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Investigating the effect of dynamic laser beam oscillations in remote fusion cutting process

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[Matteo Busatto](javascript:;) ¹: [Leonardo Caprio](javascript:;) ¹ [;](https://orcid.org/0000-0003-4964-9698) [Barbara Previtali](javascript:;) ¹

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Matteo Busatto, **D** Leonardo Caprio, **D** and Barbara Previtali **D**

AFFILIATIONS

Department of Mechanical Engineering, Politecnico di Milano, Via La Masa 1, Milan 20156, Italy

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ABSTRACT

The latest research on laser beam fusion cutting has revealed significant improvements in process productivity and cut quality through the use of dynamic beam shaping techniques. While many studies have investigated dynamic beam shaping for proximity cutting, the influence of laser beam oscillations on the remote fusion cutting process remains unexplored. The present work aims to study the effect of dynamic beam shaping on the remote fusion cutting process through analytical modeling, experimental investigations, and in situ high-speed monitoring. Initially, an analytical model based on thermodynamic analysis was developed to assess the influence of circular oscillations on the $\frac{1}{\infty}$ process zone. This model facilitates the evaluation of process performance from an energetic perspective, providing an estimate of the maximum achievable cutting speed for the remote fusion cutting process across various operating conditions. A significant increment in process productivity could be achieved through beam oscillations. Furthermore, based on theoretical findings, the effect of circular laser beam oscillations superimposed on the processing feed direction was experimentally investigated using a 1 mm thick AISI304 stainless steel $\frac{8}{5}$
Beam oscillations superimposed on the processing feed direction was exper material. A 6 kW fiber laser was utilized, alongside a high-speed camera-based system for in situ process monitoring. The experimental results demonstrate a significant increase in the process productivity under dynamic beam shaping conditions, consistent with theoretical findings. Specifically, the maximum achievable cutting speed could be increased from 0.13 to 0.20 m/s. Furthermore, the cut quality of produced samples was evaluated in terms of kerf morphology and profile. 18 November 2024 10:36:43

Key words: laser cutting, stainless steel, high speed imaging, remote cutting, beam shaping

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I. INTRODUCTION

Today, laser cutting of metal materials stands out as one of the prevalent applications in the manufacturing industry. Conventionally, the separation of metal materials relies on the combined use of a high-intensity laser beam and a high-pressure gas jet, arranged coaxially and operating in close proximity to the workpiece.^{[1,2](#page-9-0)} The laser beam melts the metal, while an inert gas jet expels the molten material from the cutting kerf. In the absence of an external gas and its blowing action, the laser may operate from a distance, resulting in a remote cutting configuration.³ Instead of utilizing high-pressure gas to expel the molten material, the remote fusion cutting process relies on a single-pass removal mechanism driven by partial vaporization and downward melt ejection induced by vapor recoil pressure.^{[4](#page-9-0)}

To date, only a few publications have addressed remote fusion cutting process in terms of productivity and part quality. Antonova et al. investigated the remote cutting of steel plates with $CO₂$ laser,^{[5](#page-9-0)} while the effect of pulsed emission was discussed by Gladush and Rodionov.^{[6](#page-9-0)} Lütke et al. compared the performance of remote fusion cutting for different material thicknesses using a fiber laser source.^{$\frac{7}{8}$ $\frac{7}{8}$ $\frac{7}{8}$ $\frac{7}{8}$ $\frac{7}{8}$} The influence of process parameters on the cutting front angle was addressed by Schober.^{[9](#page-9-0)} On the other hand, the impact of the laser beam incident angle on process performances was investigated by Villumsen and Kristiansen for the processing of stainless steel plates.^{[10](#page-9-0)} A detailed analysis on cut quality for different materials was performed by Pihlava et al. in terms of kerf width, roughness, perpendicularity, cut-edge sharp-ness, and burr height.^{[11](#page-9-0)} Furthermore, theoretical models of the

process dynamics were proposed by Matti *et al.*^{[12](#page-9-0)} and Kristiansen et al.¹

