# A detailed mathematical model of the human atrial cardiomyocyte: integration of electrophysiology and cardiomechanics

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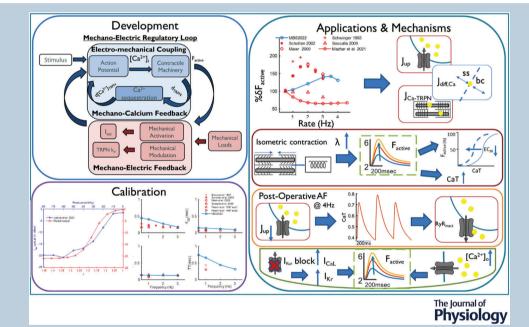
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Abstract Mechano-electric regulations (MER) play an important role in the maintenance of cardiac performance. Mechano-calcium and mechano-electric feedback (MCF and MEF) pathways adjust the cardiomyocyte contractile force according to mechanical perturbations and affects electro-mechanical coupling. MER integrates all these regulations in one unit resulting in a complex

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phenomenon. Computational modelling is a useful tool to accelerate the mechanistic understanding of complex experimental phenomena. We have developed a novel model that integrates the MER loop for human atrial cardiomyocytes with proper consideration of feedforward and feedback pathways. The model couples a modified version of the action potential (AP) Koivumäki model with the contraction model by Quarteroni group. The model simulates iso-sarcometric and isometric twitches and the feedback effects on AP and  $Ca^{2+}$ -handling. The model showed a biphasic response of  $Ca^{2+}$  transient (CaT) peak to increasing pacing rates and highlights the possible mechanisms involved. The model has shown a shift of the threshold for AP and CaT alternans from 4.6 to 4 Hz under post-operative atrial fibrillation, induced by depressed SERCA activity. The alternans incidence was dependent on a chain of mechanisms including RyRs availability time, MCF coupling, CaMKII phosphorylation, and the stretch levels. As a result, the model predicted a 10% slowdown of conduction velocity for a 20% stretch, suggesting a role of stretch in creation of substrate formation for atrial fibrillation. Overall, we conclude that the developed model provides a physiological CaT followed by a physiological twitch. This model can open pathways for the future studies of human atrial electromechanics.

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**Abstract figure legend** Three phases of novel human atrial electromechanically coupled model, MBS2023. Phase 1 deals with model development, the conceptual background, and the coupling of forward and backward pathways. Phase 2 is based on calibration of parameters related to contraction, stretch activated channels, and  $Ca^{2+}$  handling. Phase 3 is about the testing and validation of the model based on applications like rate adaptation, isometric contraction, the post-operative atrial fibrillation (PoAF) condition,  $I_{Kur}$  block, and extracellular  $Ca^{2+}$  variations.

## Key points

- With the availability of human atrial cellular data, interest in atrial-specific model integration has been enhanced.
- We have developed a detailed mathematical model of human atrial cardiomyocytes including the mechano-electric regulatory loop. The model has gone through calibration and evaluation phases against a wide collection of available human *in-vitro* data.
- The usefulness of the model for analysing clinical problems has been preliminaryly tested by simulating the increased incidence of Ca<sup>2+</sup> transient and action potential alternans at high rates in post-operative atrial fibrillation condition.
- The model determines the possible role of mechano-electric feedback in alternans incidence, which can increase vulnerability to atrial arrhythmias by varying stretch levels.
- We found that our physiologically accurate description of Ca<sup>2+</sup> handling can reproduce many experimental phenomena and can help to gain insights into the underlying pathophysiological mechanisms.

## Introduction

At each heartbeat, dynamic changes in the heart activity, induced by the mechanical load, are intrinsically regulated (via Frank Starling and Anrep effects) to maintain the cardiac output. Likewise, at the cellular level these mechanical changes regulate the electrical activity (Peyronnet et al., 2016, Quinn & Kohl, 2021) and the Ca<sup>2+</sup> handling (Calaghan et al., 1999, Neves et al., 2016) under both physiological and pathological conditions (Pfeiffer et al., 2014; Ravens, 2003; Taggart & Sutton, 1999; Varela

et al., 2021). The strong coupling of electrophysiology and cardiomechanics within the cardiomyocytes (CMs) is called electro-mechanical coupling (EM). Figure 1 shows the feedforward and feedback pathways that connect together to form the mechano-electric regulatory (MER) loop. Within the loop, the dynamic response of  $Ca^{2+}$  sequestration by troponin (TRPN) on the electrophysiology is referred to as the mechano-calcium feedback (MCF) (Rice et al., 1999). The changing contractile force affects the action potential (AP) via mechano-electric feedback (MEF) (Nazir et al., 1996; reviewed in Kohl et al., 2001).

EM coupling is a series of events triggered by membrane depolarization and the development of the AP, followed by  $Ca^{2+}$  release from the sarcoplasmic reticulum (SR) via the  $Ca^{2+}$ -induced- $Ca^{2+}$ -release (CICR) mechanism. The  $Ca^{2+}$  release is responsible for initiating the CM contraction by the activation of thin filaments, i.e. the regulatory units (RUs) followed by the cooperativity of RU proteins, and eventually leading to cross-bridge cycling (XB<sub>cy</sub>). This regulation hence depicts a complex phenomenon, making it challenging to understand and to quantify the role of individual components at the system level. The development of a sophisticated cardiac mathematical model can help to provide a better insight into these complex experimental phenomena (Vagos et al., 2018).

