOCCA,
 VIA R. COZZI 53 - BUILDING U5, 20125 MILANO, ITALY, GIANMARIO.TESSITORE@UNIMIB.IT.

growing in our notation). To complete the picture we mention, [2], [3], [4] and [13] where Ergodic BSDEs are studied in various frameworks different from the present one: namely, respectively when they are driven by a Markov chain, in the context (see [17]) of randomized control problems and BSDEs with constraints on the martingale term both in finite and in infinite dimensions and finally in the context of G - expectations theory.

In this paper we propose an alternative approach that works well in the infinite dimensional case and allows to consider degenerate multiplicative noise (Q in general non invertible and G bounded invertible but depending on x). On the other side we have to assume that $\widehat{\psi}$ has the form:

$$\widehat{\psi}(x, z, u) := \psi(x, zG^{-1}(x), u),$$

where ψ is Lipschitz and *concave* function with respect to (z, u) . Although not standard, our assumptions allow to give a stochastic representation of a relevant class of Ergodic HJB equations in Hilbert spaces (see Section 5) and of ergodic stochastic control problems for SPDEs (see Example 7.1 and Example 7.2). Notice that ψ defined above is exactly the function that naturally appears in the related HJB equation and in the applications to ergodic control.

As in all the literature devoted to the problem the main point is to prove a uniform gradient estimate (independent on α) for $v^\alpha(x) := Y^{\alpha,x}$ where $(Y^{\alpha,x}, Z^{\alpha,x}, U^{\alpha,x})$ is the solution of the discounted BSDE with infinite horizon:

$$Y_t^{\alpha,x} = Y_T^{\alpha,x} + \int_t^T [\widehat{\psi}(X_s^x, Z_s^{\alpha,x}, U_s^{\alpha,x}) - \alpha Y_s^{\alpha,x}] ds - \int_t^T Z_s^{\alpha,x} dW_s^1 - \int_t^T U_s^{\alpha,x} dW_s^2, \quad 0 \leq t \leq T < \infty.$$

Such estimate can be obtained by a change of probability argument when the noise is additive (see [19]), by energy type estimates under quantitative assumptions on the exponential decay of the forward equation (see [24]) or by regularizing properties of the Kolmogorov semigroup when the noise is multiplicative but non degenerate (see [14] and [21]).

Here we exploit concavity of ψ to introduce an auxiliary control problem and eventually obtain the gradient estimate using a decay estimate on the difference between states starting from different initial conditions, see Assumption (A.6) and, in particular, requirement (3.5). We stress the fact that the estimate in (3.5) is only in mean and not uniform (with respect to the stochastic parameter) as in the additive noise case. Moreover, as we show in Proposition 3.2, Assumption (A.6) is verified if we impose a *joint dissipativity* condition on the coefficients, see Assumption (A.7). As a matter of fact, in this case, the stronger formulation in which L^2 replaces L^1 norm holds. On the other side (A.6) allows to cover a wider class of interesting examples, see for instance Example 7.1 in which Assumption (A.7) does not seem to hold.

The structure of the paper in the following: in Section 2 we introduce the function spaces that will be used in the following, Section 3 is devoted to the infinite dimensional forward equation; in particular we state and discuss the key stability assumption (A.6). In Section 4 we present the main contribution of this work introducing the auxiliary control problem, proving the gradient estimate and the consequent existence of the solution to the ergodic BSDEs. In Section 5 we relate our ergodic BSDE to a semilinear PDE in infinite dimensional spaces (the ergodic HJB equation). In Section 6 we discuss the regularity of the solution of the ergodic BSDE, in particular we state that under quantitative conditions on the dissipativity of the forward equation similar to the ones assumed in [24], when all coefficients are differentiable then the solution of the ergodic BSDE is differentiable with respect to the initial data as well. The proof of such result adapts a similar argument in [16] and is rather technical, we have postponed it in the Appendix Section 7 we use our ergodic BSDE

68 to obtain an optimal ergodic control problem (that is with cost depending only on the asymptotic
69 behaviour of the state) for an infinite dimensional equation. We close, see Section 7.1, by two
70 examples of controlled SPDEs to which our results can be applied. In both we consider a stochastic
71 heat equation in one dimension with additive white noise. In the first, Example 7.1 the system is
72 controlled through one Dirichlet boundary condition (on which multiplicative noise also acts) while,
73 in the second one, Example 7.2, the control enters the system through a finite dimensional process
74 that affects the coefficients of the SPDE. In this last case we also give conditions guaranteeing
75 differentiability of the related solution to the Ergodic BSDE.

76 **2. General notation.** Let Ξ , H and U be real separable Hilbert spaces. In the sequel, we use
77 the notations $|\cdot|_{\Xi}$, $|\cdot|_H$ and $|\cdot|_U$ to denote the norms on Ξ , H and U respectively; if no confusion
78 arises, we simply write $|\cdot|$. We use similar notation for the scalar products. We denote the dual
79 spaces of Ξ , H and U by Ξ^* , H^* , and U^* respectively. We also denote by $L(H, H)$ the space of
80 bounded linear operators from H to H , endowed with the operator norm. Moreover, we denote by
81 $L_2(\Xi, H)$ the space of Hilbert-Schmidt operators from Ξ to H . Finally, a map $f : H \rightarrow \Xi$ is said to
82 belong to the class $\mathcal{G}^1(H, \Xi)$ if it is continuous and Gateaux differentiable with directional derivative
83 $\nabla_x f(x)h$ in $(x, h) \in H \times H$ and we denote by $\mathcal{B}(\Lambda)$ the Borel σ -algebra of any topological space Λ .

84 Given a complete probability space $(\Omega, \mathcal{F}, \mathbb{P})$ together with a filtration $(\mathcal{F}_t)_{t \geq 0}$ (satisfying the
85 usual conditions of \mathbb{P} -completeness and right-continuity) and an arbitrary real separable Hilbert
86 space V we define the following classes of processes for fixed $0 \leq t \leq T$ and $p \geq 1$:

- 87 • $L_{\mathcal{P}}^p(\Omega \times [t, T]; V)$ denotes the set of (equivalence classes) of (\mathcal{F}_s) -predictable processes $Y \in$
88 $L^p(\Omega \times [t, T]; V)$ such that the following norm is finite:

89
$$|Y|_p = \left(\mathbb{E} \int_t^T |Y_s|^p ds \right)^{1/p}.$$

- 90 • $L_{\mathcal{P}}^{p,loc}(\Omega \times [0, +\infty[; V)$ denotes the set of processes defined on \mathbb{R}^+ , whose restriction to an
91 arbitrary time interval $[0, T]$ belongs to $L_{\mathcal{P}}^p(\Omega \times [0, T]; V)$.
92 • $L_{\mathcal{P}}^p(\Omega; C([t, T]; V))$ denotes the set of (\mathcal{F}_s) -predictable processes Y on $[t, T]$ with continuous
93 paths in V , such that the norm

94
$$\|Y\|_p = \left(\mathbb{E} \sup_{s \in [t, T]} |Y_s|^p \right)^{1/p}.$$

95 is finite. The elements of $L_{\mathcal{P}}^p(\Omega; C([t, T]; V))$ are identified up to indistinguishability.

- 96 • $L_{\mathcal{P}}^{p,loc}(\Omega; C([0, +\infty[; V))$ denotes the set of processes defined on \mathbb{R}^+ , whose restriction to an
97 arbitrary time interval $[0, T]$ belongs to $L_{\mathcal{P}}^p(\Omega; C([0, T]; V))$.

98 We consider on the probability space $(\Omega, \mathcal{F}, \mathbb{P})$ two independent cylindrical Wiener processes $W^1 =$
99 $(W_t^1)_{t \geq 0}$ with values in Ξ and $W^2 = (W_t^2)_{t \geq 0}$ with values in H . By $(\mathcal{F}_t)_{t \geq 0}$, we denote the natural
100 filtration of (W^1, W^2) , augmented with the family \mathcal{N} of \mathbb{P} -null sets of \mathcal{F} . The filtration (\mathcal{F}_t) satisfies
101 the usual conditions of right-continuity and \mathbb{P} -completeness.

102 **3. Forward equation.** Given $x \in H$ and a uniformly bounded progressively measurable
103 process \mathfrak{g} with values in H , we consider the stochastic differential equation for $t \geq 0$

104 (3.1)
$$dX_t^{x, \mathfrak{g}} = AX_t^{x, \mathfrak{g}} dt + F(X_t^{x, \mathfrak{g}}) dt + QG(X_t^{x, \mathfrak{g}}) dW_t^1 + DdW_t^2 + \mathfrak{g}(t) dt, \quad X_0^{x, \mathfrak{g}} = x.$$

105 On the coefficients A, F, G, Q, D we impose the following assumptions.

106 **(A.1)** $A: \mathcal{D}(A) \subset H \rightarrow H$ is a linear, possibly unbounded operator generating a C_0 semigroup
 107 $\{e^{tA}\}_{t \geq 0}$.

108 **(A.2)** $F: H \rightarrow H$ is continuous and there exists $L_F > 0$ such that

$$109 \quad |F(x) - F(x')|_H \leq L_F |x - x'|_H,$$

110 for all $x, x' \in H$.

111 **(A.3)** $G: H \rightarrow L(\Xi)$ is a bounded Lipschitz map. Moreover, for every $x \in H$, $G(x)$ is invertible.
 112 Thus there exists three positive constants L_G , M_G and $M_{G^{-1}}$ such that for all $x, x' \in H$:

$$113 \quad |G(x)|_{L(\Xi)} \leq M_G, \quad |G(x) - G(x')|_{L(\Xi)} \leq L_G |x - x'|_H, \quad |G^{-1}(x)|_{L(\Xi)} \leq M_{G^{-1}}.$$

114 We notice that the above yields Lipschitzianity of G^{-1} , namely :

$$115 \quad |G^{-1}(x) - G^{-1}(x')|_{L(\Xi)} \leq M_{G^{-1}}^2 L_G |x - x'|_H,$$

116 **(A.4)** Q is an Hilbert-Schmidt operator from Ξ to H .

117 **(A.5)** D is a linear and bounded operator from H to H and there exist constants $L > 0$ and
 118 $\gamma \in [0, \frac{1}{2}[$:

$$119 \quad (3.2) \quad |e^{sA}D|_{L_2(H)} \leq L (s^{-\gamma} \wedge 1), \quad \forall s \geq 0.$$

121 **PROPOSITION 3.1.** *Under **(A.1) – (A.5)**, for any $x \in H$ and any \mathfrak{g} bounded and progressively*
 122 *measurable process with values in H , there exists a unique (up to indistinguishability) process $X^{x, \mathfrak{g}} =$*
 123 *$(X_t^{x, \mathfrak{g}})_{t \geq 0}$ that belongs to $L_p^{p, loc}(\Omega; C([0, +\infty[; H))$ for all $p \geq 1$ and is a mild solution of (3.1), that*
 124 *is it satisfies for every $t \geq 0$, \mathbb{P} -a.s.:*

$$125 \quad X_t^{x, \mathfrak{g}} = e^{tA}x + \int_0^t e^{(t-s)A}F(X_s^{x, \mathfrak{g}}) ds + \int_0^t e^{(t-s)A}\mathfrak{g}(s) ds + \int_0^t e^{(t-s)A}QG(X_s^{x, \mathfrak{g}}) dW_s^1$$

$$126 \quad + \int_0^t e^{(t-s)A}D dW_s^2.$$

128 *Moreover there exists a positive constant $\kappa_{\mathfrak{g}, T}$ such that*

$$129 \quad (3.3) \quad \mathbb{E}|X_t^{x, \mathfrak{g}}|^2 \leq \kappa_{\mathfrak{g}, T}(1 + |x|^2), \quad \forall t \in [0, T] \text{ and } x \in H.$$

130 Our main result will be obtained under the following exponential stability in L^1 norm requirement.
 131 We stress the fact that such assumption is much weaker in comparison with the uniform decay
 132 holding when noise is additive (see [19]).

133 **(A.6)** There exist positive constants $\kappa_{\mathfrak{g}}$, κ and μ , independent from \mathfrak{g} , such that

$$134 \quad (3.4) \quad \sup_{t \geq 0} \mathbb{E}|X_t^{x, \mathfrak{g}}| \leq \kappa_{\mathfrak{g}}(1 + |x|);$$

135

$$136 \quad (3.5) \quad \mathbb{E}|X_t^{x, \mathfrak{g}} - X_t^{x', \mathfrak{g}}| \leq \kappa e^{-\mu t}|x - x'|;$$

137 for any $x, x' \in H$ and for all $t \geq 0$.

138 Below we show that hypothesis **(A.6)** (as a matter of fact the stronger condition obtained replacing
 139 L^1 norm by L^2 norm) is verified under the usual *joint dissipative condition* **(A.7)** (see [5]). We have
 140 preferred to keep the weaker, but less intrinsic, form **(A.6)** since it allows to cover a wider class of
 141 examples, see for instance Example 7.1.

142 **(A.7)** - *Joint dissipative conditions*

143 A is dissipative i.e. $\langle Ax, x \rangle \leq \rho|x|^2$, for all $x \in \mathcal{D}(A)$, and for some $\rho \in \mathbb{R}$, moreover there
 144 exists $\mu > 0$ such that for all $x, x' \in D(A)$:

$$145 \quad (3.6) \quad 2\langle A(x - x') + F(x) - F(x'), x - x' \rangle_H + \|Q[G(x) - G(x')]\|_{L_2(\Xi, H)}^2 \leq -\mu|x - x'|_H^2,$$

146 Notice that, by adding a suitable constant to F and subtracting it from A we can always
 147 assume that ρ above is strictly negative.

148 Indeed we have that following holds

149 **PROPOSITION 3.2.** *Assume **(A.1 – A.5)** and **(A.7)** then the following estimates hold for the*
 150 *solution $X^{x, \mathfrak{g}}$ of equation (3.1):*

$$151 \quad (3.7) \quad \sup_{t \geq 0} \mathbb{E}|X_t^{x, \mathfrak{g}}|^2 \leq \kappa_{\mathfrak{g}}(1 + |x|^2);$$

152

$$153 \quad (3.8) \quad \mathbb{E}|X_t^{x, \mathfrak{g}} - X_t^{x', \mathfrak{g}}|^2 \leq e^{-\mu t}|x - x'|^2;$$

154 for any $x, x' \in H$ and for all $t \geq 0$. In particular, hypothesis **(A.6)** is verified.

155 **Proof.**

156 The proof of these estimates follows rather standard arguments, see for instance [5] where
 157 dissipative systems are widely treated. \square

158 We end this section noticing that will be mainly interested in the special case where $\mathfrak{g} \equiv 0$:

$$159 \quad (3.9) \quad dX_t = AX_t dt + F(X_t)dt + QG(X_t)dW_t^1 + DdW_t^2, \quad X_0^x = x,$$

160 and we will denote by X^x its solution through the whole paper.