Remote cutting offers numerous advantages for processing thin materials, such as the reduced risk of tool collision, elimination of cutting gas usage, narrower access paths, and increased accessibility due to the ability to operate at a distance. However, its adoption in the industrial field remains limited. This is primarily due to the significantly lower processing speeds compared to conventional fusion and reactive cutting processes. However, recent advancements in laser cutting techniques utilizing novel beam shaping methods have demonstrated substantial improvements in processing speed and part quality. These techniques refer to the manipulation of laser beam characteristics to achieve a desired beam profile and intensity distribution.

Latest research from KU Leuven and IWS Dresden has revealed notable improvements in terms of process productivity and part quality, for the fusion cutting of high-thickness materials by employing dynamic beam shaping (DBS) techniques.^{[14](#page-9-0)} These approaches rely on the high speed motion of laser beam, superimposed to the cutting direction. Laser oscillations around the kerf enable to shape the interaction zone into an arbitrary geometry while preserving the beam intensity distribution. Goppold et al. investigated the effects of different laser oscillation patterns on the part quality of high-thick stainless steel plates by employing a highdynamic scanner system $15,16$ $15,16$ $15,16$ and a tip-tilt piezo platform mirror.¹ Studies by Levichev et al. have demonstrated significant improvements in the part quality of 12 mm thick mild and stainless steel employing longitudinal linear oscillations.^{[18,19](#page-9-0)} Thermal monitoring and high-speed imaging of the cutting process were introduced by Kardan et al ^{[20](#page-9-0),[21](#page-9-0)} for evaluating the process dynamics, focusing on the impact of oscillation amplitude and frequency on the workpiece temperature and melt flow behavior.

Although significant improvements of process performances have been demonstrated for fusion cutting of thick metal materials, to date, the application of beam shaping techniques for the remote fusion cutting process has been scarcely explored in the literature, $22,23$ while the impact of laser beam oscillations has never been addressed.

The scope of this work is to explore the effect of dynamic beam shaping on the remote fusion cutting process. Initially, an analytical model based on thermodynamic analysis was developed to evaluate process performance from an energetic perspective. This provides an estimate of the maximum achievable cutting speed for the remote fusion cutting process across various operating conditions. Furthermore, the effect of circular laser beam oscillations superimposed to the processing feed direction was experimentally investigated. A 6 kW fiber laser source was employed for the remote cutting of 1 mm thick stainless steel AISI304, which represents one of the most commonly used material and thickness in industrial remote cutting applications. Moreover, a high-speed camera-based system for in situ process monitoring was used to observe the melt dynamics.

II. MODELING

An analytical model for estimating the maximum achievable cutting speed during the remote fusion cutting process was developed based on the thermodynamic model presented by Steen and Mazumder.^{[24](#page-9-0)} According to the proposed energy balance approach, the absorbed laser power P_{abs} required to melt the volume of material necessary for generating a cutting kerf can be derived as

$$
P_{abs} = P \cdot A = F \nu \rho \left[c_p \Delta T + h_S + f_\nu h_\nu \right], \tag{1}
$$

where P is the average power delivered by the laser source, A is the absorptivity coefficient, F is the cross-sectional area of the cutting kerf, ΔT is the solidus temperature variation, ρ is the material density, c_p is the specific heat capacity, h_s and h_v are the latent heat of fusion and vaporization, respectively, while f_v is the fraction of vaporized melt.

From Eq. (1), it is evident that the absorbed laser power P_{abs} can be estimated by calculating the absorptivity coefficient A and the rate of molten material per unit time, which is related to the cross-sectional area F, for the specific material of interest.

By approximating the cutting front with an inclined plane, as schematically illustrated in Fig. $1(a)$, and disregarding multiple reflections within the cutting kerf, the absorptivity coefficient A can be expressed by the Fresnel equation. This equation determines the coefficient as a function of laser beam incident angle θ on the cutting front. Based on the Fresnel equations, the laser reflectivity R varies according to material properties, angle of incidence, and polarization state of the radiation.