Several atrial computational models have been developed to improve the understanding of individual components of the EM coupling: the electrophysiology (reviewed in Heijman et al., 2021 and Wilhelms et al., 2013) and the Ca<sup>2+</sup> handling (reviewed in Sutanto et al., 2020). However, in the literature, the focus of cardiomechanics studies (reviewed in Niederer et al., 2019; Regazzoni et al., 2021) and models (reviewed in Lee et al., 2022; Trayanova et al., 2011; Varela et al., 2021) has been more on ventricular CMs (reviewed in Sutanto & Heijman, 2022). Similarly, MER integrated models were developed and were based either on ventricular

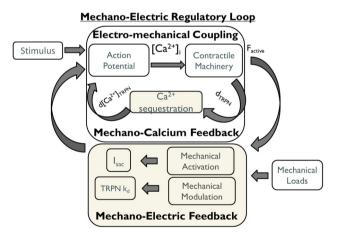


Figure 1. Mechano-electric regulatory loop along with feedforward (electro-mechanical coupling) and feedback (mechano-calcium and mechano-electric) pathways The application of action potential produces a Ca<sup>2+</sup> transient [Ca<sup>2+</sup>]<sub>i</sub> that is fed as an input to the contractile machinery. The active force ( $F_{active}$ ) generates a dynamic effect of troponin buffer (d<sub>TRPN</sub>) and triggers the mechano-calcium feedback by the sequestration of [Ca<sup>2+</sup>]<sub>i</sub> (d[Ca<sup>2+</sup>]<sub>TRPN</sub>). The varying mechanical length/loads induces change in mechanical activity either by mechanically activated channels (MACs) in the form of stretch activated current ( $I_{sac}$ ) or mechanically modulated channels (MMCs) via the  $k_d$  factor of troponin, Ca<sup>2+</sup> cytosolic buffer and is feedback to AP and Ca<sup>2+</sup> handling by MEF pathway. [Colour figure can be viewed at wileyonlinelibrary.com]

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electro-mechanics (Timmerman et al., 2017; Margara et al., 2021; Bartolucci et al., 2022) or on stem cell derived CMs (Forouzandehmehr et al., 2021). With the availability of human atrial cellular data, interest in atrial specific MER modelling has increased. The very first attempt to study the influence of human atrial mechanics during atrial fibrillation (AF) on the whole heart model was made by Land (Land et al., 2018). However, not every contraction parameter was re-tuned according to human atrial physiological contraction data. Relatively few atrial cellular MER integrated models have been developed; for instance, the integrated model by Zhan (Zhan et al., 2013; Zhan et al., 2019) is more focused on the impact of fibroblast on coupling and contraction. Another EM model was developed by Kuijpers (Kuijpers et al., 2007) and later by Brocklehurst (Brocklehurst et al., 2015); however, its cell level mechanisms were not validated quantitatively against any human atrial in-vitro data. Importantly, all the integrated models used the same human atrial AP model, i.e. the Courtemanche (CRN) model (Courtemanche et al., 1998).

Following the current MER modeling approach, in one of our previous studies (Mazhar et al., 2021), we developed an integrated model of human atrial CM by coupling the Courtemanche AP model and the contraction model developed by the Quarteroni group, the RDQ2020 model (Regazzoni et al., 2020). However, this coupled and calibrated model was not able to reproduce one important phenomenon: the force-frequency relationship for human atrial CMs. The reason for this discrepancy was identified as the Ca<sup>2+</sup> -transient rate adaptation trend, which in turn probably depends on the atrial specific physiological phenomena of CICR propagation, from the membrane to the centre of the CM in a  $Ca^{2+}$  wave-like fashion by a 'fire-diffuse-fire-mechanism' (Blatter, 2017; Bootman et al., 2006). Hence, there is a need to have a model with calibrated Ca<sup>2+</sup> handling dynamics that remains valid upon rate changes.

We aimed to develop a human atrial cellular model with MER integration that is computationally efficient and exhibits a detailed temporal Ca<sup>2+</sup> handling description, valid over a physiological range of pacing rates. In addition, to our knowledge, this is the first human atrial model that includes the MER loop and is calibrated against a wide collection of available human in-vitro atrial data. The model was designed to present coupling and mechanical load effects on electrophysiology, on  $Ca^{2+}$  transient (CaT), and on the active force ( $F_{active}$ ). In this way, it can be used to determine the possible role of feedbacks MCF and MEF in the initiation of atrial arrhythmias by varying stretch levels. The clinical utility of the model has been preliminaryly tested by simulating the increased incidence of CaT and AP alternans at high rates in post-operative atrial fibrillation (PoAF) conditions. Our model is named Mazhar-Bartolucci-Severi 2023 (MBS2023).