161 **4. Ergodic BSDEs .** In this section we study the following equation:

$$162 \quad (4.1) \quad Y_t^x = Y_T^x + \int_t^T [\psi(X_s^x, Z_s^x G^{-1}(X_s^x), U_s^x) - \lambda] ds - \int_t^T Z_s^x dW_s^1 - \int_t^T U_s^x dW_s^2, \quad 0 \leq t \leq T < \infty,$$

163 where, we recall, λ is a real number and it is part of the unknowns, and the equation has to hold
 164 for every t and every T , see for instance [19, section 4]. On the function $\psi : H \times \Xi^* \times H^* \rightarrow \mathbb{R}$ we
 165 assume:

166 **(B.1)** $(z, u) \rightarrow \psi(x, z, u)$ is a concave function at every fixed $x \in H$.

167 Moreover there exist $L_x, L_z, L_u > 0$ such that

$$168 \quad (4.2) \quad |\psi(x, z, u) - \psi(x', z', u')| \leq L_x|x - x'| + L_z|z - z'| + L_u|u - u'|, \quad x, x' \in H, z, z' \in \Xi^*, u, u' \in H^*.$$

169 Moreover $\psi(\cdot, 0, 0)$ is bounded. We denote $\sup_x |\psi(x, 0, 0)|$ by M_ψ .

170 We associate to ψ its Legendre transformation (modified according to the fact that we are dealing
 171 with concave functions):

$$172 \quad (4.3) \quad \psi^*(x, p, q) = \inf_{z \in \Xi^*, u \in H^*} \{-zp - uq - \psi(x, z, u)\}, \quad x \in H, p \in \Xi, q \in H.$$

173 Clearly ψ^* is concave w.r.t to (p, q) .

174 We collect some other properties of ψ and ψ^* we will use in the future:

175 PROPOSITION 4.1. Under hypothesis **(B.1)** we have that

$$176 \quad \psi(x, z, u) = \inf_{(p,q) \in \mathcal{D}^*(x)} \{-zp - uq - \psi^*(x, p, q)\}.$$

177 where $\mathcal{D}^*(x) = \{(p, q) : \psi^*(x, p, q) \neq -\infty\} \subset \{(p, q) \in \Xi \times H : |p| \leq L_z, |q| \leq L_u\}$.

178 Moreover $\mathcal{D}^*(x) = \mathcal{D}^*$ does not depend on $x \in H$ and the following holds

$$179 \quad (4.4) \quad |\psi^*(x, p, q) - \psi^*(x', p, q)| \leq L_x |x - x'|, \quad x, x' \in H, (p, q) \in \mathcal{D}^*.$$

180 Finally we remark that the above implies that for every $x \in H, z \in \Xi^*, u \in H^*$:

$$181 \quad \sup_{(p,q) \in \mathcal{D}} \{\psi(x, z, u) + zp + uq + \psi^*(x, p, q)\} = 0.$$

182 **Proof.** Since $\psi(x, \cdot, \cdot)$ is concave its double Legendre transform coincides with the function itself
183 and the first relation follows immediately (see [1]).

184 Then, by the definition of ψ^* :

$$185 \quad |\psi^*(x, p, q) - \psi^*(x', p, q)| \leq \sup_{z \in \Xi^*, u \in H^*} |-zp - uq - \psi(x, z, u) + zp + uq + \psi(x', z, u)| \leq L_x |x - x'|,$$

186 thus we deduce that \mathcal{D}^* doesn't depend on $x \in H$ and (4.4) holds. \square

188 As in [19] we introduce, for each $\alpha > 0$, the infinite horizon equation:

$$189 \quad (4.5) \quad Y_t^{x,\alpha} = Y_T^{x,\alpha} + \int_t^T [\psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) - \alpha Y_s^{x,\alpha}] ds - \int_t^T Z_s^{x,\alpha} dW_s^1 - \int_t^T U_s^{x,\alpha} dW_s^2,$$

190 where $0 \leq t \leq T < \infty$.

191 The next result was proved in [25, Theorem 2.1] in finite dimensions, the extension to the
192 infinite dimensional case is straightforward, see also [19, Lemma 4.2]. Notice that the random
193 function, $\widehat{\psi}(t, z, u) := \psi(X_t, G^{-1}(X_t)z, u)$, inherits the following properties:

$$194 \quad (4.6) \quad |\widehat{\psi}(t, 0, 0)| = |\psi(X_t, 0, 0)| \leq M_\psi, \quad t \geq 0, \mathbb{P}\text{-a.s.}$$

195

$$196 \quad (4.7) \quad |\widehat{\psi}(t, z, u) - \widehat{\psi}(t, z', u')| \leq L_z M_{G^{-1}} |z - z'| + L_u |u - u'| \quad t \geq 0, \quad z, z' \in \Xi^*, u, u' \in H^* .$$

197 therefore it satisfies the assumptions in [19, Lemma 4.2].

198 **THEOREM 4.1.** Let us assume **(A.1 – A.5)** and **(B.1)**. Then for every $\alpha > 0$ there exists a
199 unique solution $(Y^{x,\alpha}, Z^{x,\alpha}, U^{x,\alpha})$ to the BSDE (4.5) such that $Y^{x,\alpha}$ is a bounded continuous process,
200 $Z^{\alpha,x} \in L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; \Xi^*)$ and $U^{\alpha,x} \in L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; H^*)$.

201 Moreover

$$202 \quad (4.8) \quad |Y_t^{x,\alpha}| \leq \frac{M_\psi}{\alpha}, \quad \mathbb{P}\text{-a.s.}, \text{ for all } t \geq 0.$$

203 and

$$204 \quad (4.9) \quad \mathbb{E} \int_0^\infty |e^{-\alpha s} Z_s^{x,\alpha}|^2 ds + \mathbb{E} \int_0^\infty |e^{-\alpha s} U_s^{x,\alpha}|^2 ds < \infty.$$

205 We define

$$206 \quad (4.10) \quad v^\alpha(x) = Y_0^{\alpha,x}.$$

207 The following is the main estimate of the paper.

208 PROPOSITION 4.2. *Under (A.1 – A.6) and (B.1) one has that for any $\alpha > 0$:*

$$209 \quad (4.11) \quad |v^\alpha(x) - v^\alpha(x')| \leq \frac{C}{\mu} |x - x'|, \quad x, x' \in H.$$

210 where C depends on the constants in (A.1 – A.5) and (B.1) but not on α (nor on μ).

211 **Proof.** Since, instead of the pathwise decay estimate holding for $|X_t^x - X_t^{x'}|$ in the additive noise
 212 case (see [19, Theorem 3.2]), only the mean bound (3.5) is true here we cannot proceed as in [19,
 213 Theorem 4.4]. Moreover, being the diffusion X , in general, degenerate, it is not possible to rely
 214 on the smoothing properties of its Kolmogorov semigroup (see [21]). On the contrary, concavity
 215 assumption (B.1) allows us to use control theoretic arguments.

216 First we notice that

$$217 \quad Y_0^{x,\alpha} = e^{-\alpha t} Y_t^{x,\alpha} + \int_0^t e^{-\alpha s} \psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) ds - \int_0^t e^{-\alpha s} Z_s^{x,\alpha} dW_s^1 - \int_0^t e^{-\alpha s} U_s^{x,\alpha} dW_s^2.$$

218 Thus we have, taking also into account (4.8) and (4.9), that

$$219 \quad (4.12) \quad Y_0^{x,\alpha} = \int_0^{+\infty} e^{-\alpha s} \psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) ds - \int_0^{+\infty} e^{-\alpha s} Z_s^{x,\alpha} dW_s^1 - \int_0^{+\infty} e^{-\alpha s} U_s^{x,\alpha} dW_s^2.$$

220 Moreover being $Y_0^{x,\alpha}$ deterministic, the uniqueness in law for the system formed by equations (3.9)
 221 -(4.5) yields that it doesn't depend on the specific independent Wiener processes.

222 We fix any stochastic setting $(\hat{\Omega}, \hat{\mathcal{E}}, (\hat{\mathcal{F}}_t), \hat{\mathbb{P}}, (\hat{W}_t^1), (\hat{W}_t^2))$ where $((\hat{W}_t^1), (\hat{W}_t^2))$ are independent
 223 $(\hat{\mathcal{F}}_t)$ Wiener processes with values in Ξ and H respectively.

224 Given any $(\hat{\mathcal{F}}_t)$ progressively measurable process $\mathbf{p} := (p_t, q_t)$ with values in \mathcal{D}^* by $(\hat{X}_t^{x,\mathbf{p}})$ we
 225 denote the unique mild solution of the forward equation:

$$226 \quad (4.13) \quad d\hat{X}_t^{x,\mathbf{p}} = A\hat{X}_t^{x,\mathbf{p}} dt + F(\hat{X}_t^{x,\mathbf{p}}) dt + Dq_t dt + QG(\hat{X}_t^{x,\mathbf{p}}) p_t dt + QG(\hat{X}_t^{x,\mathbf{p}}) d\hat{W}_t^1 + Dd\hat{W}_t^2, \quad \hat{X}_0^{x,\mathbf{p}} = x.$$

227 Clearly $(\hat{X}_t^{x,\mathbf{p}})$ is also the unique mild solution of the forward equation:

$$228 \quad (4.14) \quad d\hat{X}_t^{x,\mathbf{p}} = A\hat{X}_t^{x,\mathbf{p}} dt + F(\hat{X}_t^{x,\mathbf{p}}) dt + QG(\hat{X}_t^{x,\mathbf{p}}) d\hat{W}_t^{1,\mathbf{p}} + Dd\hat{W}_t^{2,\mathbf{p}}, \quad \hat{X}_0^{x,\mathbf{p}} = x.$$

229 where

$$230 \quad (4.15) \quad \hat{W}_t^{1,\mathbf{p}} := \hat{W}_t^1 + \int_0^t G^{-1}(\hat{X}_s^{x,\mathbf{p}}) p_s ds, \quad \hat{W}_t^{2,\mathbf{p}} := \hat{W}_t^2 + \int_0^t q_s ds,$$

231 and we know that under a suitable probability $\hat{\mathbb{P}}^{\mathbf{p}}$ the processes $((\hat{W}_t^{1,\mathbf{p}}), (\hat{W}_t^{2,\mathbf{p}}))$ are independent
 232 Wiener processes with values in Ξ and H respectively.

233 Let now $(\hat{Y}^{x,\alpha,\mathbf{p}}, \hat{Z}^{x,\alpha,\mathbf{p}}, \hat{U}^{x,\alpha,\mathbf{p}})$ be the solution to:

$$234 \quad (4.16) \quad \hat{Y}_t^{x,\alpha,\mathbf{p}} = \hat{Y}_T^{x,\alpha,\mathbf{p}} + \int_t^T [\psi(\hat{X}_s^{x,\mathbf{p}}, \hat{Z}_s^{x,\alpha,\mathbf{p}} G^{-1}(\hat{X}_s^{x,\mathbf{p}}), \hat{U}_s^{x,\alpha,\mathbf{p}}) - \alpha Y_s^{x,\alpha,\mathbf{p}}] ds$$

236
237

$$- \int_t^T \hat{Z}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^{1,\mathbf{p}} - \int_t^T \hat{U}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^{2,\mathbf{p}},$$

238 where $0 \leq t \leq T < \infty$.

239 By previous considerations one has, recalling that $\{\psi(x, z) + zp + uq + \psi^*(x, p)\} \leq 0, \forall x \in H, z \in$
240 $\Xi^*, u \in H^*, (p, q) \in \mathcal{D}^*$, that for every $x \in H$

$$\begin{aligned} 241 \quad Y_0^{x,\alpha} &= \hat{Y}_0^{x,\alpha,\mathbf{p}} \\ 242 &= \int_0^\infty e^{-\alpha s} \left[\psi(\hat{X}_s^{x,\mathbf{p}}, \hat{Z}_s^{x,\alpha,\mathbf{p}} G^{-1}(\hat{X}_s^{x,\mathbf{p}}), \hat{U}_s^{x,\alpha,\mathbf{p}}) + \hat{Z}_s^{x,\alpha,\mathbf{p}} G^{-1}(\hat{X}_s^{x,\mathbf{p}}) p_s + \hat{U}_s^{x,\alpha,\mathbf{p}} q_s + \psi^*(\hat{X}_s^{x,\mathbf{p}}, p_s, q_s) \right] ds \\ 243 &- \int_0^{+\infty} e^{-\alpha s} \hat{Z}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^1 - \int_0^{+\infty} e^{-\alpha s} \hat{U}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^2 - \int_0^\infty \psi^*(\hat{X}_s^{x,\mathbf{p}}, p_s, q_s) ds \\ 244 &\leq - \int_0^{+\infty} e^{-\alpha s} \hat{Z}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^1 - \int_0^{+\infty} e^{-\alpha s} \hat{U}_s^{x,\alpha,\mathbf{p}} d\hat{W}_s^2 - \int_0^\infty \psi^*(\hat{X}_s^{x,\mathbf{p}}, p_s, q_s) ds. \end{aligned}$$

246 So:

$$247 \quad (4.16) \quad Y_0^{x,\alpha} \leq -\hat{\mathbb{E}} \int_0^\infty e^{-\alpha s} \psi^*(\hat{X}_s^{x,\mathbf{p}}, p_s, q_s) ds,$$

249 for arbitrary stochastic setting and arbitrary progressively measurable \mathcal{D}^* valued control $\mathbf{p} = (p, q)$.