The average angle of incidence at the maximum cutting speed $\frac{1}{\infty}$ [Fig. 1(a)] can be computed according to Mahrle and Beyer,^{[25](#page-9-0)} November 2024 while the contributions from parallel and perpendicular polarization states can be computed as

age angle of incidence at the maximum cutting speed
\nin be computed according to Mahrle and Beyer,²⁵
$$
\frac{1}{6}
$$

\ntributions from parallel and perpendicular polariza-
\nthe computed as
\n
$$
R_P = \frac{[n \cdot \cos(\theta) - 1]^2 + [k \cdot \cos(\theta)]^2}{[n \cdot \cos(\theta) + 1]^2 + [k \cdot \cos(\theta)]^2},
$$
\n(2)

$$
R_S = \frac{[n - \cos(\theta)]^2 + k^2}{[n + \cos(\theta)]^2 + k^2},
$$
\n(3)

FIG. 1. (a) Cutting front and angle of incidence, θ , for the maximum achievable cutting speed and (b) cross-sectional area, F, of the cutting kerf.

where $n + ik$ is the material complex refractive index. For unpolarized laser radiation, the overall reflectivity is calculated as the average value of the two contributions and the absorptivity can be expressed as

$$
A = 1 - R = 1 - \frac{R_S + R_P}{2}.
$$
 (4)

Moreover, an analytical expression for the cross-sectional area F was proposed by Lind et al.^{[26](#page-9-0)} Considering [Fig. 1\(b\)](#page-2-0), they showed that the cross-sectional area of the cutting kerf can be approximated as the area covered by the beam caustic in the plane transversal to the feeding direction,

$$
F = \int_{z=0}^{z=s} \left[d_{\text{waist}} \cdot \sqrt{1 + \frac{(z - z_0)^2}{z_R^2}} \right] dz, \tag{5}
$$

where d_{waist} is the laser beam waist diameter at z_0 axial position and z_R is the Rayleigh length.

It is worth noticing that the laser beam diameter $d(z)$ influences both the cross-sectional area of the kerf, F, and the absorptivity A, having a significant impact on the estimated absorbed laser power.

Furthermore, a fraction of the laser power absorbed by the workpiece dissipates through heat conduction phenomena, which does not contribute to the melting process. An analytical expression of these power losses was provided by Schulz et al ^{[27](#page-9-0)} as

$$
P_L = 4 \cdot K_c \cdot s \cdot \Delta T \cdot \left[\frac{Pe}{2}\right]^{0.3},\tag{6}
$$

where K_c is the thermal conductivity, s is the material thickness, and Pe is the Peclet number expressed as

TABLE II. Stainless steel AISI304: physical properties.

| Physical Properties | Values |
|-------------------------------------------------------------------------------|--------------|
| Complex refractive index, $n + ki$ | $3.6 - 5.0i$ |
| Solidus temperature, T_{solidus} (K) | 1670 |
| Material density, ρ_{liquid} (kg m ⁻³) | 6900 |
| Heat capacity, $c_{p,\text{liquid}}$ (J kg ⁻¹ K ⁻¹) | 780 |
| Thermal conductivity, K_c (W m ⁻¹ K ⁻¹) | 22. |
| Thermal diffusivity, $\alpha \cdot 10^{-6}$ (m ² s ⁻¹) | 6.06 |
| Melting latent heat, h_m (J kg ⁻¹) | 247 000 |
| Vaporized material fraction, f_{ν} (%) | 10 |
| Vaporization latent heat, h_v (J kg ⁻¹) | 6 340 000 |
| Surface tension, γ (N m ⁻¹) | 1.880 |

$$
Pe = \frac{\rho \cdot c_p}{K_c} \cdot \frac{r_0}{2} \cdot \nu,\tag{7}
$$

where c_p is the specific heat capacity, ρ is the material density, r_0 is the width of the melting front, and ν is the processing speed.