## Methods

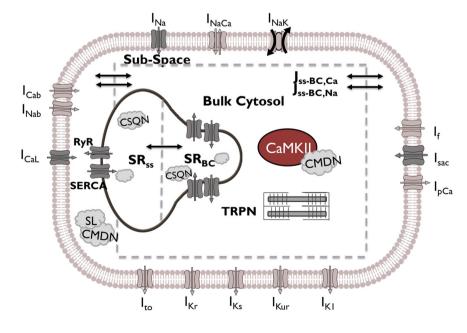
With respect to the MER loop formation between AP and contraction modules, we identified as the components of the coupling a feedforwad and a feedback pathway. The former is the input of CaT to the contraction module, the latter is the sequestration of Ca<sup>2+</sup> bound to troponin (MCF) from the free cytosolic Ca<sup>2+</sup>. In this way, a strong bidirectional coupling was ensured in the MBS2023 model. The section 'Electro-mechanical coupling and mechano-calcium feedback' highlights this stepwise approach for introducing coupling and the MCF effect in the model. Moreover, the calibration of contraction and Ca<sup>2+</sup> handling parameters was carried out to obtain a model in agreement with human atrial in-vitro data (references in Tables A2 and A3). MBS2023 is simulated in two operating modes: (1) iso-sarcometric, where the sarcomere length (SL) was fixed as is usually found in-vitro experiments, and (2) isometric, where for a given SL the preload effect was varied using the stretch ratio ( $\lambda$ ). The model was switched to isometric operating mode in order to simulate the stretch-induced MEF response and also to determine its role in alternans vulnerability.

## Cell structure of the model

In MBS2023, the bidirectional coupling was achieved with two modules: the AP module obtained from a considerable revision of the human atrial cell model Koivumäki 2011 (Koivumäki et al., 2011), KM2011; and a mean-field approximation-based contraction module, RDQ2020 (Regazzoni et al., 2020). The aim of the MBS2023 model was to combine the complex  $Ca^{2+}$ handling dynamics of atria (Blatter et al., 2003; Bootman et al., 2011; Hüser et al., 1996; Hatem et al., 1997) with a computationally efficient modelling approach for human atrial CM. Based on the KM2011 AP model, we reduced the multi-compartmental structure of cytosol and SR in KM2011 and we were left with two compartments only, i.e. the sub-space (ss) and the bulk cytosol (bc). A schematic representation of the updated compartmental organization of our model is shown in Fig. 2.

The cell dimensions for the parent AP model KM2011 were adopted from Nygren (Nygren et al., 1998). In comparison to the human atrial cell dimensions (Neef et al., 2010; n = 115), the model cell length was slightly longer (122.051  $\mu$ m vs. 101.15  $\pm$  1.5  $\mu$ m) and the radius was shorter (6.5  $\mu$ m vs. 7.4  $\pm$  0.3  $\mu$ m). The total cell capacitance (50 pF) seems to be lying near the lower bound of values reported in human atrial experiments (see Table A1). Interestingly, we found a very large variability in such data, ranging from 29.6  $\pm$  1.8 pF, as reported by the Nattel lab (Wang et al., 1999) up to 114.8  $\pm$  5.9 pF reported by Voigt et al. (Voigt et al., 2012). This variability in cell capacitance can be associated with the T-tubule density present in the cell, which affects the capacitance-to-longitudinal surface area (C/A), reported to range from 1% to 10% in human atrial cells (n = 50, Fakuade et al., 2021). The dimensions for the MBS2023 cell model were modified in accordance with the experimental values (Neef et al., 2010). Since there was a 12% increase in radius, with respect to KM2011, we made a similar increase in the capacitance, up to 56 pF. The low C/A ratio ( $\cong$  1%) implies an assumption of a low T-tubule density.

The cytosolic space was re-arranged in such a way that the ss compartment was at a distance of 0.02  $\mu$ m from the membrane (as in KM2011) and the rest of the area in the bulk/centre was summed to form the *bc* compartment



## Figure 2. Schematic diagram of human atrial myocyte model

The figure shows two cytosolic compartments, the subspace (ss) and bulk cytosol (bc). The *bc* compartment is enclosed inside a grey dashed line box. The sarcoplasmic reticulum (SR) also includes two compartments SR<sub>ss</sub> and SR<sub>bc</sub> with their own release (RyR) and uptake (SERCA) units. The cytosolic buffers are troponin (TRPN), calmodulin (CMDN), and sarcolemmal phospholipids (SL). The sarcolemmal currents and fluxes in dark grey are modified/ reformulated/added in our novel model MBS2023. [Colour figure can be viewed at wileyonlinelibrary.com]

so that the *bc* compartment was located at a distance of 6.6  $\mu$ m from the ss. In this way, the *bc* compartment volume ( $V_{bc}$ ) was 163 times larger than the volume of ss compartment ( $V_{ss}$ ). Similarly, SR was divided into two compartments, i.e. SR<sub>ss</sub> lying near the membrane and SR<sub>bc</sub> more towards the centre, each having its own release RyR and uptake SERCA units. The total volume of SR was kept the same as in KM2011 whereas the volume dedicated to ss ( $V_{SRss}$ ) and *bc* ( $V_{SRbc}$ ) was redefined.  $V_{SRbc}$  was then formed from  $\frac{3}{4}$  of the cytosolic *bc* compartment and  $V_{SRss}$  was the sum of cytosolic subspace volume  $V_{ss}$  and  $\frac{1}{4}$  of bc. Hence, the  $V_{SRbc}$  was 1.26 times larger than  $V_{SRss}$ . In this way, the contribution of Ca<sup>2+</sup> storage in SR<sub>ss</sub> was enhanced towards the formation of global cytosolic CaTs.