250 Then we fix $x \in H$ and assume, for the moment, that $\forall \varepsilon > 0$ there exists a stochastic setting

$$(\hat{\Omega}^{\varepsilon,x}, \hat{\mathcal{E}}^{\varepsilon,x}, (\hat{\mathcal{F}}_t^{\varepsilon,x}), \hat{\mathbb{P}}^{\varepsilon,x}, (\hat{W}_t^{1,\varepsilon,x}), (\hat{W}_t^{2,\varepsilon,x})),$$

251 and a couple of predictable processes $\mathbf{p}^{\varepsilon,x} = (p^{\varepsilon,x}, q^{\varepsilon,x})$ with values in \mathcal{D}^* such that (with the
252 notations introduced above) the following holds \mathbb{P} - a.s. for a.e. $s \geq 0$:

$$\begin{aligned} 254 \quad (4.17) \quad \psi(\hat{X}_s^{x,\mathbf{p}^\varepsilon}, \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} G^{-1}(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}), \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}}) + \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} G^{-1}(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}) p_s^\varepsilon + \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} q_s^\varepsilon \\ 255 \quad + \psi^*(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, p_s^\varepsilon, q_s^\varepsilon) \geq -\varepsilon. \end{aligned}$$

257 Proceeding as before we get:

$$\begin{aligned} 258 \quad (4.18) \quad Y_0^{x,\alpha} &= \hat{Y}_0^{x,\alpha,\mathbf{p}^{\varepsilon,x}} = \\ 259 &= \int_0^\infty e^{-\alpha s} \left[\psi(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} G^{-1}(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}), \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}}) \right. \\ 260 &\quad \left. + \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} G^{-1}(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}) p_s^{\varepsilon,x} + \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} q_s^{\varepsilon,x} + \psi^*(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, p_s^{\varepsilon,x}, q_s^{\varepsilon,x}) \right] ds \\ 261 &- \int_0^{+\infty} e^{-\alpha s} \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} d\hat{W}_s^{1,\varepsilon,x} - \int_0^{+\infty} e^{-\alpha s} \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} d\hat{W}_s^{2,\varepsilon,x} - \int_0^\infty \psi^*(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, p_s^{\varepsilon,x}, q_s^{\varepsilon,x}) ds \\ 262 &\geq -\frac{\varepsilon}{\alpha} - \int_0^{+\infty} e^{-\alpha s} \hat{Z}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} d\hat{W}_s^{1,\varepsilon,x} - \int_0^{+\infty} e^{-\alpha s} \hat{U}_s^{x,\alpha,\mathbf{p}^{\varepsilon,x}} d\hat{W}_s^{2,\varepsilon,x} - \int_0^\infty \psi^*(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, p_s^{\varepsilon,x}, q_s^{\varepsilon,x}) ds. \blacksquare \end{aligned}$$

264 Thus by (4.16) taking into account (4.18) and (4.4) we have:

$$265 \quad Y_0^{x',\alpha} - Y_0^{x,\alpha} \leq \int_0^\infty e^{-\alpha s} \hat{\mathbb{E}}^{\mathbf{p}^{\varepsilon,x}} |\psi^*(\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}}, p_s^{\varepsilon,x}, q_s^{\varepsilon,x}) - \psi^*(\hat{X}_s^{x',\mathbf{p}^{\varepsilon,x}}, p_s^{\varepsilon,x}, q_s^{\varepsilon,x})| ds + \varepsilon$$

266
267

$$\leq \int_0^\infty e^{-\alpha s} \hat{\mathbb{P}}^{\mathbf{p}^{\varepsilon,x}} |\hat{X}_s^{x,\mathbf{p}^{\varepsilon,x}} - \hat{X}_s^{x',\mathbf{p}^{\varepsilon,x}}| ds + \frac{\varepsilon}{\alpha},$$

268 we stress the fact that we keep the stochastic setting $(\hat{\Omega}^{\varepsilon,x}, \hat{\mathcal{E}}^{\varepsilon,x}, (\hat{\mathcal{F}}_t^{\varepsilon,x}), \hat{\mathbb{P}}^{\varepsilon,x}, (\hat{W}_t^{1,\varepsilon,x}), (\hat{W}_t^{2,\varepsilon,x}))$ and
269 control $\mathbf{p}^{\varepsilon,x}$ corresponding to the initial datum x and just replace the initial state x with a different
270 one x' .

Noticing now that both $(\hat{X}^{x,\mathbf{p}^{\varepsilon,x}})$ and $(\hat{X}^{x',\mathbf{p}^{\varepsilon,x}})$ satisfy (only the initial conditions differ):

$$d\hat{X}_t = A\hat{X}_t dt + F(\hat{X}_t) dt + Dq_t^{\varepsilon,x} dt + QG(\hat{X}_t)p_t^{\varepsilon,x} dt + QG(\hat{X}_t^{\mathbf{p}})d\hat{W}_t^{1,\varepsilon,x} + Dd\hat{W}_t^{2,\varepsilon,x}.$$

and taking into account (3.5) we can conclude that:

$$Y_0^{x',\alpha} - Y_0^{x,\alpha} \leq L_x \int_0^\infty e^{-(\alpha+\frac{\mu}{2})s} |x - x'| ds + \frac{\varepsilon}{\alpha} \leq \frac{C}{\mu} |x - x'| + \frac{\varepsilon}{\alpha}.$$

271 Interchanging the role of x with x' one gets:

$$272 \quad (4.19) \quad |Y_0^{x,\alpha} - Y_0^{x',\alpha}| \leq \frac{C}{\mu} |x - x'| + \frac{\varepsilon}{\alpha}.$$

274 where the constant C is independent of α , μ and ε and is able to conclude (4.11) being $\varepsilon > 0$
275 arbitrary.

276 We are left with the construction, for any fixed $x \in H$ and $\varepsilon > 0$ of a stochastic setting
277 $(\hat{\Omega}^{\varepsilon,x}, \hat{\mathcal{E}}^{\varepsilon,x}, (\hat{\mathcal{F}}_t^{\varepsilon,x}), \hat{\mathbb{P}}^{\varepsilon,x}, (\hat{W}_t^{1,\varepsilon,x}), (\hat{W}_t^{2,\varepsilon,x}))$ and control $\mathbf{p}^{\varepsilon,x}$ for which (4.17) holds.

We start from an arbitrary stochastic setting: $(\Omega, \mathcal{E}, (\mathcal{F}_t), \mathbb{P}, (W_t^1), (W_t^2))$. Let (X^x) be the
corresponding mild solution of equation (3.1) and $(Y^{x,\alpha}, Z^{x,\alpha}, U^{x,\alpha})$ the solution of (4.5). By a
measurable selection argument see [22, Theorem 4] we can find a couple of progressive measurable
process $\mathbf{p}^{\varepsilon,x} = (p^{\varepsilon,x}, q^{\varepsilon,x})$, (possibly depending on α as well), such that:

$$\psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) + Z_s^{x,\alpha} G^{-1}(X_s^x) p_s^{\varepsilon,x} + U_s^{x,\alpha} q_s^{\varepsilon,x} + \psi_*(X_s^x, p_s^{\varepsilon,x}, q_s^{\varepsilon,x}) \geq -\varepsilon.$$

278 Then it is enough to set:

$$279 \quad (4.20) \quad \hat{W}_t^{1,\varepsilon,x} := W_t^1 - \int_0^t G^{-1}(X_s^x) p_s^{\varepsilon,x} ds, \quad \hat{W}_t^{2,\varepsilon,x} := W_t^2 - \int_0^t q_s^{\varepsilon,x} ds,$$

280 and choose $\hat{\Omega}^{\varepsilon,x} = \Omega$, $\hat{\mathcal{E}}^{\varepsilon,x} = \mathcal{E}$, $(\hat{\mathcal{F}}_t^{\varepsilon,x}) = (\mathcal{F}_t)$ and as $\hat{\mathbb{P}}^{\varepsilon,x}$ the (unique) probability measure under
281 which $((\hat{W}_t^{1,\varepsilon,x}), (\hat{W}_t^{2,\varepsilon,x}))$ are independent Wiener processes. The claim then follows selecting the
282 above control $\mathbf{p}^{\varepsilon,x}$ and noticing that, by construction, $(\hat{X}^{x,\mathbf{p}^{\varepsilon,x}}) = (X^x)$. \square

283 Following [19] we can find a function \bar{v} and a number λ such that:

$$284 \quad (4.21) \quad [v^{\alpha_m}(x) - v^{\alpha_m}(0)] \rightarrow \bar{v}(x), \quad \forall x \in H,$$

285

$$286 \quad (4.22) \quad \alpha_m v^{\alpha_m}(0) \rightarrow \lambda.$$

287 where $\{\alpha_m\}_{m \in \mathbb{N}}$ is a suitable subsequence constructed using a diagonal method.

288 We can then proceed as in [19] to deduce from above the existence of a solution to (4.1) and the
289 uniqueness of λ .

290 THEOREM 4.2. Assume (A.1) – (A.6) and (B.1), let λ the number defined in (4.22) and set
 291 $\bar{Y}_t^x := \bar{v}(X_t^x)$, where \bar{v} is defined in (4.21). Then there exists \bar{Z}^x in $L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; \Xi^*)$ and \bar{U}^x
 292 in $L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; H^*)$ such that $(\bar{Y}^x, \bar{Z}^x, \bar{U}^x, \lambda)$ solves equation (4.1), \mathbb{P} -a.s. for all $0 \leq t \leq T$.

293 Moreover suppose that another quadruple (Y', Z', U', λ) where Y' is a progressively measurable
 294 continuous process verifying $|Y'_t| \leq c(1 + |X_t^x|)$, $Z' \in L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; \Xi^*)$, $U' \in L_{\mathcal{P}}^{2,loc}(\Omega \times$
 295 $[0, +\infty[; H^*)$ and $\lambda' \in \mathbb{R}$, satisfies (4.1). Then $\lambda' = \lambda$.

296 Finally there exists a measurable function $\bar{\zeta} : H \rightarrow \Xi^* \times H^*$ such that $(\bar{Z}_t^x, \bar{U}_t^x) = \bar{\zeta}(X_t^x)$.

297 **Proof.**

298 Once (4.11), (4.21) and (4.22) are obtained, the proof as far the first two statements is concerned
 299 follows exactly as in [19, Theorem 4.4].

300 To get the existence of a function $\bar{\zeta}$, we proceed in the following way. For arbitrary fixed
 301 $0 \leq t \leq T$ let $(\bar{Y}^{x,t,T}, \bar{Z}^{x,t,T}, \bar{U}^{x,t,T})$ be the solution to:

$$302 \quad (4.23) \quad \begin{cases} dX_s^{t,x} = AX_s^{t,x} ds + F(X_s^{t,x}) ds + QG(X_s^{t,x}) dW_s^1 + DdW_s^2, \\ X_t^{t,x} = x, \\ -dY_s^{x,t,T} = \widehat{\psi}(X_s^{x,t}, Z_s^{x,t,T}, U_s^{x,t,T}) ds - Z_s^{x,t,T} dW_s^1 - U_s^{x,t,T} dW_s^2 - \lambda ds, \\ Y_T^{x,t,T} = \bar{v}(X_T^{x,t}). \end{cases}$$

303 Then we clearly have that $(\bar{Y}^x, \bar{Z}^x, \bar{U}^x)$, restricted on $[0, T]$, coincide with $(\bar{Y}^{x,0,T}, \bar{Z}^{x,0,T}, \bar{U}^{x,0,T})$, for
 304 all $T > 0$. By [8, Prop. 3.2] we know that there exists a measurable function $\zeta^T : [0, T] \times H \rightarrow \Xi^* \times$
 305 H^* , such that $(\bar{Z}_s^{x,t,T}, \bar{U}_s^{x,t,T}) = \zeta^T(s, X_s^{x,t})$, $s \in [t, T]$. Moreover, see also [8, Remark 3.3], the map
 306 $[0, T] \ni (\tau, x) \rightarrow \zeta^T(\tau, x)$ is characterized in terms of the laws of $(\int_{\tau}^{\tau+\frac{1}{n}} \bar{Z}_s^{\tau,x,T} ds, \int_{\tau}^{\tau+\frac{1}{n}} \bar{U}_s^{\tau,x,T} ds)$,
 307 $n \in \mathbb{N}$.

308 The uniqueness in law of the solutions to the system (4.23) together with the fact that its
 309 coefficients are time autonomous, we get:

$$310 \quad \int_{\tau}^{\tau+\frac{1}{n}} \bar{Z}_s^{\tau,x,T} ds \sim \int_0^{\frac{1}{n}} \bar{Z}_s^{0,x,T-\tau} ds \sim \int_0^{\frac{1}{n}} \bar{Z}_s^x ds,$$

311 and

$$312 \quad \int_{\tau}^{\tau+\frac{1}{n}} \bar{U}_s^{\tau,x,T} ds \sim \int_0^{\frac{1}{n}} \bar{U}_s^{0,x,T-\tau} ds \sim \int_0^{\frac{1}{n}} \bar{U}_s^x ds.$$

313 So far we've proved that $\zeta^T(\tau, \cdot)$ does not depend neither from T nor from τ , thus we can define
 314 $\zeta^T(\tau, \cdot) =: \bar{\zeta}(\cdot)$ and observe that $(\bar{Z}_t^x, \bar{U}_t^x) = (\bar{Z}_t^{x,0,T}, \bar{U}_t^{x,0,T}) = \zeta^T(t, X_t^{x,0}) = \bar{\zeta}(X_t^x)$. \square

315 **REMARK 4.1.** Concerning the uniqueness of the Markovian solution to the Ergodic BSDE (4.1)
 316 and consequently of the mild solution to the ergodic HJB equation (5.1) only partial results are
 317 available even in the additive case (beside the obvious consideration that adding a constants to Y
 318 and consequently to v transforms solutions into solutions). In particular an argument based on
 319 recurrence of the solution X to (1.1) is developed in [12] (see also [19]) to obtain a control theoretic
 320 representation of v and consequently its uniqueness up to an additive constant. Such arguments seem
 321 inapplicable in the present context due to possible degeneracy of the noise.

322 **5. Ergodic Hamilton-Jacobi-Bellman.** We wish now to prove that function \bar{v} satisfies, in
 323 a suitable way, the following Hamilton Jacobi Bellman elliptic partial differential equation:

$$324 \quad (5.1) \quad \frac{1}{2} \text{tr}[QG(x)G^*(x)Q\nabla^2 \bar{v}(x)] + \frac{1}{2} \text{tr}[DD^*(x)Q\nabla^2 \bar{v}(x)] + \langle Ax + F(x), \nabla \bar{v}(x) \rangle =$$

328 Since the proof of differentiability of \bar{v} requires quantitative conditions that we were able to avoid in
 329 Theorem 4.2 we firstly formulate the PDE in a weaker sense involving the *Generalized directional*
 330 *gradient* introduced in [11]. The following is the version of Theorem 3.1 in [11] adapted to the
 331 present autonomous and Lipschitz case. The proof is identical to the one in [11] and is omitted.

