

A Comparison Study of Modelling Techniques for Permanent Magnet Machines

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Abstract—In this paper, four different modelling techniques for permanent magnet (PM) machines are compared for their accuracy and computational complexity. The considered techniques are primarily based on conformal mapping and harmonic modelling. In conformal mapping, the slotted air gap is mapped into a simpler canonical shape, where the field solution is calculated and then mapped back to the original domain. In harmonic modelling, the regions of the machine cross section are represented as Fourier series and coupled with each other by means of boundary conditions. The field solution is obtained by solving the boundary value problem. In order to quantify the accuracy of the field solutions, global parameters such as cogging torque and flux linkage are computed. The effectiveness of the modelling techniques are evaluated by comparing the global parameters and the simulation time with finite element analysis (FEA) results.

Index Terms—Harmonic Modelling, Relative Permeance, Complex relative Permeance, Conformal Mapping, Schwarz-Christoffel Toolbox, FEA

NOMENCLATURE

B_r	Magnet remanence (T)
B_{slr}	Slotless air-gap radial flux density (Wb/m ²)
$B_{sl\theta}$	Slotless air-gap tangential flux density
B_s	Slotted air-gap flux density (Wb/m ²)
$B_{s\theta}$	slotted air gap tangential flux density
J_z	Current density distribution z-direction (A/m ²)
H	Magnetic field strength (A/m)
A_z	Magnetic vector potential z-direction (Vs/m)
$\tilde{\lambda}$	Relative permeance
λ	Complex relative permeance
μ_0	Permeability of air (Vs/(Am))
μ_r	Relative permeability of iron
p	Number of pole pairs
α_p	Magnet arc/pole pitch ratio
Q_s	Slot number
l_m	Magnet radial thickness (m)
g	Air gap length (m)
R_g	Radius of air gap center (m)
R_r	Radius of the rotor surface (m)
R_m	Radius of the magnet surface (m)
R_s	Stator outer radius (m)

l_a	stack length (m)
b_0	Slot opening (m)
d_s	Slot depth (m)
N_c	No. of conductors in a slot
T_c	Cogging Torque (Nm)
r, θ	Coordinates of air-gap evaluation points

I. INTRODUCTION

Permanent Magnet (PM) machines are widely used for their high efficiency and power density. The motor requirements on volume, axial length, outer diameter, efficiency, weight and cost are different [1], [2] for each of these applications. Optimal selection of the motor parameters is essential to meet these requirements effectively. In order to use a motor model in an optimization routine, it has to be accurate and computationally cheap. There are several modelling techniques available in literature, which have their own merits and demerits. In this study, four such models are analyzed and the results are compared with FEM.

In the mathematical models, for simplicity, the end windings are neglected, the iron is assumed to have infinite permeability, and the magnet end effects are not considered. With these assumptions, the electro-magnetic problem can be solved for the magnetic vector or scalar potential in two dimension and the Laplacian or Poissonian equations are solved for the field solutions [3]–[6]. The field solutions can further be used to calculate torque, back-emf, and losses.

The modelling techniques considered in this study are relative permeance model (RP) [7], complex permeance model (CP) [8], Schwarz-Christoffel Toolbox model (SC) [9], and Harmonic model(HM) [5]. The first three models are fundamentally based on conformal mapping. For the SC toolbox model, MATLAB toolbox, which allows numerical mapping of complex polygons to simple rectangles, can be utilized. The harmonic model allows one to use the exact solutions of Laplacian or Poissonian equations that are represented as Fourier series in the tangential direction [10], [11].

II. MODELLING TECHNIQUES

A twenty pole permanent magnet machine is considered in the analysis as a test bench, shown in Fig.1, and its parameters are given in Table I.