The diffusion of  $Ca^{2+}$  and  $Na^+$  in the cytosolic compartments is also shown in Fig. 2. In addition, all of the intracellular electrolyte concentrations were dynamic except  $[K^+]_i$  which was fixed at 134 mM since it was necessary to have a steady-state stable value for the electrolyte concentrations (Barral et al., 2022).

### Ca<sup>2+</sup> in the cytosol

In KM2011, cytosolic Ca<sup>2+</sup> diffusion from the ss to bc region was an analytical diffusion equation and the diffusion distance  $x_{ss-bc}$  was from the centre of ss to the centre of the first bc compartment. Using the same definition of  $x_{ss-bc}$  for the simplified bc compartment of MBS2023 results into a 4-fold increase in the diffusion distance, which yields a significant slowing of diffusion. In contrast, Ca<sup>2+</sup> diffusion is faster in atria than in ventricles and therefore it cannot be a slow process physiologically (Tanaami et al., 2005). Hence, in our model  $x_{\text{ss-bc}}$  was kept as it is in KM2011, i.e. 0.8125  $\mu$ m to give a fast diffusion of Ca<sup>2+</sup> in the cytosol. KM2011 has a single arbitrary cytosolic buffer whose dissociation coefficient  $K_d$  and concentration values correspond to the characteristics of calmodulin (CMDN). For bidirectional coupling, TRPN was included in the cytosol with the existing CMDN buffer (Ji et al., 2015). The bidirectional coupling of the AP and contraction modules has been described in detail in the section on 'Electromechanical coupling and mechano-calcium feedback'. Apart from this, the Ca<sup>2+</sup> diffusion via mobile buffer in the bc compartment was no longer needed in our model since MBS2023 now has a single bc compartment. The absence of mobile buffers and functional T-tubules in the model makes Ca-SR release the main source of spatial spread of Ca<sup>2+</sup> signal other than the diffusion seen in experiments (Walden et al., 2008). Hence, the differential equations for both free intracellular  $Ca^{2+}$  ([ $Ca^{2+}$ ]<sub>i</sub>) and  $Ca^{2+}$  in SR ( $[Ca^{2+}]_{SR}$ ) were modified and are listed in Appendix 1.

#### Ca<sup>2+</sup> transport in the SR

Within SR, the bidirectional SERCA pumping and buffering has the same formulation as the KM2011 model. Other than this, KM2011 adopted the numerical value of the SERCA buffer affinity from ventricular myocytes reported in Shannon et al. (2000). However, an increased SERCA activity with fast uptake dynamics was found in atria (Walden et al., 2009). In KM2011 model, the SERCA affinity formulation is based on the regulatory protein phospholamban (PLB) and SERCA expression levels (Koivumäki et al., 2009). A low level of SERCA inhibitory protein (PLB) along with high SERCA expression eventually leads to the raised affinity of SERCA buffer for human atria (Bokník et al., 1999). Moreover, less PLB inhibition of SERCA directly affects the SERCA dynamics too. Thus, in our model, the affinity of SERCA buffer was modified to meet human atrial SERCA activity levels (Koivumäki et al., 2014, Supplementary Section).

On the release side, the  $Ca^{2+}$  release from RyR channels has a Hodgkin-Huxley formulation in KM2011 that consists of activation, inactivation and an adaptation gate. The adaptation gate adapts the RyR open probability according to the intracellular Ca<sup>2+</sup> environment. Since we have modified the intracellular Ca<sup>2+</sup> levels, the RyR adaptation level was modified accordingly, to maintain a physiological Ca<sup>2+</sup> concentration both in the cytosol and in the SR. In the bc compartment, MCF was introduced by the inclusion of troponin-like fast buffer and this resulted in a depressed Ca<sup>2+</sup> peak and slower relaxation kinetics (Michailova et al., 2002; Smith et al., 2019). This lower  $Ca^{2+}$  peak was the consequence of less release flux,  $J_{\rm rel}$ , from the RyR in the bc compartment. Hence, an adjustment of  $Ca^{2+}$  release was required to give a physiological level of CaT. For the systolic value, we modified adaptation parameters of the bc compartment, i.e. the saturation level and the sensitivity of adaptation. The saturation value was reduced so that the release gate activation can saturate at lower  $[Ca^{2+}]_{bc}$  values and the adaptation sensitivity was increased. On the other hand, the relaxation kinetics were more dependent on the contraction dynamics modelled by RDQ2020. Therefore, an automatic optimization was carried out for the time related biomarkers relevant to contraction. This is elaborated in the section on 'Optimization of contraction parameters'.