332 **THEOREM 5.1.** *Given any Lipschitz function v on H there exists a couple of bounded and Borel*
 333 *measurable functions $\zeta^1 : H \rightarrow \Xi^*$, $\zeta^2 : H \rightarrow H^*$ such that denoting, for all $\xi = (\xi^1, \xi^2) \in \Xi \times H$,*
 334 *by $W_s^\xi := \langle (W_s^1, W_s^2), \xi \rangle$ the real Brownian Motion obtained projecting (W_s^1, W_s^2) along direction ξ ,*
 335 *then we have the following relation, for any $x \in H$ and any $\rho > 0$*

$$336 \quad \langle v(X_t^x), W_t^\xi \rangle_{[0, \rho]} = \int_0^\rho \zeta^1(X_t^x) \xi^1 dt + \int_0^\rho \zeta^2(X_t^x) \xi^2 dt, \quad \mathbb{P} - a.s.$$

337 **DEFINITION 5.1.** *The family of functions $\zeta = (\zeta^1, \zeta^2)$ satisfying the above will be called the*
 338 *generalized (QG, D) directional gradients of u (denoted by $\tilde{\nabla}^{QG, D}$).*

339 **REMARK 5.1.** *Concerning uniqueness we can only say that if ζ and $\hat{\zeta}$ both belong to $\tilde{\nabla}^{QG, D}$*
 340 *then $\zeta^1(X_t^x) = \hat{\zeta}^1(X_t^x)$ and $\zeta^2(X_t^x) = \hat{\zeta}^2(X_t^x)$, \mathbb{P} -a.s. for almost every $t \geq 0$. See [11]. It is*
 341 *also clear that, by Ito rule, if u is regular enough, including twice continuously differentiable, then*
 342 *$(\nabla u(\cdot)QG(\cdot), \nabla u(\cdot)D)$ is in $\tilde{\nabla}^{QG, D}$.*

343 We are therefore led to the following definition of generalized solution to HJB equation. see [11,
 344 Section 5]:

345 **DEFINITION 5.2.** *A pair (v, λ) is a mild solution in the sense of generalized directional gradient*
 346 *of the HJB equation (5.1) if $v : H \rightarrow \mathbb{R}$ is Lipschitz and, for every $T > 0$ and for all $0 \leq t \leq T$ and*
 347 *$x \in H$ it holds*

$$348 \quad (5.2) \quad v(x) = P_{T-t}[v](x) + \int_t^T (P_{s-t}[\psi(\cdot, \zeta^1(\cdot)G^{-1}, \zeta^2(\cdot))](x) - \lambda) ds.$$

349 where $\zeta = (\zeta^1, \zeta^2)$ is an arbitrary element of the generalized gradient $\tilde{\nabla}^{(QG, D)}$ and $(P_t)_{t \geq 0}$ is the
 350 transition semigroup corresponding to the diffusion X^x , see equation (3.9), that is:

$$351 \quad (5.3) \quad P_t[\phi](x) := \mathbb{E} \phi(X_t^x), \quad \phi : H \rightarrow \mathbb{R} \text{ measurable and bounded.}$$

We notice that function \bar{v} defined in (4.21) is Lipschitz. Moreover recalling, see Theorem 4.2, that
 $(\bar{Y}_t^x, \bar{Z}_t^x, \bar{U}_t^x) = (\bar{v}(X_t^x), \bar{\zeta}^1(X_t^x), \bar{\zeta}^2(X_t^x))$ we have that then equation (4.1) is satisfied, in particular,
 for $t = 0$ and all $T > 0$ we immediately deduce that $\bar{\zeta} = (\bar{\zeta}^1, \bar{\zeta}^2)$ is in $\tilde{\nabla}^{(QG, D)}$. Finally recalling
 once more equation (4.1) now interpreted as a finite horizon BSDE:

$$-d\bar{Y}_t^x = \psi(X_s^x, \bar{Z}_s^x G^{-1}(X_s^x), \bar{U}_s^x) ds - \lambda ds - \bar{Z}_s^x dW_s^1 - \bar{U}_s^x dW_s^2, \quad \bar{Y}_T^x = \bar{v}(X_T^x)$$

352 we can conclude the following, proceeding exactly as in [11] Theorem 5.1

353 **THEOREM 5.2.** *Assume (A.1 – A.6), (B.1) then the couple (\bar{v}, λ) , characterized in Theorem*
 354 *4.2, is a mild solution, in the sense of the generalized directional gradient of equation (5.1).*

355 Whenever \bar{v} is differentiable then we can switch to the more classical notion of mild solution to
 356 equation (5.1) (\bar{v}, λ) , see [20, Section 6]:

357 DEFINITION 5.3. A pair (v, λ) is a mild solution to the HJB equation (5.1) if $v \in \mathcal{G}^1(H, \mathbb{R})$ with
 358 bounded derivative and, for all $0 \leq t \leq T$, $x \in H$ it holds:

$$359 \quad (5.4) \quad v(x) = P_{T-t}[v](x) + \int_t^T (P_{s-t}[\psi(\cdot, \nabla v(\cdot)Q, \nabla v(\cdot)D)](x) - \lambda) ds.$$

360 We have the following result.

361 THEOREM 5.3. Assume (A.1 – –A.6), (B.1) and that \bar{v} is of class \mathcal{G}^1 . Then (\bar{v}, λ) , defined in
 362 (4.21) is a mild solution of the HJB equation (5.1). On the other hand if (v', λ') is a mild solution
 363 of (5.1) then setting $Y_t^x := v'(X_t^x)$, $Z_t^x = \nabla v'(X_t^x)QG(X_t^x)$ and $U_t^x = \nabla v'(X_t^x)D$, we obtain that
 364 (Y^x, Z^x, U^x, λ) is a solution to equation (4.1).

365 Moreover if (v', λ') is another solution with v' Gateaux differentiable with linear growth then
 366 $\lambda = \lambda'$.

367 **Proof.** The existence part follows from [10, Theorem 6.2], while the uniqueness of λ in the class of
 368 solutions that are Gateaux differentiable with linear growth follows as [20, Theorem 4.6]. \square

369 REMARK 5.2. The differentiability of function \bar{v} is proved in Theorem 6.1 under quantitative
 370 assumptions on the coefficients. Although the argument essentially follows the classical paths of L^2
 371 estimates on infinite horizon see, for instance, [6] it is not completely standard since exploits in
 372 several points an a priori L^∞ estimate on Z and U descending from Proposition 4.2. In particular
 373 the uniform bounds for Z is essential in getting (A.9).

374 We conclude this section proving the following asymptotic expansion result for parabolic solu-
 375 tions to the HJB equation.

376 PROPOSITION 5.1. Let $v(\cdot, \cdot)$ be a mild solution of the parabolic HJB equation:

$$377 \quad (5.5) \quad \begin{cases} \partial_t v(t, x) = \frac{1}{2} [tr[QG(x)G^*(x)Q\nabla^2 v(t, x)] + tr[DD^*(x)Q\nabla^2 v(t, x)] + \langle Ax + F(x), \nabla v(t, x) \rangle \\ \quad + \psi(x, \nabla v(t, x)Q, \nabla v(t, x)D), \\ v(0, x) = \phi(x). \end{cases}$$

378 where $\phi : H \rightarrow \mathbb{R}$ is function of class \mathcal{G}^1 with bounded derivative and by mild solution of equation
 379 5.5 we mean a function $v : \mathbb{R}^+ \times H \rightarrow \mathbb{R}$ of class $\mathcal{G}^{0,1}$ (see [10]) verifying for all $t > 0$, $x \in H$:

$$380 \quad (5.6) \quad v(t, x) = P_t[\phi](x) + \int_0^t P_{t-s}[\psi(\cdot, \nabla v(\cdot)Q, \nabla v(\cdot)D)](x) ds.$$

381 Then

$$382 \quad (5.7) \quad \lim_{T \rightarrow \infty} \frac{v(T, x)}{T} = \lambda.$$

383 **Proof.** We fix $T > 0$ and consider the following finite horizon BSDE

$$384 \quad (5.8) \quad \begin{cases} -d\bar{Y}_s^{T,x} = \psi(X_s^x, G^{-1}(X_s^x)Z_s^{T,x}, U_s^{T,x}) ds - Z_s^{T,x} dW_s^1 - U_s^{T,x} dW_s^2 - \lambda ds \\ \bar{Y}_T^{T,x} = \phi(X_T^x). \end{cases}$$

385 By standard results on finite horizon BSDEs and mild solution of parabolic Hamilton-Jacobi-Bellman
 386 equations (see [10]) we have that $\bar{Y}_s^{T,x} = v(T-s, X_s^x) - \lambda(T-s)$, $s \in [0, T]$.

387 Set $\tilde{Y}_t^{T,x} = \bar{Y}_t^x - \bar{Y}_t^{T,x}$, for all $t \in [0, T]$, then $\tilde{Y}^{T,x}$ verifies:

388 (5.9)
$$\begin{cases} -d\tilde{Y}_s^{T,x} = [\psi(X_s^x, G^{-1}(X_s^x)\bar{Z}_s^x, \bar{U}_s^x) - \psi(X_s^{x,t}, G^{-1}(X_s^x)Z_s^{T,x}, U_s^{T,x})] ds - (\bar{Z}_s^x - Z_s^{T,x}) dW_s^1 \\ \quad - (\bar{U}_s^x - U_s^{T,x}) dW_s^2 - \lambda ds \\ \tilde{Y}_T^{T,x} = \bar{v}(X_T^x) - \phi(X_T^x). \end{cases}$$

389 We rewrite (5.9) as:

390 (5.10)
$$\begin{cases} -d\tilde{Y}_s^{T,x} = \gamma_t^1(\bar{Z}_s^x - Z_s^{T,x}) ds + \gamma_t^2(\bar{U}_s^x - U_s^{T,x}) ds - (\bar{Z}_s^x - Z_s^{T,x}) dW_s^1 \\ \quad - (\bar{U}_s^x - U_s^{T,x}) dW_s^2 - \lambda ds \\ \tilde{Y}_T^{T,x} = \bar{v}(X_T^x) - \phi(X_T^x). \end{cases}$$

391 where γ^1 and γ^2 are the typical uniformly bounded processes that arise in the linearization trick.

392 Hence, by a Girsanov argument, we get that

393 (5.11)
$$\tilde{Y}_0^{T,x} = \mathbb{E}^{\gamma^1, \gamma^2}(\bar{v}(X_T^x)),$$

394 where the probability measure $\mathbb{P}^{\gamma^1, \gamma^2}$ is the one under which $W_t^{\gamma^1, \gamma^2} = (W_t^1 - \int_0^t \gamma_s^1 ds, W_t^2 - \int_0^t \gamma_s^2 ds)$
395 is a cylindrical Wiener process in $\Xi \times H$ in $[0, T]$. Therefore by (3.4) and having \bar{v} Lipschitz, we get
396 that

397 (5.12)
$$\tilde{Y}_0^{T,x} = \mathbb{E}^{\gamma^1, \gamma^2}(\bar{v}(X_T^x) - \phi(X_T^x)) \leq \kappa_{\gamma^1, \gamma^2}(1 + |x|),$$

398 for some constant $\kappa_{\gamma^1, \gamma^2}$ independent of T . Thus, noticing that $\tilde{Y}_0^{T,x} = \bar{v}(x) - v(T, x) + \lambda T$ we get
399 that:

400 (5.13)
$$\lim_{T \rightarrow \infty} \frac{v(T, x)}{T} = \lim_{T \rightarrow \infty} \frac{\bar{v}(x)}{T} + \lambda = \lambda.$$

401 □

402 **REMARK 5.3.** *A more precise description of the asymptotic behaviour of $\frac{v(T, x)}{T}$ is obtained in*
403 *[15] when the noise is non degenerate by techniques involving Girsanov change of probability and*
404 *coupling estimates. Due to the possible non-invertibility of Q we do not know whether similar results*
405 *can be true in the present framework. We do think that, in any case, the proof of such results would*
406 *require different arguments.*

407 **6. Differentiability with respect to initial data.** In this section we wish to present suffi-
408 cient conditions under which the function \bar{v} defined in the section above is differentiable.

409 Throughout the section we assume the following:

410 **(C.1)** F is of class $\mathcal{G}^1(H, H)$ and G is of class $\mathcal{G}^1(H, L(\Xi, H))$

411 We start from a straightforward result in the non-degenerate case.

412 **PROPOSITION 6.1.** *Beside (A.1 – A.6), (B.1) and (C.1) assume that the operator $\mathcal{Q} :=$*
413 *$(\mathcal{Q}, D) : \Xi \times H \rightarrow H$ admits a right inverse \mathcal{Q}^{-1} then \bar{v} belongs to class $\mathcal{G}^1(H)$.*

Proof. We fix $T > 0$ and notice that $(\bar{Y}, \bar{Z}, \bar{U}, \lambda)$ satisfies (see (4.1) and the definition of \bar{Y}_t in Theorem 4.2):

$$Y_t^x = \bar{v}(X_T^x) + \int_t^T [\hat{\psi}(X_s^x, \bar{Z}_s^x, \bar{U}_s^x) - \lambda] ds - \int_t^T \bar{Z}_s^x dW_s^1 - \int_t^T \bar{U}_s^x dW_s^2, \quad 0 \leq t \leq T < \infty,$$

where, we recall $\widehat{\psi}(x, z, u) = \psi(x, zG^{-1}(x), u)$ is lipschitz with respect to z and u . Moreover the forward equation (3.9) solved by X^x can be rewritten as

$$dX_t^x = AX_t^x dt + F(X_t^x)dt + \tilde{Q}(X_t^x)d\mathcal{W}_t, \quad X_0^x = x.$$

414 where $\mathcal{W}_t := \begin{pmatrix} W_t^1 \\ W_t^2 \end{pmatrix}$ is a $\Xi \times H$ valued Wiener process and $\tilde{Q}(x) = (QG(x), D)$.

Under the present assumptions $\tilde{Q}(x)$ turns out to be invertible with bounded right inverse:

$$[\tilde{Q}(x)]^{-1} = \begin{pmatrix} G^{-1}(x) & 0 \\ 0 & I \end{pmatrix} Q^{-1}.$$

415 It is then straight forward to verify that all the assumptions in [9, Theorem 3.10] are satisfied and
416 consequently \bar{v} (that coincides with the map $x \rightarrow \bar{Y}_x$) is in class \mathcal{G}^1 \square

417 When the noise in the diffusion can be degenerate the situation is less simple and we will need
418 quantitative conditions on the coefficients (see, for instance, [24]).