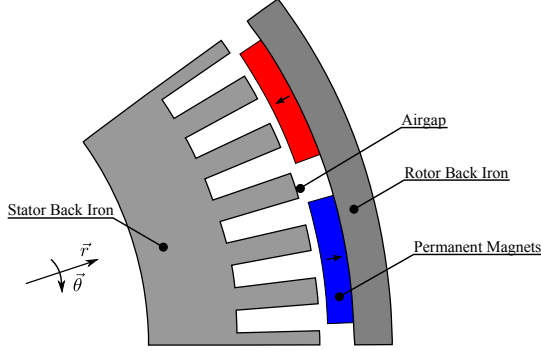


Fig. 1. Test bench machine for the comparison of modelling techniques

TABLE I
PARAMETERS OF EXTERNAL ROTOR SURFACE PM MOTOR

Parameters,	Symbol,	Value,	Unit
Number of poles	$2p$	20	-
Slot number	Q_s	60	-
Magnet arc/pole pitch ratio	α_p	0.75	-
Air gap length	g	1.2	mm
Radius of the rotor surface	R_r	85	mm
Radius of the magnet surface	R_m	76.2	mm
Stator outer radius	R_s	75	mm
Magnet remanence	B_r	1.19	T
Core length	l_a	60	mm
Slot depth	d_s	12	mm
No. of conductors in a slot	N_c	6	-

A. Relative permeance model

The radial and tangential components of the slotless air gap field solution are obtained in the polar coordinates using (1) to (4). In order to capture the slotting effect, the radial field solution is multiplied with relative permeance. The influence of slotting on tangential component of the air gap field is neglected.

The slotless air gap field solution given in [12]:

When $(np \neq 1)$

$$B_{slr} = \sum_{n=1,3,\dots}^{\infty} \frac{-4npB_r \sin \frac{n\pi\alpha_p}{2}}{n\pi\mu_r((np)^2 - 1)} \left[\left(\frac{r}{R_m}\right)^{np-1} + \left(\frac{R_s}{R_m}\right)^{np-1} \left(\frac{R_s}{r}\right)^{np+1} \right]$$

$$\left\{ \frac{(np-1) \left(\frac{R_m}{R_r}\right)^{2np} + 2 \left(\frac{R_m}{R_r}\right)^{np-1} - (np+1)}{\frac{\mu_r+1}{\mu_r} \left[1 - \left(\frac{R_s}{R_r}\right)^{2np} \right] - \frac{\mu_r-1}{\mu_r} \left[\left(\frac{R_s}{R_m}\right)^{2np} - \left(\frac{R_m}{R_r}\right)^{2np} \right]} \right\} \cos(np\theta) \quad (1)$$

$$B_{sl\theta} = \sum_{n=1,3,\dots}^{\infty} \frac{-4npB_r \sin \frac{n\pi\alpha_p}{2}}{n\pi\mu_r((np)^2 - 1)} \left[\left(\frac{-r}{R_m}\right)^{np-1} + \left(\frac{R_s}{R_m}\right)^{np-1} \left(\frac{R_s}{r}\right)^{np+1} \right]$$

$$\left\{ \frac{(np-1) \left(\frac{R_m}{R_r}\right)^{2np} + 2 \left(\frac{R_m}{R_r}\right)^{np-1} - (np+1)}{\frac{\mu_r+1}{\mu_r} \left[1 - \left(\frac{R_s}{R_r}\right)^{2np} \right] - \frac{\mu_r-1}{\mu_r} \left[\left(\frac{R_s}{R_m}\right)^{2np} - \left(\frac{R_m}{R_r}\right)^{2np} \right]} \right\} \sin(np\theta) \quad (2)$$

and when $(np = 1)$

$$B_{slr} = \sum_{n=1,3,\dots}^{\infty} \frac{2npB_r \sin \frac{n\pi\alpha_p}{2}}{n\pi\mu_r((np)^2 - 1)} \left[1 + \left(\frac{R_s}{r}\right)^2 \right] \left\{ \frac{np \left(\frac{R_m}{R_s}\right)^2 - np \left(\frac{R_r}{R_s}\right)^2 + \left(\frac{R_r}{R_s}\right)^2 \ln \left(\frac{R_m}{R_r}\right)^2}{\frac{\mu_r+1}{\mu_r} \left[1 - \left(\frac{R_r}{R_s}\right)^2 \right] - \frac{\mu_r-1}{\mu_r} \left[\left(\frac{R_m}{R_r}\right)^2 - \left(\frac{R_r}{R_m}\right)^2 \right]} \right\} \cos(np\theta) \quad (3)$$