Finally, according to Lind et $al.^{26}$ $al.^{26}$ $al.^{26}$, the maximum cutting speed achievable with a laser power P is reached when the absorbed power is no longer sufficient to melt the material volume per unit of time required to generate the cutting kerf. An estimate of this cutting speed under various processing conditions can be derived as follows:

$$
v_{\text{max}} = \frac{P \cdot A - P_L}{F \cdot \rho \cdot (c_p \Delta T + h_S + f_v h_v)}.
$$
 (8)

III. MATERIALS AND METHODS

A. Material

The material employed was 1 mm thick stainless steel AISI304. Its nominal chemical composition and physical properties are summarized in Tables I and II, respectively.²

B. Laser system

The remote fusion cutting experiments were conducted on a robotic laser system originally designed for welding applications

FIG. 2. Schematic representation of the experimental setup for the remote fusion cutting process with a high-speed monitoring system.

TABLE III. Laser system specifications.

| Laser cutting system | Values | |
|---------------------------------------------------------------|--------|--|
| Wavelength, λ (nm) | 1070 | |
| Beam quality factor, M^2 | 11.2 | |
| Laser power, $P(W)$ | 6000 | |
| Feeding fibre core diameter, $d_{\text{core}} (\mu \text{m})$ | 100 | |
| Focal length, f_{foc} (mm) | 300 | |
| Collimation length, f_{col} (mm) | 200 | |
| Beam waist diameter, $d_{\text{waist}}(\mu m)$ | 150 | |

and subsequently adapted to perform remote separation operations (BLMGroup, Cucciago, Italy). Within the same laser system, both welding and remote cutting processes can be successfully performed. This system consisted of a 6-degree-of-freedom robotic arm and a 2-degree-of-freedom rotary/tilting table. An industrial high-power multimode fiber laser source capable of delivering up to 6 kW of power at 1070 nm (YLS-6000-CUT, IPG Photonics Coorp., Oxford, Massachusetts) was employed to cut the material. The transport fiber, having a core diameter of $100 \mu m$, was coupled to a wobble head capable of oscillating the beam in various trajectories with up to 500 Hz frequency and 3 mm amplitude (IPG D50, Oxford, MA, USA). A schematic representation of the experimental setup is shown in [Fig. 2](#page-3-0).

The optical path consisted of a 200 mm collimating and a 300 mm focusing lens producing a $150 \mu m$ beam waist diameter. Laser system specifications are reported in Table III.

Moreover, a high-speed camera and a secondary illumination light source were employed to observe the process from an off-axis perspective. The illumination light was a CAVILUX HF (Cavitar, Tampere, Finland) laser source emitting at 640 nm and synchronized with the shutter of the camera. The high-speed camera equipped with a CMOS sensor was a Fastcam Mini AX200 (Photron, Tokyo, Japan) enabling elevated acquisition rates up to 900 000 fps.

C. Experimental design

Initially, the conventional remote fusion cutting process employing linear trajectories was addressed. Previous literature publications highlighted the influence of the laser beam size on the ejection of the molten material from the kerf, which directly impacts the feasibility of material separation. Accordingly, the effect of the laser beam dimension and processing speed was experimentally evaluated in the first part of this study. The laser power

FIG. 3. Schematic representation of the remote fusion cutting process: (a) defocused laser beam on a linear trajectory and (b) focused laser beam oscillating with a circular pattern. (c) Schematic representation of a circular trajectory.

was kept fixed at 5 kW, while the focal position was varied between 0 and 10 mm above the workpiece surface. In these working condi- 18 November 2024 10:36:43 $\vec{8}$ tions, the spot sizes on the material surface varied from 150 up to $1028 \,\mu$ m. Furthermore, 11 cutting speed levels between 0.05 and 0.15 m/s were examined, evenly spaced at intervals of 0.01 m/s. Process parameters are summarized in Table IV.