At higher beating frequency, the incidence of CaT and AP alternans was observed in the model. In general, the vulnerability of  $Ca^{2+}$  induced alternans is often associated with steep non-linear  $[Ca^{2+}]_{SR}$  dependence on  $J_{rel}$  (Díaz et al., 2004; Xie et al., 2008). However, cytosolic  $Ca^{2+}$  alternans were also observed without any significant concurrent oscillations in diastolic  $[Ca^{2+}]_{SR}$  in single myocytes (Hüser et al., 2000; Picht et al., 2006). For

instance, Shkryl et al. (2012) observed experimentally in rabbit atrial myocytes that the inactivation kinetics and the refractoriness of RyR release was the key mechanism behind Ca<sup>2+</sup> alternans generation. Similar results were observed in experiments in human atrial cells (Llach et al., 2011) where the mechanism proposed for alternans was the slow recovery of release. For this reason, we modified the refractoriness of RyR gate, both in ss and bc compartments by increasing the recovery from inactivation time from its standard value of 200 ms (Stern et al., 1999). This was achieved in the model by increasing the inactivation time constants for the recovery period from 15 to 450 ms and from 87.5 to 656 ms for ss and bc compartments, respectively. Moreover, the inactivation time constant for the ss compartment was also updated from 15 to 60 ms. In this way, the  $J_{\rm rel}$  fluxes for ss and bc had a slow decay time and hence a comparable total time duration.

#### **Alternans protocol**

To characterize the AP and CaT alternans, MBS2023 requires faster rates under control condition. To quantify the role of potential parameters both from AP and contractility in the incidence of alternans at higher rates, we used a pacing protocol similar to the one used in a clinical study (Narayan et al., 2011) and in computational modelling (Zile & Trayanova, 2018). Under iso-sarcometric conditions, at a basic cycle length (BCL) of 750 ms, the model is paced for 1000 beats to reach the steady state condition. Using this steady state value as the initial condition, the model was paced further for 74 beats with BCL decreasing from 500 to 200 ms in 50 ms steps. From our literature study, the parameters that were analysed using this pacing protocol were Ca<sup>2+</sup>-handling-related (the RyR inactivation time (tau<sub>inact</sub>), SERCA pumping rate ( $k_{4-\text{SERCA}}$ ), SERCA expression ([SERCA])), contraction-related (thin filament  $Ca^{2+}$  binding sensitivity ( $k_{d-TRPN}$ ), thin filament rate constant ( $K_{\text{off}}$ ), cross bridge cycling (XB<sub>cy</sub>), rate constant  $(K_{\text{basic}})$ ), the coupling effect, i.e. MCF, and CaMKII-related (expression of CaMKII ([CaMKII])). For the analysis, all the parameters were varied between a range of values obtained by increasing values up to 2-fold and reducing them 0.5-fold.

Based on a comprehensive alternans analysis (Zile & Trayanova, 2018), we defined a similar measure to quantify the magnitude of AP and CaT alternans. Alternans normalized magnitude (ANM) was calculated as the mean of change over the consecutive last 10 pairs of beats during a BCL divided by the mean value over the last 10 beats and is given as:

ANM = 
$$\frac{\sum_{b=63}^{73} |X_b - X_{b+1}| / 10}{\operatorname{avg}(X)_{b=64 \ to \ 74}}$$

where X is AP or CaT biomarkers like APD<sub>90</sub>, APD<sub>50</sub>, and CaT<sub>peak</sub>, thus giving the three alternans measures APD<sub>90</sub>-ANM, APD<sub>50</sub>-ANM and CaT-ANM. The threshold used for alternans to be considered significant was ANM > 5%.

#### Sarcolemmal ion currents

Few of the sarcolemmal currents were re-formulated, modified or newly included in comparison to the KM2011 model, as shown in dark grey in the schematic of Fig. 2. The inward sodium current  $(I_{Na})$  was reformulated according to new experimental data as described in Skibsbye et al. (2016). The L-type  $Ca^{2+}$  current ( $I_{CaL}$ ) gating was modified as explained in Koivumäki et al. (2014) (KM2014). In comparison to KM2014, steady state inactivation for the fast VDI gate was adopted from KM2011 to avoid irregular reactivation of  $I_{CaL}$  at the low frequency of 0.5 Hz. As a result, the I-V curves for both KM2011 and KM2014 versions were compared with data (Li et al., 1997) under voltage clamp protocol given in Fig. A1. Our modified formulation of  $I_{CaL}$  is in better agreement with the experimental data in comparison to KM2014. Apart from this, all other minor modifications in conductance values have been listed in Table A5, section A.