$$B_{sl\theta} = \sum_{n=1,3,\dots}^{\infty} \frac{2npB_r \sin \frac{n\pi\alpha_p}{2}}{n\pi\mu_r((np)^2 - 1)} \left[-1 + \left(\frac{R_s}{r}\right)^2 \right] \left\{ \frac{np \left(\frac{R_m}{R_s}\right)^2 - np \left(\frac{R_r}{R_s}\right)^2 + \left(\frac{R_r}{R_s}\right)^2 \ln \left(\frac{R_m}{R_r}\right)^2}{\frac{\mu_r+1}{\mu_r} \left[1 - \left(\frac{R_r}{R_s}\right)^2 \right] - \frac{\mu_r-1}{\mu_r} \left[\left(\frac{R_m}{R_r}\right)^2 - \left(\frac{R_r}{R_m}\right)^2 \right]} \right\} \sin(np\theta) \quad (4)$$

The relative permeance ($\tilde{\lambda}$) is derived as a function of flux path length and modification factor (γ) [7]. The radial variation of the field solution can be captured by varying the modification factor (γ) that changes the effective flux path length as shown in Fig.2.

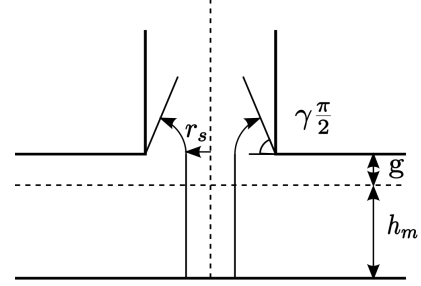


Fig. 2. Relative permeance model

So the relative permeance can be expressed as,

$$\tilde{\lambda} = \begin{cases} \frac{g+l_m/\mu_R}{g+l_m/\mu_R+\gamma\frac{\pi}{2}\left[\frac{b_0}{2}-r_s\right]}, & \text{for region of the slot opening} \\ 1, & \text{for region of the tooth} \end{cases} \quad (5)$$

where,

$$\gamma = \frac{g+l_m/\mu_R}{\pi b_0/4} \left(\frac{B_{max}}{B_{min}} - 1 \right)$$

The cogging torque (T_c) is calculated by summing the lateral forces acting on the teeth as in [7]:

$$T_c = \sum_{k=1}^{Q_s} l_a \int_0^{\frac{b_0}{2}} \left(\frac{B_1^2 - B_2^2}{2\mu_0} \right) r_t dy \quad (6)$$

where b_0 is slot opening, r_t is equal to $R_s - r_s$, and B_1 and B_2 are the flux densities along opposite sides of the slot walls.

B. Complex relative permeance model

The relative permeance model fairly well estimates the radial component of the field solution in the slotted air gap, but it does not include the tangential component which can be useful to derive the closed form solution for the cogging torque and electromagnetic torque based on integration of

the Maxwell's stress tensor. The real and imaginary parts of the complex relative permeance are derived to take into account the influence of slotting on both radial and tangential components of the air gap flux density. The field distribution in the slotted air gap (B_s) is obtained by multiplying the slotless air gap flux density (B_s) and the complex conjugate of the relative permeance (λ^*).

$$B_s = B_k \lambda^* = B_k \left(\frac{\partial k}{\partial s} \right)^* = B_k \left(\frac{\partial k}{\partial t} \frac{\partial t}{\partial w} \frac{\partial w}{\partial z} \frac{\partial z}{\partial s} \right)^* \quad (7)$$

Each partial derivative in (7) is defined by conformal transformation explained in [8] and the final expression is

$$B_s = B_k \left[\frac{k}{s} \frac{(w-1)}{(w-a)^{\frac{1}{2}}(w-b)^{\frac{1}{2}}} \right] \quad (8)$$

so the complex relative permeance is written as

$$\lambda = \lambda_a + j\lambda_b = \frac{k}{s} \frac{(w-1)}{(w-a)^{\frac{1}{2}}(w-b)^{\frac{1}{2}}} \quad (9)$$

The radial and tangential components of the slotted air gap field solution are derived as

$$B_s = B_{sr} + jB_{s\theta} = B_k \lambda^* = (B_r \lambda_a + B_\theta \lambda_b) + j(B_\theta \lambda_a + B_r \lambda_b). \quad (10)$$

In order to employ equation for the slotless air gap field solution [12], the mapped evaluation points are assumed to form a circular arc and the magnets are considered to retain their shapes. Although these two assumptions allow one to derive the closed form solution for flux density and cogging torque, they will impair the accuracy [8].