Conventional remote fusion cutting process employs highly defocused laser beams to enlarge the kerf to facilitate melt ejection. However, as an alternative solution, a well-focused and oscillating $\tilde{\ddot{\theta}}$ laser beam can be also used. In the second part of this study, the effect of circular laser beam oscillations was experimentally investigated. A laser beam focused on the material surface was used, while circular oscillation patterns with variable amplitudes were superimposed to the linear trajectory. A comparison between the two operating conditions is represented in Figs. $3(a)$ and $3(b)$.

The oscillation amplitudes were selected to achieve spot size dimensions on the material surface equal to the one obtained with defocused laser beams [Fig. $3(c)$]. The equivalent spot diameter varied between 150 and 1028 μ m, corresponding to oscillation amplitudes between 0 and $878 \mu m$. The linear speed was varied between 0.05 and 0.25 mm/s, with intervals of 0.05 m/s. Based on a preliminary investigation, a fixed oscillation frequency value of 450 Hz was considered. Process parameters are summarized in [Table V.](#page-5-0)

D. Characterization and measurement equipment

Each processing condition was replicated two times resulting in a total of 224 samples. These were visually inspected, and three categorical conditions were identified.

• No separation: the molten material is not ejected but resolidified, forming a remolten seam.

(a) $\mu(w_{\text{terf}})$ (b) $\sigma(w_{\text{kerf}})$ 1_{mm}

FIG. 5. Kerf profile identification via image processing algorithm. (a) Average width $\mu(w_{\text{kerf}})$ and (b) standard deviation $\sigma(w_{\text{kerf}})$.

concession separation. The moneth material is ejected, reading to Representative optical microscope images of these conditions complete material separation. are reported in Fig. 4 from a top view perspective. No separation conditions are indicated in red, while unstable and successful separations are indicated in yellow and green, respectively.

> The Echo-Lab optical microscope (Echo-Lab SM 535 H, Devco S.r.l., Milan, Italy) was employed for high-resolution imaging of samples, while the kerf profile of successful conditions was analyzed through a MATLAB image processing algorithm. The top view images of the cutting kerf taken with the Mitutoyo optical microscope (Mitutoyo Quick Vision PRO ELF QV-202) were converted into binary matrices and rotated for misalignment compensation (see Fig. 5). Subsequently, the upper and lower kerf profiles were identified and the $\frac{2}{8}$ average width $\mu(w_{\text{kerf}})$ of the kerf was computed as the difference $\frac{3}{8}$ between the two [Fig. 5(a)]. Additionally, the kerf variability was deterr 2024 1 mined as the standard deviation of the width profile $\sigma(w_{\text{kerf}})$ [Fig. 5(b)].

IV. RESULTS

A. Linear cutting trajectories

Based on the previously presented analytical model (see Sec. [II](#page-2-0)), the absorptivity A and cross-sectional area F are determined as a

FIG. 6. Absorptivity A (in green) and cross-sectional area F (in blue), as a function of beam spot diameter on the material surface.

- Unstable separation: the molten material is only partially ejected, leading to partial separation of the material.
- Successful separation: the molten material is ejected, leading to

FIG. 4. Qualitative cutting results: (a) no separation (red), (b) unstable separation (yellow), and (c) successful separation (green).

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18 November 2024 10:36:43

10:36:43

FIG. 7. Feasibility map for linear trajectory: no separation (red), unstable separation (yellow), and successful separation (green). The maximum achievable cutting speed curve, derived from the analytical model, is depicted as a dashed green line.

function of laser beam dimension on material's surface as illustrated in [Fig. 6](#page-5-0). It is evident that, as the laser beam dimension increases, the absorptivity coefficient decreases, while the cross-sectional area increases.

According to Eq. [\(8\),](#page-3-0) these parameters have a significant impact on the maximum processing speed, which was calculated and is represented in Fig. 7, as a function of beam dimension. Additionally, Fig. 7 displays the feasibility map for linear cutting trajectory experiments, superimposed to the maximum achievable cutting speed curve derived from the analytical model.