The mechano-activated channels (MACs) allow the stretch-activated currents  $(I_{sac})$  to flow. On the cellular scale, MACs are the stretch activated gates that are classified according to the stimulus type by which they are activated, i.e. either by using mechanical stimulus alone or by a combination of different stimuli. However, in this work, we have considered those channels that are activated by mechanical stimulus alone. Very often, Isac is subdivided into  $K^+$  selective ( $I_{K0}$ ) and non-selective  $(I_{\rm ns})$  cations channels.  $I_{\rm ns}$  allows Na<sup>+</sup>, Ca<sup>2+</sup> and K<sup>+</sup> to enter the cells and its reversal potential, based on the relative permeability of ion channels, is halfway between the resting membrane potential (RMP) and the plateau potential, i.e. 0 and -50 mV, and has a near-linear voltage dependence (Kim, 1993). Therefore, the activation of  $I_{ns}$  can depolarize the cells, but during the plateau phase can cause the repolarization of the AP. In contrast, the reversal potential of  $I_{K0}$  is close to RMP; therefore it can modulate the repolarization phase more than the RMP of the AP. In the model, the  $I_{\rm ns}$  current is permeable to Na<sup>+</sup>, K<sup>+</sup> and Ca<sup>2+</sup> and it is defined as:

$$I_{\rm ns} = I_{\rm ns,Na} + I_{\rm ns,K} + I_{\rm ns,Ca}.$$

where  $I_{ns}$  is the ion-specific contribution of Na<sup>+</sup>, K<sup>+</sup> or Ca<sup>2+</sup>. The driving force of each  $I_{ns,X}$  (where X refers to

ion contribution of Na<sup>+</sup>, K<sup>+</sup> or Ca<sup>2+</sup>) is modelled by the Goldman-Hodgin-Katz (GHK) based formulation:

$$I_{\rm ns,X} = P_{\rm X} \ g_{\rm ns} \frac{z_{\rm X}^2 \ F^2 V_{\rm m}}{RT} \ \frac{[X]_{\rm i} - [X]_{\rm o} \exp\left(-\frac{z_{\rm X} F \ V_{\rm m}}{RT}\right)}{1 - \exp\left(-\frac{z_{\rm X} F \ V_{\rm m}}{RT}\right)},$$

where  $P_X$  is the permeability of X ion Na<sup>+</sup>, K<sup>+</sup> and Ca<sup>2+</sup>,  $z_{\rm X}$  is the ion valency for each ion, and the rest of the constants are defined as in KM2011. For our model the permeability ratio  $P_{\text{Na}}: P_{\text{Ca}}: P_{\text{K}}$  was fixed as 1:1:1, which gives a linear I-V relationship for  $I_{ns}$  current with a reversal potential near 0 mV (Fig. A7, red triangles) as was shown for rat atrial myocytes (Kim, 1993). A sensitivity analysis, based on a 3-fold increment of each ion permeability at a time, predicted that changes in  $P_{Na}$  and  $P_K$  can increase (25 mV) and reduce (-25 mV) the reversal potential of  $I_{\rm ns}$ , respectively, as shown in Fig. A7 (yellow plus marks and blue squares, respectively). Moreover, Kim suggested the  $P_{Ca}/P_{K}$  ratio to be equal to 0.9 (Kim, 1993) and we have adopted this ratio from Kuijper et al. as 1 (Kuijper et al., 2007). Thus, we opted for an equal permeability ratio for all the three ion channels in our model. The effect of stretch change was introduced in the  $I_{ns}$  current by a conductance  $g_{ns}$  equation, given as,

$$g_{\rm ns} = \frac{g_{\rm ns,max}}{1 + K_{\rm ns} \exp\left[-\alpha_{\rm ns}\left(\lambda - 1\right)\right]}$$

where  $g_{\rm ns,max}$  is the maximal conductance,  $K_{\rm ns}$  is the parameter that sets the amount of current in the absence of stretch ( $\lambda = 1$ ), and  $\alpha_{\rm ns}$  is the sensitivity to stretch. All the parameter values are listed in Table A5, section C in Appendix 3. Furthermore, the  $I_{\rm ns}$  current formulation was included in electrolyte equations as:

$$\frac{d[Na^{+}]_{i}}{dt} = \frac{-(I_{Na} + I_{Nab} + 3I_{NaK} + 3I_{NaCa} + I_{fNa} + I_{ns,Na}}{V_{ss} F}$$
$$\frac{d[K^{+}]}{dt} = 0,$$
$$\frac{d[Ca^{2+}]_{ss}}{dt} = \frac{-I_{CaL} - I_{Cab} - I_{CaP} + 2I_{NaCa} - I_{ns,Ca}}{2V_{ss}F},$$

where  $V_{ss}$  is the volume occupied by the ss compartment, and the rest of the constants are the same as the original model. For  $I_{K0}$  current, we have considered the formulation for rat ventricular myocytes by Niederer and Smith (Niederer & Smith, 2007) with a change in offset for  $\gamma_{K0}$  that is the measure of strain dependence for conductance. A small offset was included to nullify the effect of  $I_{ns}$  current at basal rate when CMs were not in a stretched state. Briefly, the conductance is a function of strain and is modeled as:

$$\gamma_{\rm K0} = \beta_{\rm K0} \ (\lambda - 1) + 0.001,$$

and  $I_{\rm K0}$  becomes

$$I_{\rm K0} = g_{\rm K0} \; \frac{\gamma_{\rm K0}}{1 + \exp{-\frac{10+V}{45}}} \left(V - E_{\rm K}\right).$$

All the other constants were the same as in Niederer & Smith (2007).