C. Schwarz-Christoffel Toolbox

The field solution can be calculated with better accuracy by capturing the distortions of the magnet shape and evaluation points in the slotless domain. For this purpose, the numerical conformal mapping using SC Toolbox [9] is employed and Hague's [3] equation is solved.

The canonical domain field solution is given as [3]

$$\Omega = \frac{\mu_0 I}{2\pi} \left(\arctan \left[\frac{\tan \left(\frac{\pi}{2\Delta y} (y + I_y) \right)}{\tanh \left(\frac{\pi}{2\Delta y} (x - \Delta x k - I_x) \right)} \right] + \arctan \left[\frac{\tan \left(\frac{\pi}{2\Delta y} (y - I_y) \right)}{\tanh \left(\frac{\pi}{2\Delta y} (x - \Delta x k - I_x) \right)} \right] \right) \quad (11)$$

and

$$B_{wr} = -\mu_0 \frac{\partial \Omega}{\partial x} \quad (12)$$

$$B_{w\theta} = -\mu_0 \frac{\partial \Omega}{\partial y} \quad (13)$$

The solution (B_w) can be mapped back to the original slotted air gap using the permeance function

$$B_s = B_w \left(\frac{\partial w}{\partial s} \right)^* = B_w \left(\frac{\partial w}{\partial z} \frac{\partial z}{\partial s} \right)^* \quad (14)$$

where

$$\frac{\partial w}{\partial z} = \text{evaldiff}(f, w)$$

$$\frac{\partial z}{\partial s} = \frac{1}{R_g e^{j\theta}}$$

The function *evaldiff* is a MATLAB function introduced in its workspace by SC Toolbox.

D. Harmonic Modelling

The geometry of the test bench machine, shown in Fig 1, has well defined regions like slots, air gap, and rotor permanent magnets. Due to this fact, the model parameters are described by periodic functions for each region. For a 2D problem, using the definition of magnetic vector potential and Ampere's circuital law, the following system of differential equations can be derived in the polar coordinate system:

$$\mu_r H_r = \frac{1}{r} \frac{\partial A_z}{\partial \theta}; \quad (15)$$

$$\mu_\theta H_\theta = -\frac{\partial A_z}{\partial r}; \quad (16)$$

$$J_z = \frac{1}{r} H_\theta + \frac{\partial H_\theta}{\partial r} - \frac{1}{r} \frac{\partial H_r}{\partial \theta}. \quad (17)$$

where A_z is the component of vector magnetic potential in the axial direction (z), H_r and H_θ are the radial and tangential field strength respectively, μ_r , μ_θ and J_z are the distribution of the iron permeability and current density, which is zero in this study. The equations (15-17) are solved for A_z ,

$$A_z = W r^{\lambda_m} a + W r^{-\lambda_m} b + r^2 G_1. \quad (18)$$

where λ_m and W are the results of eigen-decomposition of the root of the quadratic equation resulting from the differential equation, G_1 is the constant which contains the source components from the region, and a and b are the unknowns. For coupling multiple regions with each other, continuous boundary conditions are applied [5].

E. Finite Element Analysis

The results from the analytical models are compared with the results from a commercial FEA software (FLUX2D). The geometry of the model is finely meshed so that further refinement of the mesh will not improve the results. The assumptions considered in the FEA model are identical to the ones in the analytical models. The permeability of the iron is assigned to a large value to approximate infinite permeability. The same computation is repeated for different number of mesh nodes in order to study their influence on computational time.

III. DISCUSSION ON THE RESULTS

A. Magnetic field solution

From Fig.3 it can be observed that the radial flux density waveforms of the analytical models, except the relative permeance model, are matching well with FEA result. The tangential field, in Fig. 4, is accurately calculated by the Harmonic and SC Toolbox models. The complex permeance model has some deviations.