A smaller beam allows for higher laser intensity and greater absorptivity, while also decreasing the cross-sectional area and, thus, the rate of molten material per unit time. As depicted in Fig. 7, the experimental and theoretical results demonstrated that the combined increase in absorptivity and reduction in the crosssectional area gives the possibility to cut faster. A value of 0.08 m/s processing speed was achieved with a spot diameter of $1082 \mu m$, whereas, with a spot of $677 \mu m$, the speed increased to 0.13 m/s. From energetic considerations, further reducing the laser beam dimensions enables an even greater increase in the processing speed. However, the experimental results have shown that process conditions with lower spot dimensions result in failure to achieve material separation.

This is related to the fact that the beam size on the material surface is a crucial characteristic to enable molten material ejection out of the kerf. Figure 8 presents a visual comparison between the three separation conditions through high-speed video frames. The successful and unstable separation conditions correspond to beam spot diameters of $667 \mu m$ at 0.10 and 0.15 m/s processing speed, respectively [Figs. $8(a) - 8(b)$], while the no separation condition refers to a spot size of 150 μ m and a cutting speed of 0.10 m/s [Fig. 8(c)].

FIG. 8. High-speed video frames for (a) successful, (b) unstable, and (c) no

Without the blowing action of the assist gas, hydrodynamic conditions of the molten material play a significant role in the melt ejection and, consequently, on the material separation. The successful expulsion of the molten material relies on the pressure field generated by the recoil pressure, as well as material viscosity and surface tension as discussed by Hirano et al.^{[30](#page-9-0)} Small laser beams generate a restricted molten area, which severely limits the ejection of the molten material. This leads to the formation of an inconsistent remolten seam, ultimately resulting in a failure to achieve separation. Conversely, enlarging the laser beam enables a wider molten area, thereby enhancing melt ejection and, thus, enabling complete material separation even in the absence of the gas jet.

separation conditions.

Furthermore, the cut quality of successful separation conditions was characterized in terms of average kerf width $\mu(w_{\text{kerf}})$ and

FIG. 9. Cut quality evaluation as a function of cutting speed and spot diameter for linear trajectories: (a) kerf width $\mu(w_{\text{kerf}})$ and (b) kerf variability $\sigma(w_{\text{kerf}})$.

variability $\sigma(w_{\text{kerf}})$. These are plotted in Fig. 9 as a function of cutting speed and spot dimension on the material surface. Both beam spot dimension and cutting speed have a notable impact on the kerf width and variability. Larger beam dimensions on the material surface result in a considerable widening of the kerf, whereas faster processing speeds lead to a significant reduction in the kerf width and a slight increase in kerf variability.

From Fig. 9, it is worth noticing that the kerf width is smaller than the laser spot dimensions on the material surface. This is because part of the molten material is not ejected from the kerf but remains attached to the sample cut-edges. This is evident from the high-speed video image in [Fig. 8](#page-6-0) and the sample cross sections, where the cut-edges exhibit a rounded shape caused by the

FIG. 10. Linear cutting trajectory: (a) sample cross section along the transversal direction and (b) microscope photo of the sample top surface.

FIG. 11. Feasibility map for circular oscillations: no separation (red), unstable separation (yellow), and successful separation (green). The maximum speed curve, derived from analytical model, is depicted as a dashed green line.

resolidified material adhering to the kerf due to surface tension. An example is reported in Fig. $10(a)$, while the corresponding micro- \vec{a} scope photo of the kerf profile is shown in Fig. 10(b). November 2024

B . Laser beam oscillations

As already mentioned, a well-focused, and oscillating laser beam can be employed as an alternative solution to widen the kerf. This offers significantly higher absorptivity compared to a defocused beam (see [Fig. 6\)](#page-5-0). By selecting oscillation amplitudes to achieve a spot dimension on the material surface comparable to that of defocused beam under linear trajectory conditions [as illustrated in Fig. $3(c)$], an estimate of the maximum achievable cutting speed for laser beam oscillations was retrieved. Although the molten volume is equal to the one obtained for linear trajectories, the higher absorptivity leads to a more efficient process, giving the possibility to cut faster. The feasibility map for the experiments under circular laser oscillations is reported in Fig. 11, together with the maximum achievable cutting speed curve derived from the analytical model.