Parameter validation for  $I_{ns}$ . We have validated the  $I_{ns}$ current parameters based on the average current values on negative pressure ranging from 0 to 80 mmHg extracted from Jakob et al. (2021). The measurements are obtained from 11 tissue donors and 110 cells. The model output for the  $I_{ns}$  (red line) fit to the experimental data is shown in Fig. A5 (blue line), and it increases up to  $\lambda = 1.25$  followed by a plateau that fits qualitatively with the experimental average current values, i.e. -20 pA. The calibrated set of parameters thus obtained are listed in Table A5, section C, which shows a high sensitivity towards stretch, i.e.  $\alpha_{ns} = 30$  and a maximum conductance of  $g_{nsmax} = 0.461e-6$  nS. The permeability ratio is made the same for all three ion channels. Using these calibrated parameters of  $I_{ns}$ , we evaluated the effect of increasing stretch level on the MBS2023 model in the Results section.

# Electro-mechanical coupling and mechano-calcium feedback

Using our modified version of KM2011 as the AP module and RDQ2020 as the contraction module, we built a coupled model of human atrial CMs with MER loop. The CaT from the AP module served as an input to the contraction module in a feedforward manner and the dynamic effect of Ca<sup>2+</sup> bound to TRPN ( $\frac{d[Ca^{2+}]_{TRPN}}{dt}$ ) from the contraction module was the feedback to the free CaT in the AP module, i.e. the MCF. In this way, a bidirectional strong electro-mechanical coupling was ensured. Mathematically, the integration of the AP and contraction modules was performed in accordance with Ji et al. (2015) and later by Timmermann et al. (2017). We defined the amount of calcium bound to TRPN  $([Ca^{2+}]_{TRPN})$  as the fraction of TRPN units with Ca<sup>2+</sup> bound to its regulatory binding site  $(Ca_{TRPN})$  multiplied by a constant maximum concentration of Ca<sup>2+</sup> that can bind to TRPN ( $[Ca^{2+}]_{TRPN,max} = 0.56 \text{ mM}$ ), leading to:

$$\frac{\mathrm{d}[\mathrm{C}a^{2+}]_{\mathrm{TRPN}}}{\mathrm{d}t} = [\mathrm{C}a^{2+}]_{\mathrm{TRPN},\mathrm{max}} \frac{\mathrm{d}\mathrm{C}a_{\mathrm{TRPN}}}{\mathrm{d}t},\qquad(1)$$

where the  $[Ca^{2+}]_{TRPN,max}$  was set to the minimum concentration allowed for 50% of accessible volume (Fabiato et al., 1983; Robertson et al., 1981). The differential equation for  $[Ca^{2+}]$  was modified after adding the troponin feedback from the contraction module.

 $Ca_{TRPN}$  from Eqn (1) is computed from RDQ2020 as follows:

$$Ca_{\text{TRPN}} = B_{\text{so}}X_{\text{SO}}\left(SL\right) + B_{\text{nso}}\left(1 - X_{\text{SO}}\left(SL\right)\right), \quad (2)$$

where  $B_{so}$  and  $B_{nso}$  are the ratio of bounded TRPN units in the single-overlap zone (SO) and in the non-single-overlap zone (NSO), respectively, and  $X_{so}(SL)$ is the function of the size of the single overlap zone that models the effect of SL changes. On the one hand,  $B_{so}$  was already modelled in the original RDQ2020 model as:

$$B_{\rm so} = \sum_{\alpha,\beta,\delta} \pi \; (\alpha,\beta,\delta,B)$$

where  $\pi(\alpha, \beta, \delta, B)$  is the ratio of TRPN being permissive, regardless of the state of the tropomyosin (Tm) triplet  $\alpha, \beta, \delta$ . On the other hand, we introduced the variable  $B_{nso}$ , whose dynamic is modelled as:

$$\frac{\mathrm{d}B_{\mathrm{nso}}}{\mathrm{d}t} = k_{\mathrm{off}} \frac{\left[\mathrm{Ca}^{2+}\right]_{i}}{k_{\mathrm{D}}\left(\mathrm{SL}\right)} \left(1 - B_{\mathrm{nso}}\right) - k_{\mathrm{off}} B_{\mathrm{nso}}.$$

Experimental data. The contraction parameters in the RDQ2020 model were calibrated according to ventricular human data recorded at body temperature (see Table 3, Regazzoni et al., 2020). Hence, to adopt human atrial-like physiology, a re-calibration of contraction parameters was required. This calibration was based on a few biomarkers extracted from the *in-vitro* experimental data available for human atrial trabeculae, summarized in Table A2. Instead of employing the *in-vitro* data from single cardiac myocytes, the calibration data for our model was based on tissue level preparations. The removal of CM from its normal environment has a significant impact on its electrical and mechanical functionality because of its loss of interaction with several other cells (Pfeiffer et al., 2014). Therefore, tissue level data allow the myocytes to be studied in an environment that more closely mimics in-vivo conditions.