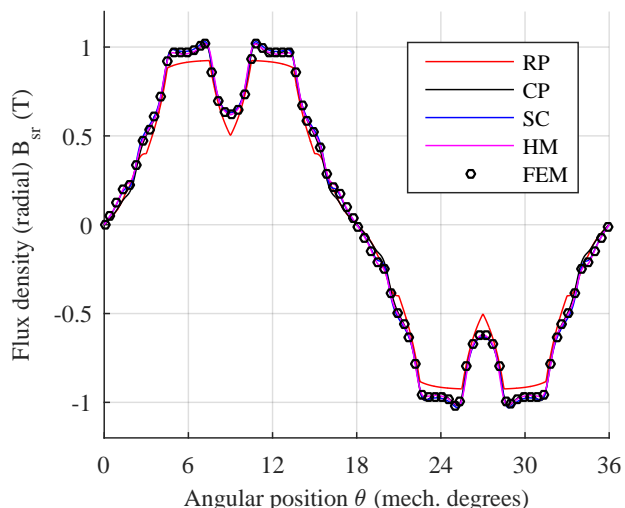


Fig. 3. Radial flux density in the air gap

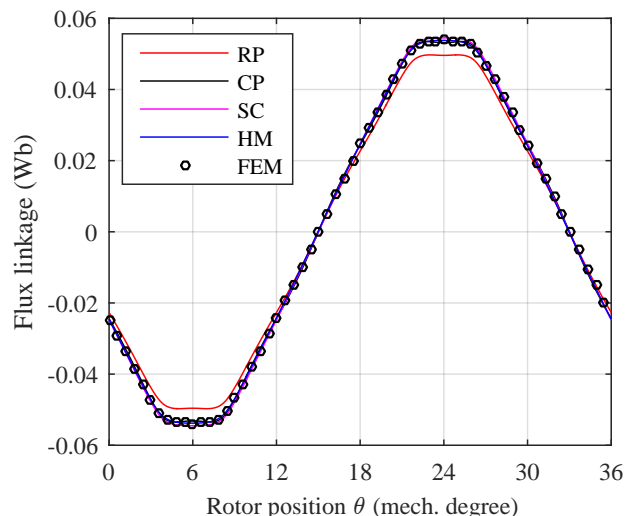


Fig. 5. Flux linkage of phase A

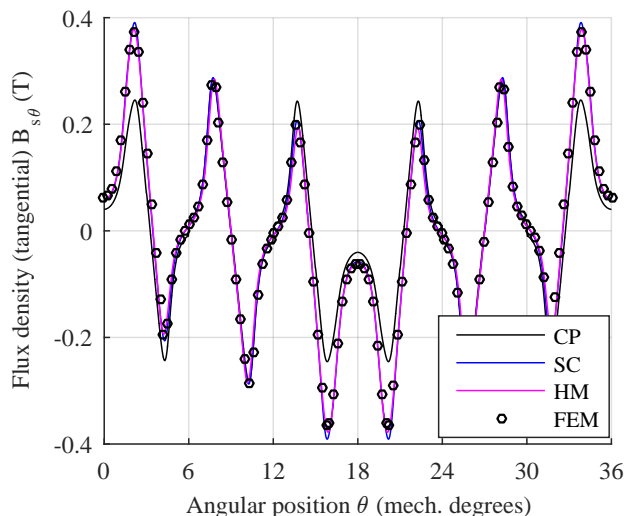


Fig. 4. Tangential flux density in the air gap

B. Flux linkage

The flux linkage of phase A winding is computed by integrating the radial flux density component across the coil pitch and multiplying it by the number of coils connected in series. The analytical models, except relative permeance model, have good correspondence with FEA model as shown in Fig.5.

C. Cogging torque

The cogging torque is more sensitive to the accuracy of the field solutions. A closed form solution based on Maxwell stress tensor is used for calculating the cogging torque in the complex permeance model. In harmonic and SC Toolbox models, tangential component of Maxwell stress tensor is numerically integrated. In the relative permeance model, the

cogging torque is calculated by integrating the lateral forces acting on the teeth walls as in (6). The FEM software uses virtual work method for computing the forces and torque. From Fig. 6 it is evident that the field solutions from harmonic and SC Toolbox models are more accurate. Although the radial field component of the complex permeance model is accurate, the deviation in tangential component produces higher cogging torque values. Due to the inaccuracy of the radial flux density solution produced by the relative permeance model, its cogging torque values are also high.