419 We will now work under the *joint dissipative condition* (A.7) that, taking into account differentia-
420 bility of F and G becomes:

$$421 \quad (6.1) \quad 2\langle Ay + \nabla_x F(x)y, y \rangle_H + \|Q\nabla_x G(x)y\|_{L_2(\Xi, H)}^2 \leq -\mu|y|_H^2, \quad \forall y \in D(A), \forall x \in H.$$

422 Under the above assumptions the following well known differentiability result for the forward
423 equation (3.1) holds:

424 LEMMA 6.1. *Under (A.1 – A.5), (A.7) and (C.1) the map $x \rightarrow X^x$ is Gâteaux differentiable.*
425 *Moreover, for every $h \in H$, the directional derivative process $\nabla_x X^x h$, solves, \mathbb{P} - a.s., the equation*
426 (6.2)

$$426 \quad \nabla_x X_t^x h = e^{tA}h + \int_0^t e^{(t-s)A} \nabla_x F(X_s^x) \nabla_x X_s^x h ds + \int_0^t e^{(t-s)A} Q \nabla_x G(X_s^x) \nabla_x X_s^x h dW_s, \quad t \geq 0,$$

427 *Moreover*

$$428 \quad (6.3) \quad \mathbb{E}|\nabla_x X_t^x h|^2 \leq e^{-\mu t}|h|^2.$$

429 **Proof.** Our hypotheses imply the Hypotheses 3.1 of [10], therefore we can apply [10, Prop 3.3]. The
430 estimate (6.3) follows applying the Itô formula to $|\nabla_x X_t^x h|^2$ and arguing as in Proposition 3.1. \square

431

432 We will need the following additional assumption to state the last result

433 (C.2) G and G^{-1} are of class $\mathcal{G}^1(H, L(\Xi))$ and ψ is of class $\mathcal{G}^1(H \times \Xi^*, \mathbb{R})$

434 We eventually have:

435 THEOREM 6.1. *Assume that (A.1 – A.5), (A.7) and (B.1) hold with $\mu > 2(L_z^2 M_{G^{-1}}^2 + L_u^2)$,*
436 *moreover we assume (C.1) and (C.2). Then the function \bar{v} defined in (4.21) is of class $\mathcal{G}^1(H, \mathbb{R})$.*

437 **Proof.** The proof is detailed in the Appendix. \square

438 **7. Application to optimal control.** Let Γ be a separable metric space, an admissible control
439 γ is any \mathcal{F}_t - progressively measurable Γ -valued process. The cost corresponding to a given control
440 is defined as follows. Let $R_1 : \Gamma \rightarrow \Xi$, $R_2 : \Gamma \rightarrow H$ and $L : H \times \Gamma \rightarrow \mathbb{R}$ measurable functions such
441 that, for some constant $c > 0$, for all $x, x' \in H$ and $\gamma \in \Gamma$:

442 **(E.1)** $|R_1(\gamma)| \leq c, \quad |R_2(\gamma)| \leq c, \quad |L(x, \gamma)| \leq c, \quad |L(x, \gamma) - L(x', \gamma)| \leq c|x - x'|.$

443

444 Let for every $x \in H$ be X^x the solution to (3.9), then for every $T > 0$ and every control γ we
445 consider the Girsanov density:

446
$$\rho_T^\gamma = \exp \left(\int_0^T G^{-1}(X_s^x) R_1(\gamma_s) dW_s^1 + \int_0^T R_2(\gamma_s) dW_s^2 - \frac{1}{2} \int_0^T [|G^{-1}(X_s^x) R_1(\gamma_s)|_\Xi^2 + |R_2(\gamma_s)|_H^2] ds \right)$$

447 and we introduce the following ergodic cost corresponding to x and γ :

448
$$J(x, \gamma) = \limsup_{t \rightarrow \infty} \frac{1}{T} \mathbb{E}^{\gamma, T} \int_0^T L(X_s^x, \gamma_s) ds,$$

where $\mathbb{E}^{\gamma, T}$ is the expectation with respect to $\mathbb{P}^\gamma := \rho_T^\gamma \mathbb{P}$. Notice that with respect to \mathbb{P}^γ the processes

$$W_t^{1, \gamma} := - \int_0^t G^{-1}(X_s^x) R_1(\gamma_s) ds + dW_s^1, \quad W_t^{2, \gamma} := - \int_0^t R_2(\gamma_s) ds + dW_s^2,$$

are independent cylindrical Wiener processes and with respect to them X^x verifies:

$$\begin{cases} dX_t^x = AX_t^x dt + F(X_t^x) dt + QR_1(\gamma_s) ds + DR_2(\gamma_s) ds + QG(X_t^x) dW_t^{1, \gamma} + DdW_t^{2, \gamma}, & t \geq 0, \\ X_0^x = x, \end{cases}$$

449 and this justifies the above (weak) formulation of the control problem.

450 We introduce the *usual* Hamiltonian:

451 (7.1)
$$\psi(x, z, u) = \inf_{\gamma \in \Gamma} \{L(x, \gamma) + zR_1(\gamma) + uR_2(\gamma)\}, \quad x \in H, z \in \Xi^*, u \in H^*,$$

452 that by construction is a concave function and, under **(E.1)**, fullfills assumption **(B.1)**. The forward
453 backward system associated to this problem, is the following:

454 (7.2)
$$\begin{cases} dX_t^x = AX_t^x dt + F(X_t^x) dt + QG(X_t^x) dW_t^1 + DdW_t^2, & t \geq 0, \\ X_0^x = x, \\ -dY_t^x = [\psi(X_t^x, Z_t^x G^{-1}(X_t^x), U_t^x) - \lambda] dt - Z_t^x dW_t^1 - U_t^x dW_t^2. \end{cases}$$

455 By Theorem 4.2 under **(A.1 – A.6)** and **(E.1)** for every $x \in H$ there exists a solution:

456 (7.3)
$$(\bar{Y}^x, \bar{Z}^x, \bar{U}^x, \lambda) = (\bar{v}(X^x), \bar{\zeta}_1(X^x), \bar{\zeta}_2(X^x), \lambda),$$

457 where \bar{Y} is a progressive measurable continuous process, $\bar{Z} \in L_p^{2, loc}(\Omega \times [0, +\infty[; \Xi^*)$, $\bar{U} \in L_p^{2, loc}(\Omega \times$
458 $[0, +\infty[; H^*)$, $\lambda \in \mathbb{R}$, \bar{v} is Lipschitz and $\bar{\zeta}_1, \bar{\zeta}_2$ are measurable.

459 Once we have solved the above ergodic BSDE the proof of the following result containing the
460 synthesis of the optimal control for the ergodic cost is identical to the one of [19, Theorem 7.1].

461 **THEOREM 7.1.** *Assume **(A.1 – A.6)** and **(E.1)**. Then the following holds:*

(i) *For arbitrary control γ we have $J(x, \gamma) \geq \lambda$, and equality holds if and only if the following holds \mathbb{P} - a.s. for a.e. $t \geq 0$:*

$$L(X_t^x, \gamma_t) + \bar{\zeta}_1(X_t^x) G^{-1}(X_t^x) R_1(\gamma_t) + \bar{\zeta}_2(X_t^x) R_2(\gamma_t) = \psi(X_t^x, \bar{\zeta}_1(X_t^x) G^{-1}(X_t^x), \bar{\zeta}_2(X_t^x)).$$

- 462 (ii) If the infimum is attained in (7.1) and $\rho : \Xi^* \times H^* \rightarrow \Gamma$ is any measurable function realizing
463 the minimum (that always exists by Filippov selection theorem, see [22]) then the control
464 $\bar{\gamma}_t = \rho(X_t^x, \bar{\zeta}_1(X_t^x), \bar{\zeta}_2(X_t^x))$ is optimal, that is $J(x, \bar{\gamma}) = \lambda$.
465 (iii) \bar{v} admits a generalized directional gradient and (\bar{v}, λ) is the mild solution of the equation
466 (5.1), in the sense of definition (5.2) and $\bar{\zeta}_1, \bar{\zeta}_2 \in \tilde{\nabla}^{(QG, D)}$.
467 (iv) Finally if \bar{v} is in class \mathcal{G}^1 then (\bar{v}, λ) is a mild solution of equation (5.1), in the sense of
468 definition (5.3) and $\bar{\zeta}_1 = \nabla \bar{v} QG$ and $\bar{\zeta}_2 = \nabla \bar{v} D$.

469 7.1. Examples.

470 EXAMPLE 7.1. We consider an ergodic control problem for a stochastic heat equation controlled
471 through the boundary

$$472 \quad (7.4) \quad \begin{cases} d_t x(t, \xi) = \frac{\partial}{\partial \xi^2} x(t, \xi) dt + d(\xi) \dot{W}(t, \xi) dt, & t \geq 0, \xi \in (0, \pi), \\ x(t, 0) = y(t), \quad x(t, \pi) = 0, \\ x(0, \xi) = x_0(\xi), & \xi \in (0, \pi) \\ dy(t) = b(y(t)) dt + \sigma(y(t)) \rho(\gamma(t)) dt + \sigma(y(t)) dB_t, & t \geq 0, \\ y(0) = x \in \mathbb{R}. \end{cases}$$

473 where W is the space-time white noise on $[0, +\infty) \times [0, \pi]$ and B is a Brownian motion. An admissible
474 control γ is a predictable process $\gamma : \Omega \times [0, +\infty) \rightarrow \mathbb{R}$. The cost functional is

$$475 \quad (7.5) \quad J(x_0, \gamma) = \liminf_{T \rightarrow +\infty} \frac{1}{T} \mathbb{E} \int_0^T \int_0^\pi \ell(x(t, \xi), \gamma(t)) d\xi dt.$$

476 We assume that

- 477 1. $b : \mathbb{R} \rightarrow \mathbb{R}$ is a measurable function such that

$$478 \quad |b(y) - b(y')| \leq L_b |y - y'|,$$

480 for a suitable positive constant L_b , for every $y, y' \in \mathbb{R}$.

- 481 2. $\sigma : \mathbb{R} \rightarrow \mathbb{R}$ is a measurable and bounded function, such that

$$482 \quad |\sigma(y) - \sigma(y')| \leq L_\sigma |y - y'|,$$

484 for suitable positive constants L_σ and there exists a suitable positive δ such that:

$$485 \quad |\sigma(y)| \geq \delta > 0,$$

486 for every $y \in \mathbb{R}$.

- 487 3. there exists $\mu > 0$ such that for all $y, y' \in \mathbb{R}$:

$$488 \quad (7.6) \quad 2(b(y) - b(y'), y - y') + |\sigma(y) - \sigma(y')|^2 \leq -\mu |y - y'|^2,$$

489 4. $d : [0, \pi] \rightarrow \mathbb{R}$, $\rho : \mathbb{R} \rightarrow \mathbb{R}$ are bounded and measurable functions.

490 5. $\ell : \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ is a measurable and bounded function such that

$$491 \quad |\ell(x, \gamma) - \ell(x', \gamma)| \leq L |x - x'|,$$

492 for a suitable positive constant L , for every $x, x', \gamma \in \mathbb{R}$.

493 Under these hypotheses, see [18], the above equation can be reformulated in an infinite dimen-
494 sional space as:

$$495 \quad (7.7) \quad \begin{cases} d_t \mathcal{X}_t = \Delta \mathcal{X}_t dt - \Delta \mathbf{r}y(t)dt + \tilde{D}d\tilde{W}_t, & t \geq 0, \xi \in [0, \pi], \\ \mathcal{X}_0 = x_0(\cdot), & \xi \in (0, \pi) \\ dy(t) = b(y(t))dt + \sigma(y(t))\rho(u(t))dt + \sigma(y(t))dB(t), & t \geq 0, \\ y(0) = y_0 \in \mathbb{R}. \end{cases}$$

where $\mathcal{X}_t := x(\cdot)$ is in $L^2(0, \pi)$, \tilde{W} is a cylindrical Wiener process in $L^2(0, \pi)$, \tilde{D} is the bounded operator in $L^2(0, \pi)$ corresponding to multiplication by a bounded function d , Δ is the realisation of the Laplace operator with Dirichlet boundary conditions in $L^2(0, \pi)$, that is (denoting by $\mathcal{D}(\Delta)$ the domain of the operator)

$$\mathcal{D}(\Delta) = H^2(0, \pi) \cap H_0^1(0, \pi), \quad \Delta f = \frac{\partial^2 f}{\partial \xi^2}, \quad \forall f \in \mathcal{D}(\Delta)$$

496 Finally $\mathbf{r}(\xi) = 1 - \frac{\xi}{\pi}$, $\xi \in [0, \pi]$ is the solution to

$$497 \quad (7.8) \quad \begin{cases} \frac{\partial^2 \mathbf{r}}{\partial \xi^2}(\xi) = 0, & \xi \in (0, \pi), \\ \mathbf{r}(0) = 1, \quad \mathbf{r}(\pi) = 0. \end{cases}$$

498 It is well known that Δ generates an analytic semigroup of contractions (of negative type -1)
499 moreover, for any $\delta > 0$, $\mathbf{r} \in \mathcal{D}((-\Delta)^{1/2-\delta})$ (where $(-\Delta)^\alpha$ denotes the fractional power). Standard
500 results on analytic semigroups then yield:

$$501 \quad (7.9) \quad |(-\Delta)e^{t\Delta}\mathbf{r}|_{L^2(0,\pi)} \leq c_{\mathbf{r}}e^{-t}t^{-(\frac{1}{2}+\delta)}, \quad t > 0.$$

502 We are now in a position to rephrase the problem according to our general framework. Indeed
503 setting $H = L^2(0, \pi) \times \mathbb{R}$, $\Xi = \mathbb{R}$ and $X_t = (\mathcal{X}_t, y(t))$ equation (7.7) becomes

$$504 \quad (7.10) \quad \begin{cases} dX_t^x = AX_t^x dt + F(X_t^x)dt + QG(X_t^x)\rho(\gamma_t)dt + QG(X_t^x)dW_t^1 + DdW_t^2, & t \geq 0, \\ X_0^x = x. \end{cases}$$

505 where:

$$506 \quad 1. \quad A = \begin{pmatrix} -\Delta & -\Delta R \\ 0 & 0 \end{pmatrix} \text{ where } R : \mathbb{R} \rightarrow D((-\Delta)^{\frac{1}{2}-\delta}), \text{ is defined as } Ry = \mathbf{r}(\cdot)y, y \in \mathbb{R}$$

507 It is easy to verify that A generates a C_0 -semigroup in H .

$$508 \quad 2. \quad F : H \rightarrow H, \text{ is defined as: } F \begin{pmatrix} \mathcal{X} \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ b(y) \end{pmatrix},$$

$$509 \quad Q : \Xi \rightarrow H \text{ is defined as: } Qy = \begin{pmatrix} 0 \\ y \end{pmatrix},$$

$$510 \quad G : \Xi \rightarrow \Xi, \text{ is defined as: } G(y) = \sigma(y)$$

$$511 \quad D : H \rightarrow H \text{ is defined as: } D \begin{pmatrix} \mathcal{X} \\ y \end{pmatrix} = \begin{pmatrix} \tilde{D}\mathcal{X} \\ 0 \end{pmatrix}.$$

512 3. $W^1(t) = B(t)$ and (W^2) is a cylindrical Wiener process in H .