The experimental results show an increase in the processing speed for all laser spot dimensions considered, consistent with theoretical findings. Specifically, the maximum achievable cutting speed could be increased from 0.13 to 0.20 m/s for 677 μ m laser spot diameter.

For successful conditions, the average kerf width $\mu(w_{\text{kerf}})$ and its profile variability $\sigma(w_{\text{kerf}})$ are plotted in [Fig. 12](#page-8-0) as a function of cutting speed and spot dimensions on the material surface.

It is evident that wider laser spots on the material surface led to an expansion in the kerf dimension, while kerf variability remained unaffected by spot dimensions or cutting speed. Moreover, a notable difference between the variability of the left

18 November 2024 10:36:43

FIG. 12. Cut quality evaluation as a function of cutting speed and spot diameter for circular oscillations: (a) kerf width $\mu(w_{\rm ker})$; (b) left and (c) right kerf variability $\sigma(W_{\text{kerf}})$.

and right kerf profiles can be observed in Fig. 12. This difference is attributed to the circular movement of the laser beam, which significantly impacts the quality of the cut-edge profile.

In particular, the right profile exhibits significant deterioration due to the accumulation of remolten material, as highlighted in the cross-sectional image shown in Fig. 13(a). High-speed monitoring videos of the remote cutting process under circular beam oscillation [Fig. $13(b)$], reveal that a portion of the molten material is not expelled downward, as in the static condition, but it is pushed backward and deposited on the right cut-edge profile.

While it is evident that circular oscillations can significantly increase the processing speed of the remote fusion cutting process, further tests and analysis are needed to ensure that they can also enhance part quality.

FIG. 13. Circular oscillation for a focused laser beam. Beam diameters of 677 μ m and speed of 0.1 m/s: (a) cross section; (b) high-speed video frame.

V. CONCLUSIONS

The present work explores the influence of dynamic beam shaping on the remote fusion cutting process. An analytical model based on thermodynamic analysis was developed to evaluate process performances by providing an estimate of the maximum achievable cutting speed for the remote fusion cutting process across various operating conditions.

The impact of the laser beam size on the separation process was investigated through theoretical analysis and experimentation. The results demonstrated that a smaller beam enables higher absorptivity and reduced cross-sectional areas, ultimately increasing processing speed. However, in the absence of assistant gas, the restricted molten area resulting from small laser beams severely limits the ejection of molten material. This leads to the formation of an inconsistent remolten seam, ultimately resulting in a failure to achieve separation.

Furthermore, the effect of circular laser beam oscillations superimposed on the processing feed direction was investigated. A well-focused and oscillating laser beam was employed to widen the kerf, while providing higher absorptivity compared to defocused beams. By means of circular oscillations, the results demonstrated a significant increase in the cutting speed, with a maximum rise from 0.13 to 0.20 m/s.

Future developments of the present research will be aimed at disclosing the effects of the different oscillating pattern on the process dynamics.

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AUTHOR DECLARATIONS

Conflict of Interest

The authors have no conflicts to disclose.

Author Contributions

Matteo Busatto: Conceptualization (equal); Data curation (equal); Formal analysis (equal); Investigation (equal); Methodology (equal); Writing – original draft (equal); Writing – review & editing (equal). Leonardo Caprio: Conceptualization (equal); Methodology (equal); Writing - review & editing (equal). Barbara Previtali: Conceptualization (equal); Supervision (equal); Writing – review & editing (equal).

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