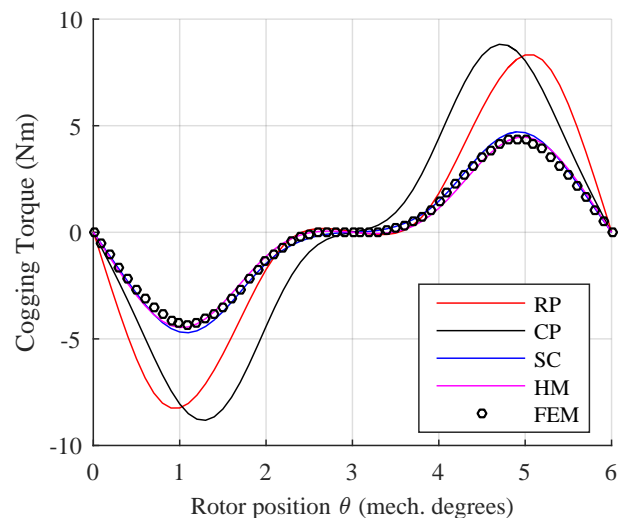


Fig. 6. Cogging torque of the given test bench machine

D. Performance analysis

The accuracy of the field solutions and calculation time are influenced by the number of harmonics and discretization points. The number of harmonics represents the maximum

order of harmonics used in the slotless air gap field solution calculated using (1) to (4) and in the Fourier equation (18) that defines the parameters in the harmonic model. The number of discretization points denotes the number of evaluation points in the middle of the air gap along an angular span of two pole pitches used for calculation of the air-gap field.

A parameter sweep analysis is performed to obtain the minimum time required to get a stable result in terms of variation of RMS value of the cogging torque waveform. The reference RMS value is the one obtained using FEA with maximum number of nodes. For example, in Fig. 7 the RMS value of the cogging torque gets stabilized after 30 harmonics and the corresponding calculation time is approx. 6 seconds. The offset between the stabilized RMS value and the reference RMS value from FEM indicates the accuracy. The fastest convergence is observed for the harmonic model and its accuracy is also high.

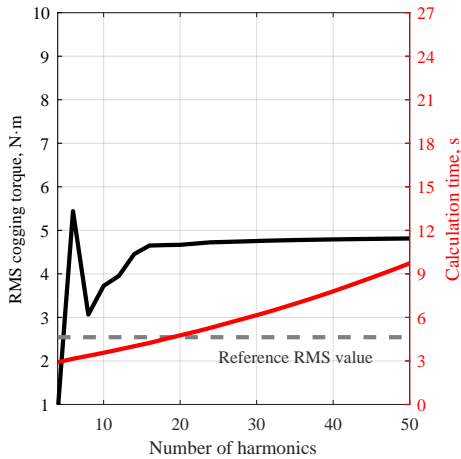


Fig. 7. Performance analysis based on number of harmonics for RP model (Discretization points = 1000).

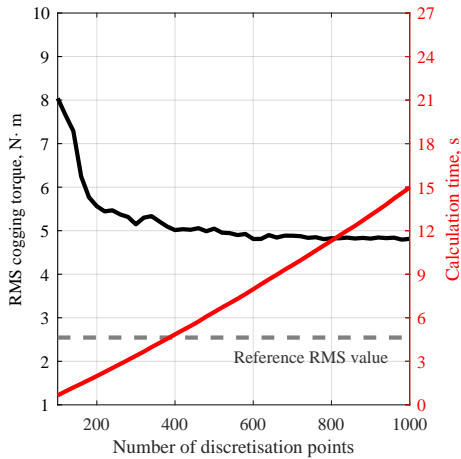


Fig. 8. Performance analysis based on number of discretization points for RP model (Harmonics = 50).

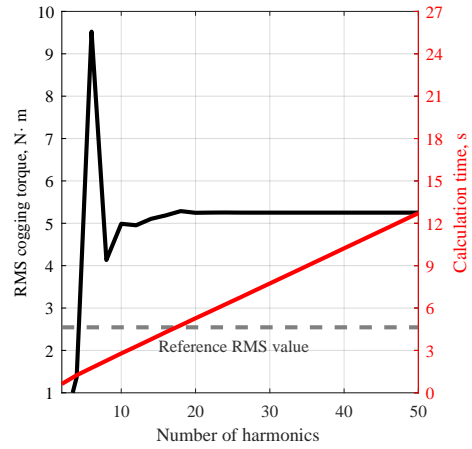


Fig. 9. Performance analysis based on number of harmonics for CP model (Discretization points = 50).