513 Hypotheses **(A.1 – A.5)** are immediately verified, we have to check **(A.6)**. We come back to
514 the formulation (7.7) and start with the second component y (that only depends on y_0). By (7.6),
515 Proposition 3.2 gives:

$$516 \quad (7.11) \quad \mathbb{E}|y^{y_0}(t) - y^{y'_0}(t)|^2 \leq e^{-2\mu t}|y_0 - y'_0|^2.$$

517

518 Coming now to the first component we have that it fullfills in $L^2(0, \pi)$ the following mild formulation:

$$519 \quad \mathcal{X}_t^{x_0, y_0} = e^{t\Delta} x_0 - \int_0^t \left[\Delta e^{(t-s)\Delta} \mathbf{r} \right] y^{y_0}(s) ds + \int_0^t e^{(t-s)\Delta} D dW_s.$$

520

521 Thus considering two different initial data

$$522 \quad \mathcal{X}_t^{x_0, y_0} - \mathcal{X}_t^{x'_0, y'_0} = e^{t\Delta} (x_0 - x'_0) - \int_0^t \Delta e^{(t-s)\Delta} (\mathbf{r} y^{y_0}(s) - \mathbf{r} y^{y'_0}(s)) ds.$$

523

524 By (7.9) and (7.11) choosing $\mu_0 \in (0, 1 \wedge \mu)$

$$525 \quad \mathbb{E} |\mathcal{X}_t^{x_0, y_0} - \mathcal{X}_t^{x'_0, y'_0}| \leq e^{-t} |x_0 - x'_0| + \int_0^t e^{-(t-s)} (t-s)^{-\left(\frac{1}{2} + \delta\right)} e^{-\mu s} |y_0 - y'_0| ds$$

$$526 \quad \leq e^{-t} |x_0 - x'_0| + e^{-\mu_0 t} \left[\int_0^t e^{-(1-\mu_0)(t-s)} (t-s)^{-\left(\frac{1}{2} + \delta\right)} ds \right] |y_0 - y'_0|.$$

527

528 That implies that (3.5) holds. In the same way one gets the proof of (3.4).

529 We notice that it is not at all obvious that the stronger versions (3.7), (3.8) holds in this case.

530 As far as the control functional is concerned it is enough to set $L(X, \gamma) = \int_0^\pi \ell(\xi, \mathcal{X}(\xi), \gamma) d\xi$
 531 and to verify in a straightforward way that (E.1) holds (in this case $R_1 = \rho$, $R_2 = 0$, $\Gamma = \mathbb{R}$).

532 Thus all the hypotheses of Theorem 7.1 hold and points (i) and (ii) in its thesis give the optimal
 533 ergodic cost and strategy in terms of the solution to the ergodic BSDE in (7.2). Moreover by point
 534 (iii) of Theorem 7.1 we have that (\bar{v}, λ) is the mild solution of the equation (5.1), in the sense
 535 of definition (5.2) and the optimal feedback law can be characterized in terms of the generalized
 536 directional gradient of \bar{v} .

537 **EXAMPLE 7.2.** We consider an ergodic control problem for a stochastic heat equation with
 538 Dirichlet boundary conditions with nonlinearity controlled through a one dimensional process y .

$$539 \quad (7.12) \quad \begin{cases} d_t x(t, \xi) = \frac{\partial}{\partial \xi^2} x(t, \xi) dt + f(x(t, \xi), y(t)) + d(\xi) \dot{W}(t, \xi) dt, & t \geq 0, \xi \in (0, 1), \\ x(t, 0) = x(t, 1) = 0, \\ x(0, \xi) = x_0(\xi), & \xi \in (0, 1) \\ dy(t) = b(y(t)) dt + \sigma(y(t)) \gamma(t) dt + \sigma(y(t)) dB_t, & t \geq 0, \\ y(0) = y_0 \in [-1, 1]. \end{cases}$$

540 where \mathcal{W} is the space-time white noise on $[0, +\infty) \times [0, 1]$ and B is a Brownian motion. An admissible
 541 control γ is a predictable process $\gamma : \Omega \times [0, +\infty) \rightarrow [-1, 1]$. The cost functional is

$$542 \quad (7.13) \quad J(x_0, \gamma) = \liminf_{T \rightarrow +\infty} \frac{1}{T} \mathbb{E} \int_0^T \left[\int_0^1 (\ell(x(t, \xi), y(t)) d\xi + \gamma^2(t)) dt.$$

543 We assume:

544 1. $f : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a Lipschitz map. We fix two constants $L_f > 0$ and $\mu_f \in \mathbb{R}$ such that

$$545 \quad |f(x, y) - f(x', y)| \leq L_f (|x - x'| + |y - y'|), \quad \langle f(x, y) - f(x, y'), x - x' \rangle \leq -\mu_f |x - x'|^2,$$

546

547 for every $x, x', y, y' \in \mathbb{R}$.

548 2. $b : \mathbb{R} \rightarrow \mathbb{R}$ is Lipschitz. We fix a constant $\mu_b \in \mathbb{R}$ such that:

$$549 \quad \langle b(y) - b(y'), y - y' \rangle \leq -\mu_b |y - y'|^2, \quad \forall y, y' \in \mathbb{R}$$

551 3. $\sigma : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a Lipschitz and bounded. We fix L_σ such that

$$552 \quad |\sigma(y) - \sigma(y')| \leq L_\sigma |y - y'|, \quad \forall y, y' \in \mathbb{R},$$

554 We also assume that there exists a suitable positive δ such that:

$$555 \quad |\sigma(y)| \geq \delta > 0, \quad \forall y \in \mathbb{R}.$$

556 4. $d : [0, 1] \rightarrow \mathbb{R}$ is a bounded and measurable function.

557 5. $\ell : \mathbb{R}^2 \rightarrow \mathbb{R}$ is bounded and Lipschitz

As in the previous example the above equation can be reformulated in an infinite dimensional space as:

$$\begin{cases} d_t \mathcal{X}_t = \Delta \mathcal{X}_t dt + f(\mathcal{X}_t, y(t)) dt + \tilde{D} d\tilde{W}_t, & t \geq 0, \xi \in [0, 1], \\ \mathcal{X}_0 = x_0(\cdot), & \xi \in [0, 1] \\ dy(t) = b(y(t)) dt + \sigma(y(t)) \gamma(t) dt + \sigma(y(t)) dB(t), & t \geq 0, \\ y(0) = y_0 \in \mathbb{R}. \end{cases}$$

558 where $\mathcal{X}_t := x(\cdot)$ is in $L^2(0, 1)$, \tilde{W} is a cylindrical Wiener process in $L^2(0, 1)$, Δ is the realisation of
559 the Laplace operator with Dirichlet boundary conditions in $L^2(0, 1)$, \tilde{D} is the bounded operator in
560 $L^2(0, 1)$ corresponding to multiplication by a bounded function d .

561 Finally setting $H = L^2(0, 1) \times \mathbb{R}$, $\Xi = \mathbb{R}$, $\Gamma = [-1, 1]$ and $X_t = (\mathcal{X}_t, y(t))$ equation (7.4) becomes

$$562 \quad (7.14) \quad \begin{cases} dX_t^x = AX_t^x dt + F(X_t^x) dt + QG(X_t^x) \gamma_t dt + QG(X_t^x) dW_t^1 + DdW_t^2, & t \geq 0, \\ X_0^x = x. \end{cases}$$

and the cost takes our general form:

$$J(x_0, \gamma) = \liminf_{T \rightarrow +\infty} \frac{1}{T} \mathbb{E} \int_0^T L(X(t), \gamma(t)) dt.$$

563 where

564 1. $A = \begin{pmatrix} -\Delta & 0 \\ 0 & 0 \end{pmatrix}$ generates a C_0 -semigroup in H . We also have that

$$565 \quad \langle AX, X \rangle_H = \langle \Delta \mathcal{X}, \mathcal{X} \rangle_{L^2(0,1)} \leq -\mu_\Delta |\mathcal{X}|_{L^2(0,1)}^2,$$

566 for some $\mu_\Delta > 0$.

567 2. $F : H \rightarrow H$, is defined as: $F \begin{pmatrix} \mathcal{X} \\ y \end{pmatrix} = \begin{pmatrix} f(\mathcal{X}, y) \\ b(y) \end{pmatrix}$,

568 $Q : \Xi \rightarrow H$ is defined as: $Qy = \begin{pmatrix} 0 \\ y \end{pmatrix}$,

569 $G : \Xi \rightarrow \Xi$, is defined as: $G(y) = \sigma(y)$,

570 $D : H \rightarrow H$ is defined as: $D \begin{pmatrix} \mathcal{X} \\ y \end{pmatrix} = \begin{pmatrix} \tilde{D}\mathcal{X} \\ 0 \end{pmatrix}$.

571 3. $W^1(t) = B(t)$ and (W^2) is a cylindrical Wiener process in H .

572 4. $L : H \times \Gamma \rightarrow \mathbb{R}$, $L(X, \gamma) = \int_0^1 \ell(\mathcal{X}(\xi), y) d\xi + |\gamma|^2$.

573 We also notice that in this case the Hamiltonian defined as in (7.1) becomes:

$$574 \quad (7.15) \quad \psi \left(\begin{pmatrix} \mathcal{X} \\ y \end{pmatrix}, z \right) = -\frac{z^2}{4} I_{[-2,2]}(z) + (1 - |z|) I_{[-2,2]^c}(z) + \int_0^1 \ell(\mathcal{X}(\xi), y) d\xi.$$

575 We also assume that there exists $\bar{\mu} > 0$ such that

$$576 \quad (7.16) \quad \begin{pmatrix} -\mu_\Delta - \mu_f & \frac{1}{2} L_f \\ \frac{1}{2} L_f & -\mu_b + \frac{1}{2} L_\sigma \end{pmatrix} \leq -\bar{\mu} I_{\mathbb{R}^2}.$$

577 Hypotheses (A.1 – A.5) are immediately verified. Moreover relation (7.16) ensures that (A.7)
 578 holds as well. Finally (E.1) is straight forward (in this case $R_1 = id$, $R_2 = 0$). Thus the hypotheses
 579 of Theorem 7.1 hold and points (i), (ii) and (iii) in its thesis give the optimal ergodic cost, the
 580 strategy in terms of the solution to the ergodic BSDE in (7.2) and we have that (\bar{v}, λ) is the mild
 581 solution of the equation (5.1), in the sense of definition (5.2) and the optimal feedback law can be
 582 characterized in terms of the generalized directional gradient of \bar{v} .

583 We finally wish to apply the differentiability result in Theorem 6.1 to this specific example. We
 584 notice that by (7.15) the Hamiltonian ψ is concave and differentiable with respect to z with $\nabla_z \psi \leq 1$.
 585 Thus (B.1) holds and we can choose $L_z = 1$ in (4.2). If we assume that f , b , σ and ℓ are of class
 586 C^1 in all their variables then (C.1) and (C.2) hold, moreover if we impose that $\bar{\mu} > 2\delta^{-2}$ (here,
 587 comparing with Theorem 6.1, $L_u = 0$, $M_{G^{-1}} = \delta^{-1}$) then all the assumptions of Theorem 6.1 are
 588 verified and we can conclude that function \bar{v} in Theorem 7.1 is differentiable. Consequently point
 589 (iv) in Theorem 7.1 as well applies here and we obtain that (\bar{v}, λ) is a mild solution of equation (5.1),
 590 in the sense of definition (5.3), and that the optimal feedback law can be characterized in terms of
 591 the gradient of \bar{v} .

592 Appendix A. Proof of Theorem 6.1.

593 We will need to use some results from [23, Theorem 5.21 and Section 5.6]. The first concerns
 594 finite horizon BSDEs and the estimate of their solution, while the second concerns the infinite horizon
 595 case. We restate them in our setting as follows:

596 LEMMA A.1. *Let us consider the following equation:*

$$597 \quad (A.1) \quad -dY_t = (\phi(t, Z_t, U_t) dt - \alpha Y_t) dt - Z_t dW_t^1 - U_t dW_t^2, \quad Y_T = \eta, \quad t \in [0, T], \quad \alpha \geq 0.$$

598 assume that:

- 600 1. $|\phi(t, z, u) - \phi(t, z', u')| \leq \ell(t)(|z - z'|^2 + |u - u'|^2)^{1/2}$, $\forall z, z' \in \Xi^*$, $u, u' \in H^*$, \mathbb{P} -a.s. for
 601 some $\ell \in L^2([0, T])$;
- 602 2. for $\nu_t := \int_0^t \ell^2(s) ds$, one has

$$603 \quad (A.2) \quad \mathbb{E} \left(e^{2\nu_T - 2\alpha T} |\eta|^2 \right) < \infty, \quad \mathbb{E} \left(\int_0^T e^{\nu_s - \alpha s} |\phi(s, 0, 0)| ds \right)^2 < \infty.$$

604 Then there exists a unique solution $(Y, Z, U) \in L^2_{\mathcal{P}}(\Omega; C([0, T]; \mathbb{R})) \times L^2_{\mathcal{P}}(\Omega \times [0, T]; \Xi^*) \times L^2_{\mathcal{P}}(\Omega \times$
 605 $[0, T]; H^*)$ and it verifies for all $0 \leq t \leq T$:

$$606 \quad \mathbb{E}^{\mathcal{F}_t} \left(\sup_{s \in [t, T]} e^{2(\nu_s - \alpha s)} |Y_s|^2 \right) + \mathbb{E}^{\mathcal{F}_t} \left(\int_t^T e^{2(\nu_s - \alpha s)} |Z_s|^2 ds \right) + \mathbb{E}^{\mathcal{F}_t} \left(\int_t^T e^{2(\nu_s - \alpha s)} |U_s|^2 ds \right) \leq$$

$$(A.3) \quad \mathbb{E}^{\mathcal{F}_t} (e^{2\nu T - 2\alpha T} |\eta|^2) + \mathbb{E}^{\mathcal{F}_t} \left(\int_t^T e^{V_s - \alpha s} |\phi(s, 0, 0)| ds \right)^2, \quad \mathbb{P} - a.s., \quad t \in [0, T].$$

LEMMA A.2. *Let us consider the following equation for $\alpha \geq 0$:*

$$(A.4) \quad -dY_t = (\phi(t, Z_t, U_t) dt - \alpha Y_t) dt - Z_t dW_t^1 - U_t dW_t^2, \quad t \geq 0, .$$

Assume that:

1. $|\phi(t, z, u) - \phi(t, z', u')| \leq \ell(t)(|z - z'|^2 + |u - u'|^2)^{1/2}, \forall z, z' \in \Xi^*, u, u' \in H^*, \mathbb{P} - a.s.$ for some $\ell \in L_{loc}^2([0, +\infty[);$
2. for $\nu_t := \int_0^t \ell^2(s) ds$, one has

$$(A.5) \quad \mathbb{E} \left(\int_0^\infty e^{\nu_s} |\phi(s, 0, 0)| ds \right)^2 < \infty.$$

Then there exists a unique triple of processes (Y, Z, U) with $Y \in L_P^{2,loc}(\Omega; C([0, +\infty[; \mathbb{R}))$, $Z \in L_P^{2,loc}(\Omega \times [0, +\infty[; \Xi^*)$, $U \in L_P^{2,loc}(\Omega \times [0, +\infty[; H^*)$, such that

$$(A.6) \quad \mathbb{E} \left(\sup_{t \in [0, T]} e^{2\nu_t} |Y_t|^2 \right) < +\infty, \quad \forall T \geq 0, \quad \lim_{T \rightarrow \infty} \mathbb{E}(e^{2\nu_T} |Y_T|^2) = 0.$$

Moreover

$$(A.7) \quad \mathbb{E}^{\mathcal{F}_t} \left(\sup_{s \geq t} e^{2\nu_s} |Y_s|^2 \right) + \mathbb{E}^{\mathcal{F}_t} \left(\int_t^\infty e^{2\nu_s} (|Z_s|^2 + |U_s|^2) ds \right) \leq C \mathbb{E}^{\mathcal{F}_t} \left(\int_t^\infty e^{\nu_s} |\phi(s, 0, 0)| ds \right)^2,$$

for some positive constant C .