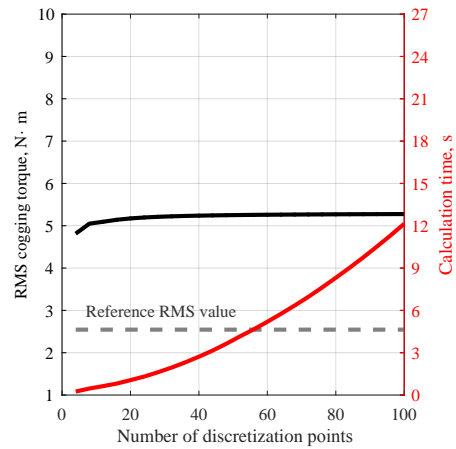


Fig. 10. Performance analysis based on number of discretization points for CP model (Harmonics = 50).

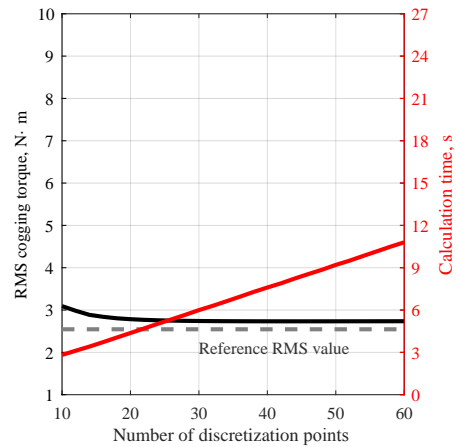


Fig. 11. Performance analysis based on number of discretization points for SC model.

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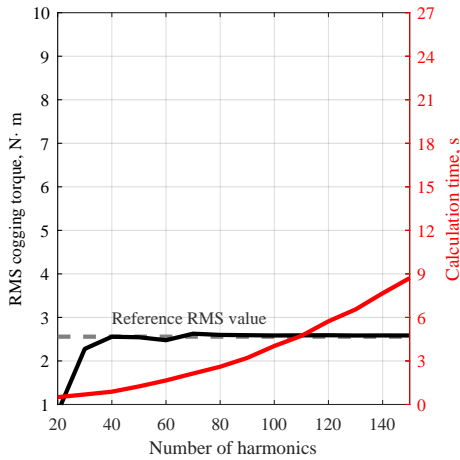


Fig. 12. Performance analysis based on number of harmonics for HM model.

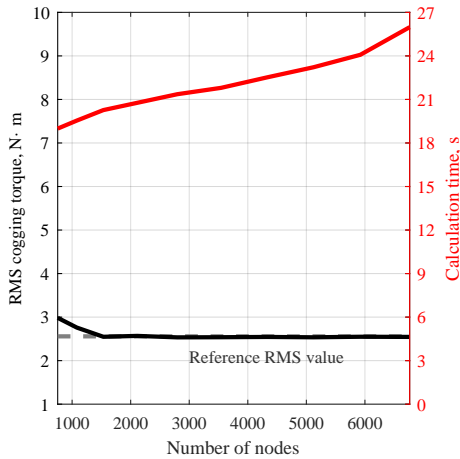


Fig. 13. Performance analysis based on number of nodes for FEM model.

IV. CONCLUSIONS

Open circuit analysis was performed on a 20 pole surface PM machine using four different modelling techniques and the results are compared with FEA model. The accuracy of the modelling techniques are studied by comparing the cogging torque and flux linkage waveforms. Similarly, the modelling complexities are evaluated based on their computation time. The parameter sweep analysis was performed to find the minimum number of harmonics or minimum number of discretization points required to obtain a stable result in terms of RMS value of the waveform. It can be observed that the most accurate result compared to FEA is obtained by the harmonic model, which is followed by the SC Toolbox based model. The complex permeance model gives the radial component of the field solution with good accuracy, but its tangential component has some deviations. The results from the relative permeance model are not accurate enough. In terms of computation time, the harmonic model's convergence is the fastest and the complex permeance model comes next.

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