Proof of Theorem 6.1. The proof is split into two parts. The first deals with approximating functions v^α defined in (4.10)

Part I - Differentiability of v^α

We first have to come back to the elliptic approximations:

$$(A.8) \quad Y_t^{x,\alpha} = Y_T^{x,\alpha} + \int_t^T [\psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) - \alpha Y_s^{x,\alpha}] ds - \int_t^T Z_s^{x,\alpha} dW_s^1 - \int_t^T U_s^{x,\alpha} dW_s^2,$$

and for those equations we prove that:

PROPOSITION A.1. *Under the same assumptions of Theorem 6.1 we have that, for each $\alpha > 0$, the map $x \rightarrow Y_0^{x,\alpha}$ belongs to $\mathcal{G}^1(H, \mathbb{R})$.*

Proof. We fix $n \in \mathbb{N}$ and introduce the following finite horizon approximations where $0 \leq t \leq n$:

$$Y_t^{x,\alpha,n} = \int_t^n [\psi(X_s^x, Z_s^{x,\alpha,n} G^{-1}(X_s^x), U_s^{x,\alpha,n}) - \alpha Y_s^{x,\alpha,n}] ds - \int_t^n Z_s^{x,\alpha,n} dW_s^1 - \int_t^n U_s^{x,\alpha,n} dW_s^2.$$

For such equations [16, Prop. 3.2] holds true, moreover we have from [10, Propositions 5.6 and 5.7] that $x \rightarrow Y_0^{x,\alpha,n} := v^{\alpha,n}(x)$ belongs to $\mathcal{G}^1(H, \mathbb{R})$ and $Z_t^{x,\alpha,n} = \nabla_x v^{\alpha,n}(X_t^x) G(X_t^x)$ and $U_t^{x,\alpha,n} = \nabla_x v^{\alpha,n}(X_t^x) D$.

Hence, arguing as in Proposition 4.2, we deduce that $|Z_t^{x,\alpha,n}| \leq |\nabla_x v^{\alpha,n}(X_t^x) G(X_t^x)| \leq C/\mu$ and $|U_t^{x,\alpha,n}| \leq |\nabla_x v^{\alpha,n}(X_t^x) D| \leq \frac{C}{\mu}$, with C independent of n and α .

638 Moreover, see [10, Prop 5.2], the map $x \rightarrow (Y_t^{x,\alpha,n}, Z_t^{x,\alpha,n}, U_t^{x,\alpha,n})$ is Gateaux differentiable and the
 639 equation for the derivative in the direction $h \in H$, $|h| = 1$, is the following:

$$640 \quad \nabla_x Y_t^{x,\alpha,n} h = \int_t^n [\phi^{h,\alpha}(s, \nabla_x Z_s^{x,\alpha,n} h, \nabla_x U_s^{x,\alpha,n} h) - \alpha \nabla_x Y_s^{x,\alpha,n} h] ds - \int_t^n \nabla_x Z_s^{x,\alpha,n} h dW_s^1$$

$$641 \quad - \int_t^n \nabla_x U_s^{x,\alpha,n} h dW_s^2, \quad 0 \leq t \leq n.$$

643 where

$$644 \quad \phi^{h,\alpha,n}(s, z, u) = \nabla_x \psi(X_s^x, Z_s^{x,\alpha,n} G^{-1}(X_s^x), U_s^{x,\alpha,n}) \nabla_x X_s^x h + \nabla_u \psi(X_s^x, Z_s^{x,\alpha,n} G^{-1}(X_s^x), U_s^{x,\alpha,n}) u h$$

$$645 \quad + \nabla_z \psi(X_s^x, Z_s^{x,\alpha,n} G^{-1}(X_s^x), U_s^{x,\alpha,n}) [Z_s^{x,\alpha,n} \nabla_x G^{-1}(X_s^x) \nabla_x X_s^x h + z h G^{-1}(X_s^x)].$$

647 Notice that $\phi^{h,\alpha}(t, z, u)$ is affine in z and u and :

$$648 \quad |\phi^{h,\alpha,n}(s, z, u) - \phi^{h,\alpha,n}(s, 0, 0)| \leq L_u |u| + L_z M_{G^{-1}} |z| \leq (L_z^2 M_{G^{-1}}^2 + L_u^2)^{1/2} (|z|^2 + |u|^2)^{1/2}, \quad \mathbb{P} - a.s.$$

649 where here and in the following the constant C may change from line to line but always independently
 650 from n , ε and from α .

651 We can apply Lemma A.1 with $\nu_s = (L_z^2 M_{G^{-1}}^2 + L_u^2)s =: Ks$, indeed for $\varepsilon = \frac{1}{2}(\mu - 2K)$, we
 652 have, recalling also that $U_s^{x,\alpha,n}$ and $Z_s^{x,\alpha,n}$ are bounded uniformly in s , α and n

$$653 \quad (\text{A.9}) \quad \mathbb{E} \left[\int_0^n |\phi^{h,\alpha,n}(s, 0, 0)| e^{(-\alpha+K)s} dt \right]^2 \leq \frac{C}{\varepsilon} \int_0^n e^{(\varepsilon-2\alpha+2K)s} \mathbb{E} |\nabla_x X_s^x h|^2 dt \leq \frac{C}{\mu - 2K}.$$

654 Therefore the following estimate holds, arguing as before in (A.9), for all $0 \leq t \leq n$:

$$655 \quad (\text{A.10}) \quad \mathbb{E} \sup_{s \in [t, n]} e^{2(-\alpha+K)s} |\nabla_x Y_s^{x,\alpha,n} h|^2 + \mathbb{E} \int_t^n e^{2(-\alpha+K)s} [|\nabla_x Z_s^{x,\alpha,n} h|^2 + |\nabla_x U_s^{x,\alpha,n} h|^2] dt$$

$$657 \quad \leq C \mathbb{E} \left[\int_t^n e^{(-\alpha+K)s} |\phi^{h,\alpha,n}(s, 0, 0)| ds \right]^2 \leq \frac{C e^{(-2\alpha - \frac{1}{2}\mu + K)t}}{\mu - 2K}, \quad t \leq s \leq n.$$

659 In particular, we have for all $t \geq 0$:

$$660 \quad (\text{A.11}) \quad \mathbb{E} \left(e^{2Kt} |\nabla_x Y_t^{x,\alpha,n} h|^2 \right) \leq C e^{(-\frac{1}{2}\mu + K)t}.$$

661 From estimate (A.10) we deduce that $(\nabla_x Y^{x,\alpha,n} h, \nabla_x Z^{x,\alpha,n} h, \nabla_x U^{x,\alpha,n} h)$ weakly converges in the
 662 Hilbert space $L^2(\Omega \times (0, T); \mathbb{R} \times \Xi^* \times H^*)$ to some $(R^{x,\alpha,h}, V^{x,\alpha,h}, M^{x,\alpha,h})$, for every $T > 0$. From
 663 (A.11) we also have that $\nabla_x Y_0^{x,\alpha,n} h$ converge in \mathbb{R} to $\xi^{x,\alpha,h}$.

We define for every $t \geq 0$

$$\tilde{R}_t^{x,\alpha,h} = \xi^{x,\alpha,h} + \int_0^t [\phi^{h,\alpha}(s, V_s^{x,\alpha,h}, M_s^{x,\alpha,h}) - \alpha R_s^{x,\alpha,h}] ds - \int_0^t V_s^{x,\alpha,h} dW_s^1 - \int_0^t M_s^{x,\alpha,h} dW_s^2.$$

664 Compare the above with the forward equation fulfilled by $(\nabla_x Y^{x,\alpha,n} h, \nabla_x Z^{x,\alpha,n} h, \nabla_x U^{x,\alpha,n} h)$,
 665 namely:

$$666 \quad \nabla_x Y_t^{x,\alpha,n} h = \nabla_x Y_0^{x,\alpha,n} h + \int_0^t [\phi^{h,\alpha,n}(s, \nabla_x Z_s^{x,\alpha,n}, \nabla_x U_s^{x,\alpha,n}) - \alpha \nabla_x Y_s^{x,\alpha,n} h] ds$$

667
668

$$-\int_0^t \nabla_x Z_s^{x,\alpha,n} h dW_s^1 - \int_0^t \nabla_x U_s^{x,\alpha,n} h dW_s^2, \quad \mathbb{P} - a.s..$$

Since every term in the R.H.S., passing to a subsequence if necessary, weakly converges in $L^2(\Omega \times (0, T); \mathbb{R})$, see also [16, Theo. 3.1], we have that $\tilde{R}_t^{x,\alpha,h} = R_t^{x,\alpha,h}$, \mathbb{P} -a.s. for a.e. $t \geq 0$. Thus the triplet processes $(\tilde{R}^{x,\alpha,h}, V^{x,\alpha,h}, M^{x,\alpha,h})$ verifies for all $t > 0$, \mathbb{P} -a.s.:

$$\tilde{R}_t^{x,\alpha,h} = \tilde{R}_0^{x,\alpha,h} + \int_0^t [\phi^{h,\alpha}(s, V_s^{x,\alpha,h}, M_s^{x,\alpha,h}) - \alpha \tilde{R}_s^{x,\alpha,h}] ds - \int_0^t V_s^{x,\alpha,h} dW_s^1 - \int_0^t M_s^{x,\alpha,h} dW_s^2.$$

669 where

$$\begin{aligned} \phi^{h,\alpha}(s, z, u) &= \nabla_x \psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) \nabla_x X_s^x h + \nabla_u \psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) u h \\ &+ \nabla_z \psi(X_s^x, Z_s^{x,\alpha} G^{-1}(X_s^x), U_s^{x,\alpha}) [Z_s^{x,\alpha} \nabla_x G^{-1}(X_s^x) \nabla_x X_s^x h + z h G^{-1}(X_s^x)]. \end{aligned}$$

673 Moreover, thanks to (A.10) and (A.11) we have that

$$(A.12) \quad \mathbb{E} \sup_{s \in [0, T]} e^{2Ks} |\tilde{R}_s^{x,\alpha,h}|^2 < +\infty \quad \text{and} \quad \mathbb{E} e^{2Ks} |\tilde{R}_s^{x,\alpha,h}|^2 \leq \tilde{C} e^{(-\mu+2K)s},$$

675 therefore, $(\tilde{R}^{x,\alpha,h}, V^{x,\alpha,h}, M^{x,\alpha,h})$ is the unique solution of equation:

$$(A.13) \quad d_s R_s = [\phi^{h,\alpha}(s, V_s, M_s) - \alpha R_s] ds - V_s dW_s^1 - M_s dW_s^2,$$

677 in the class of processes with the regularity imposed in Lemma A.2 verifying:

$$(A.14) \quad \mathbb{E} \sup_{t \in [0, T]} |\tilde{R}_t^{x,\alpha,h}|^2 < +\infty \quad \text{and} \quad \lim_{T \rightarrow +\infty} \mathbb{E} e^{2K2T} |\tilde{R}_T^{x,\alpha,h}|^2 = 0, \quad \forall T > 0.$$

679 We then closely follow the proof of [16, Prop 3.2], indeed we get that $\lim_{n \rightarrow +\infty} \nabla_x Y_0^{\alpha,n,x} h =$
680 $\tilde{R}^{\alpha,x,h}(0)$, defines a linear and bounded operator $\tilde{R}^{\alpha,x}(0)$ from H to H , by (A.11), such that
681 $\tilde{R}^{\alpha,x}(0)h = \tilde{R}^{x,\alpha,h}(0)$, moreover for every fixed $h \in H$, $x \rightarrow \tilde{R}^{\alpha,x}(0)h$ is continuous in x , we
682 will sketch the argument by the the end of the proof in a similar point. Therefore, by dominated
683 convergence, we get that:

684

$$(A.15) \quad \lim_{\ell \downarrow 0} \frac{Y_0^{x+\ell h, \alpha} - Y_0^{x, \alpha}}{\ell} = \lim_{\ell \downarrow 0} \lim_{n \rightarrow \infty} \frac{Y_0^{x+\ell h, \alpha, n} - Y_0^{x, \alpha, n}}{\ell} = \lim_{\ell \downarrow 0} \lim_{n \rightarrow \infty} \int_0^1 \nabla_x Y_0^{x+\theta \ell h, \alpha, n} h d\theta$$

$$= \lim_{\ell \downarrow 0} \int_0^1 \tilde{R}^{x+\theta \ell h, \alpha}(0) h d\theta = \tilde{R}^{x, \alpha}(0) h.$$

686
687

688 Thus v^α is differentiable and since $Y_t^{x,\alpha} = v^\alpha(X_t^x)$ we have $\nabla_x Y_t^{x,\alpha} h = v^\alpha(X_t^x) \nabla_x X_t^x h$.

689 Fixing $T > 0$ we can see the equation satisfied by $(Y^{x,\alpha}, Z^{x,\alpha}, U^{x,\alpha})$ as a BSDE on $[0, T]$ with
690 final condition $v^\alpha(X_T^x)$ and we can apply standard results on the differentiability of markovian, finite
691 horizon BSDEs (see, for instance, [10]) to deduce that the map $x \rightarrow Y^{x,\alpha}$ is of class \mathcal{G}^1 from H to
692 $L_P^2(\Omega; C([0, T]; \mathbb{R}))$ and $x \rightarrow Z^{x,\alpha}$ is of class \mathcal{G}^1 from $L_P^2([0, T] \times \Omega; \Xi^*)$. Moreover for every $h \in H$,
693 for every $0 \leq t \leq T$ it holds that:

$$(A.16) \quad \begin{aligned} \nabla_x Y_t^{x,\alpha} h &= \nabla_x Y_T^{x,\alpha} h + \int_t^T [\phi^h(s, \nabla_x Z_s^{x,\alpha} h, \nabla_x U_s^{x,\alpha} h) - \alpha \nabla_x Y_s^{x,\alpha} h] ds \\ &- \int_t^T \nabla_x Z_s^{x,\alpha} h dW_s^1 - \int_t^T \nabla_x U_s^{x,\alpha} h dW_s^2, \quad 0 \leq t \leq n. \end{aligned}$$

Comparing the above with (A.13) and noticing that for all $T > 0$:

$$\mathbb{E}e^{2KT}|\nabla_x Y_T^{x,\alpha}h|^2 = \mathbb{E}e^{2KT}|\nabla_x v^\alpha(X_T^x)\nabla_x X_T^x h|^2 \leq Ce^{(2K-\mu)T},$$

the uniqueness part of Lemma A.2 tells us that $(\nabla_x Y^{x,\alpha}h, \nabla_x Z^{x,\alpha}h, \nabla_x U^{x,\alpha}h)$ coincides with $(\tilde{R}^{x,h,\alpha}, V^{x,h,\alpha}, M^{x,h,\alpha})$ and is the unique solution of equation (A.13) in the sense of Lemma A.2.

Part II - Differentiability of \bar{v}

We also introduce the following infinite horizon BSDE:

$$(A.17) \quad -dR_s^{x,h} = \phi^h(s, V_s^{x,h}, M_s^{x,h})ds - V_t^{x,h}dW_t^1 - M_t^{x,h}dW_t^2 \quad t \geq 0.$$

with

$$\begin{aligned} \phi^h(s, z, u) = & [\nabla_x \psi(X_s^x, \bar{Z}_s^x G^{-1}(X_s^x), \bar{U}_s^x) + \nabla_z \psi(X_s^x, \bar{Z}_s^x G^{-1}(X_s^x), \bar{U}_s^x) \bar{Z}_s^x \nabla_x G^{-1}(X_s^x)] \nabla_x X_s^x h \\ & + \nabla_u \psi(X_s^x, \bar{Z}_s^x G^{-1}(X_s^x), \bar{U}_s^x) u + \nabla_z \psi(X_s^x, \bar{Z}_s^x G^{-1}(X_s^x), \bar{U}_s^x) z. \end{aligned}$$

By Lemma A.2 has a unique solution in the class of processes $R^{x,h} \in L_{\mathcal{P}}^{2,loc}(\Omega; C([0, +\infty[; \mathbb{R}))$, $V^{x,h} \in L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; \Xi^*)$, $M \in L_{\mathcal{P}}^{2,loc}(\Omega \times [0, +\infty[; H^*)$ verifying:

$$(A.18) \quad \lim_{T \rightarrow +\infty} e^{2KT} \mathbb{E}|R_T^{x,h}|^2 = 0, \quad \forall T > 0.$$

As in [19, Theorem 5.1] we claim that, along the sequence (α_m) introduced in (4.21), it holds:

$$(A.19) \quad \nabla_x v^{\alpha_m}(x)h = \nabla_x Y_0^{\alpha_m, x}h = R_0^{x, \alpha_m, h} \rightarrow R_0^{x, h},$$

as $m \rightarrow \infty$.

Let us introduce again some parabolic approximations. For $s \in [0, n]$ consider:

$$\begin{cases} -dR_s^{x, \alpha, n, h} = \phi^{h, \alpha}(s, V_s^{x, \alpha, n, h}, M_s^{x, \alpha, n, h})ds - \alpha R_s^{x, \alpha, n, h} ds - V_s^{x, \alpha, n, h} dW_s^1 - M_s^{x, \alpha, n, h} dW_s^2, \\ R_n^{x, \alpha, n, h} = 0. \end{cases}$$

and

$$\begin{cases} -dR_s^{x, n, h} = \phi^h(s, V_s^{x, h, n}, M_s^{x, n, h})ds - V_s^{x, h, n} dW_s^1 - M_s^{x, h} dW_s^2, \\ R_n^{x, h, n} = 0, \end{cases}$$

Since along the sequence (α_m) selected in Section 4 we have

$$\mathbb{E} \sup_{s \in [0, n]} |\bar{Y}_s^x - Y_s^{x, \alpha_m}|^2 + \mathbb{E} \int_0^n [|\bar{Z}_s - Z_s^{x, \alpha_m}|^2 + |\bar{U}_s^x - U_s^{x, \alpha_m}|^2] ds \rightarrow 0.$$

and consequently

$$\mathbb{E} \int_0^n |\phi^{h, \alpha_m}(s, 0, 0) - \phi^h(s, 0, 0)|^2 ds \rightarrow 0 \quad \text{as } m \rightarrow \infty.$$

standard estimates on finite horizon BSDEs give:

$$(A.20) \quad \mathbb{E} \sup_{s \in [0, n]} |R_s^{x, n, h} - R_s^{x, \alpha_m, n, h}|^2 \rightarrow 0, \quad \text{as } m \rightarrow \infty.$$

718 Moreover if we compare with the solution $(\tilde{R}^{x,\alpha,h}, V^{x,\alpha,h}, M^{x,\alpha,h})$ of equation (A.13)

(A.21)

$$\begin{cases} -d(R_s^{x,\alpha,n,h} - \tilde{R}_s^{x,\alpha,h}) = [\phi^{h,\alpha}(s, V_s^{x,\alpha,n,h} - V_s^{x,\alpha,h}, M_s^{x,\alpha,n,h} - M_s^{x,\alpha,h}) - \alpha(R_s^{x,\alpha,n,h} - \tilde{R}_s^{x,\alpha,h})] ds \\ -[V_s^{x,\alpha,n,h} - V_s^{x,\alpha,h}] dW_s^1 - [M_s^{x,\alpha,n,h} - M_s^{x,\alpha,h}] dW_s^2, \\ R_n^{x,\alpha,n,h} - \tilde{R}_n^{x,\alpha,h} = -\nabla_x v^\alpha(X_n^x) \nabla_x X_n^x h \end{cases}$$

719

720 Thus Lemma A.1 estimate (A.3) yields:

721 (A.22) $|R_0^{x,\alpha,n,h} - \tilde{R}_0^{x,\alpha,h}|^2 \leq \mathbb{E}(e^{2Kn} |\nabla_x v^\alpha(X_n^x) \nabla_x X_n^x h|^2) \leq Ce^{(2K-\mu)n} \rightarrow 0, \text{ as } n \rightarrow +\infty.$

722 Notice that the right hand side does not depend on α . Finally

723 (A.23)
$$\begin{cases} -d(R_s^{x,n,h} - R_s^{x,h}) = \phi^h(s, V_s^{x,n,h} - V_s^{x,h}, M_s^{x,n,h} - M_s^{x,h}) ds \\ -[V_s^{x,n,h} - V_s^{x,h}] dW_s^1 - [M_s^{x,n,h} - M_s^{x,h}] dW_s^2, \\ R_n^{x,n,h} - R_n^{x,h} = -\tilde{R}_n^{x,h}, \end{cases}$$

724 and taking into account (A.18), one has, again by Lemma A.1 relation (A.3):

725 (A.24) $|R_0^{x,n,h} - R_0^{x,h}|^2 \leq \mathbb{E}(e^{2Kn} |R_n^{x,h}|^2) \leq Ce^{(2K-\mu)n} \rightarrow 0, \text{ as } N \rightarrow +\infty.$

726 Therefore summing up (A.22), (A.24) and (A.20) we have that:

727
$$R_0^{x,\alpha_m,h} \rightarrow R_0^{x,h}, \text{ as } m \rightarrow +\infty.$$

728 Finally the continuity with respect to x of $R_0^{x,h}$ descends immediately from (A.24) and from the
729 continuity of the map $x \rightarrow R_0^{x,n,h}$ proved in [10, Prop. 4.3].

730 We can now conclude as above (and as in [16, Prop 3.2]); $R^{x,h}(0)$, defines a linear and bounded
731 operator $R^x(0)$ from H to H , such that $R^x(0)h = R^{x,h}(0)$, and we have:

732
$$\lim_{t \downarrow 0} \frac{\bar{v}(x+th) - \bar{v}(x)}{t} = \lim_{t \downarrow 0} \frac{\bar{Y}_0^{x+th} - \bar{Y}_0^x}{t} = \lim_{t \downarrow 0} \lim_{m \rightarrow 0} \frac{Y_0^{x+th,\alpha_m} - Y_0^{x,\alpha}}{t} =$$

733
$$= \lim_{t \downarrow 0} \lim_{m \rightarrow 0} \int_0^1 \nabla_x Y_0^{x+\theta th,\alpha_m} h d\theta = \lim_{t \downarrow 0} \lim_{m \rightarrow 0} \int_0^1 R^{x+\theta th,\alpha_m,h}(0) h d\theta =$$

734
$$= \lim_{t \downarrow 0} \int_0^1 R^{x+\theta th}(0) h d\theta = R^x(0)h.$$

735

736 □

737 **Acknowledgements.** The authors would like to thank both referees for their careful reading
738 of our paper and for their valuable insights.

739 REFERENCES

740 [1] J.P. Aubin. *Applied functional analysis*. John Wiley & Sons, New York-Chichester-Brisbane, 1979. Translated
741 from the French by Carole Labrousse, With exercises by Bernard Cornet and Jean-Michel Lasry.

742 [2] S.N. Cohen and Y. Hu. Ergodic bsdes driven by markov chains. *SIAM Journal on Control and Optimization*,
743 51(5):4138–4168, 2013. cited By 6.

744 [3] A. Cosso, M. Fuhrman, and H. Pham. Long time asymptotics for fully nonlinear Bellman equations: a backward
745 SDE approach. *Stochastic Process. Appl.*, 126(7):1932–1973, 2016.

- 746 [4] A. Cosso, G. Guatteri, and G. Tessitore. Ergodic control of infinite-dimensional stochastic differential equations
747 with degenerate noise. *ESAIM Control Optim. Calc. Var.*, 25:Art. 12, 29, 2019.
- 748 [5] G. da Prato and J. Zabczyk. *Ergodicity for infinite-dimensional systems*, volume 229 of *London Mathematical*
749 *Society Lecture Note Series*. Cambridge University Press, Cambridge, 1996.
- 750 [6] R. W. R. Darling and E. Pardoux. Backwards SDE with random terminal time and applications to semilinear
751 elliptic PDE. *Ann. Probab.*, 25(3):1135–1159, 1997.
- 752 [7] A. Debussche, Y. Hu, and G. Tessitore. Ergodic BSDEs under weak dissipative assumptions. *Stochastic Process.*
753 *Appl.*, 121(3):407–426, 2011.
- 754 [8] M. Fuhrman. A class of stochastic optimal control problems in Hilbert spaces: BSDEs and optimal control laws,
755 state constraints, conditioned processes. *Stochastic Process. Appl.*, 108(2):263–298, 2003.
- 756 [9] M. Fuhrman and G. Tessitore. The Bismut-Elworthy formula for backward SDEs and applications to nonlinear
757 Kolmogorov equations and control in infinite dimensional spaces. *Stoch. Stoch. Rep.*, 74(1-2):429–464, 2002.
- 758 [10] M. Fuhrman and G. Tessitore. Nonlinear Kolmogorov equations in infinite dimensional spaces: the backward
759 stochastic differential equations approach and applications to optimal control. *Ann. Probab.*, 30(3):1397–
760 1465, 2002.
- 761 [11] M. Fuhrman and G. Tessitore. Generalized directional gradients, backward stochastic differential equations and
762 mild solutions of semilinear parabolic equations. *Appl. Math. Optim.*, 51(3):279–332, 2005.
- 763 [12] B. Goldys and B. Maslowsky. Ergodic control of semilinear stochastic equations and the Hamilton-Jacobi
764 equation. *J. Math. Anal. Appl.*, 234(2):592–631, 1999.
- 765 [13] M. Hu and F. Wang. Ergodic bsdes driven by g-brownian motion and applications. *Stochastics and Dynamics*,
766 18(6), 2018. cited By 0.
- 767 [14] Y. Hu and F. Lemonnier. Ergodic BSDE with an unbounded and multiplicative underlying diffusion and appli-
768 cation to large time behavior of viscosity solution of HJB equation. *arXiv e-prints*, page arXiv:1801.01284,
769 Jan 2018.
- 770 [15] Y. Hu, P.Y. Madec, and A. Richou. A probabilistic approach to large time behavior of mild solutions of HJB
771 equations in infinite dimension. *SIAM J. Control Optim.*, 53(1):378–398, 2015.
- 772 [16] Y. Hu and G. Tessitore. BSDE on an infinite horizon and elliptic PDEs in infinite dimension. *NoDEA Nonlinear*
773 *Differential Equations Appl.*, 14(5-6):825–846, 2007.
- 774 [17] I. Kharroubi and H. Pham. Feynman-Kac representation for Hamilton-Jacobi-Bellman IPDE. *Ann. Probab.*,
775 43(4):1823–1865, 2015.
- 776 [18] I. Lasiecka and R. Triggiani. *Differential and algebraic Riccati equations with application to boundary/point*
777 *control problems: continuous theory and approximation theory*, volume 164 of *Lecture Notes in Control and*
778 *Information Sciences*. Springer-Verlag, Berlin, 1991.
- 779 [19] Y. Hu M. Fuhrman and G. Tessitore. Ergodic BSDEs and optimal ergodic control in Banach spaces. *SIAM J.*
780 *Control Optim.*, 48(3):1542–1566, 2009.
- 781 [20] Y. Hu M. Fuhrman and G. Tessitore. Stochastic maximum principle for optimal control of SPDEs. *C. R. Math.*
782 *Acad. Sci. Paris*, 350(13-14):683–688, 2012.
- 783 [21] P.Y. Madec. Ergodic bsdes and related pdes with neumann boundary conditions under weak dissipative as-
784 sumptions. *Stochastic Processes and their Applications*, 125(5):1821–1860, 2015. cited By 3.
- 785 [22] E. J. McShane and R. B. Warfield, Jr. On Filippov’s implicit functions lemma. *Proc. Amer. Math. Soc.*,
786 18:41–47, 1967.
- 787 [23] E. Pardoux and A. Răşcanu. *Stochastic differential equations, backward SDEs, partial differential equations*,
788 volume 69 of *Stochastic Modelling and Applied Probability*. Springer, Cham, 2014.
- 789 [24] A. Richou. Ergodic bsdes and related pdes with neumann boundary conditions. *Stochastic Processes and their*
790 *Applications*, 119(9):2945–2969, 2009. cited By 17.
- 791 [25] M. Royer. BSDEs with a random terminal time driven by a monotone generator and their links with PDEs.
792 *Stoch. Stoch. Rep.*, 76(4):281–307, 2004.