**Optimization of contraction parameters.** The major difference between atrial and ventricular contraction lies in the contractile protein myosin and its iso-enzymatic variety (Reiser et al., 2001; Narolska et al., 2005; Morano et al., 1988). Contractile protein myosin is the major determinant of contractility since it can modulate the ATPase activity and maximum shortening velocity ( $v_{max}$ ). Accordingly, the contraction and relaxation kinetics are faster in atria than ventricles (Piroddi et al., 2007). Moreover, the Ca<sup>2+</sup> sensitivity of atrial fibre is also less than ventricles (Morano et al., 1988).

Based on these observations, the contraction dynamics of regulatory unit (RU), i.e. the TRPN and Tm units were modified. The dynamics were based on transition rates ( $k_{\text{off}}$  and  $k_{\text{basic}}$ ), the cooperativity ( $\gamma$ ) and the  $Ca^{2+}$  sensitivity ( $k_d$ ). The optimized set of parameters was obtained on the basis of time-related biomarkers of  $F_{\text{active}}$ , i.e. relaxation time at 50% and 90% of peak values ( $rt_{50Factive}$ ,  $rt_{90Factive}$ ), time to  $F_{max}$  ( $ttp_{Factive}$ ), and twitch time (TT<sub>Factive</sub>). This optimization was carried out at basal frequency of 1 Hz and then the optimized model parameters were used for simulations at other frequencies (see the comparison of model output with experimental data normalized by 1 Hz in Fig. A2). The late relaxation phase of  $F_{\text{active}}$ , i.e.  $rt_{90\text{Factive}}$  demonstrates rate dependency and the model has reproduced it. In contrast, the ttp<sub>Factive</sub> and rt<sub>50Factive</sub> shows a flat behaviour with respect to rate and it should be noted that the only available experimental data (Sossalla et al., 2010) shows a flat rate dependency for overall relaxation phases of  $F_{\text{active}}$ . The model output for  $rt_{50\text{Factive}}$  shows a contrasting behaviour with respect to the data and this was a expected result after looking into the rate dependent shortening of rt<sub>90Factive</sub>. The human *in-vitro* data collection used for this calibration is listed in Table A2. Similarly, biomarkers were also computed for the CaT - the diastolic level ( $CaT_{min}$ ), the systolic value ( $CaT_{max}$ ), the amplitude  $(CaT_{amp})$  relaxation time at 50% of  $CaT_{max}$  (rt<sub>50Ca</sub>), time to  $CaT_{max}$  (ttp<sub>Ca</sub>), total time duration (TT<sub>Ca</sub>), and the decay time  $(\tau)$  – and compared with a small set of data from Table A2. The calibration was performed using the Matlab (The MathWorks, Natick, MA) fminsearch optimization function to find local minima. In this way, the model parameters were tuned in accordance with the human atrial *in-vitro* data to adopt atrial-like physiology.

Other parameters based on cross-bridge cycling (XB<sub>cy</sub>), i.e.  $\mu_{f_p}^0$ ,  $\mu_{f_p}^1$ ,  $\alpha$ , and  $r_0$  were also calibrated by following the procedure presented in Regazzoni et al. (2021). This calibration was based on difference in  $\nu_{max}$  of XB<sub>cy</sub> for human atria from ventricles (Kuijpers et al., 2007). Using the new  $\nu_{max}$  value, based on the best fit with force-velocity hyperbolic curve we re-calibrated the XB<sub>cy</sub> parameters (Table A2).

#### Modelling mechano-electric feedback

The model simulates MEF-based phenomena by including mechanical activity of the CMs. MEF was modelled by using a three-element rheological scheme (Solovyova et al., 2002) also shown in Fig. 3. The contractile element (CE) is responsible for generating  $F_{\text{active}}$  and is the one produced by RDQ2020. The element has a length of SL and a resting length, SL<sub>r</sub>, i.e. the length when the segment is relaxed and no force is applied. The force developed by CE unit is given as:

$$F_{\rm CE} = f_{\rm CE} F_{\rm norm} \left( {\rm SL}, {\rm Ca}^{2+}, \nu \right), \qquad (3)$$

where  $f_{CE}$  is a scaling factor and  $F_{norm}$  is the normalized active force generated by RDQ2020 that is dependent

on SL, Ca<sup>2+</sup> dynamics, and the shortening velocity v of the XBs. The MEF-based mechanical modulation of Ca<sup>2+</sup> was already modelled in RDQ2020, since Ca-TRPN sensitivity  $k_d$  was a function of SL. The passive elastic effect is introduced in two forms, i.e. by the series (SE) and parallel (PE) elements. Mathematically, the passive forces developed by SE and PE were modelled as non-linearly related to lengths  $l_{\text{PE}}$  and  $l_{\text{SE}}$ :

$$F_{\rm SE} = f_{\rm SE} \left\{ \exp \left[ k_{\rm SE} \left( l_{\rm SE} - l_{\rm SEr} \right) \right] - 1 \right\}$$
(4)

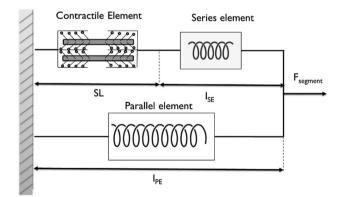
and

$$F_{\rm PE} = f_{\rm PE} \, \left\{ \exp \left[ k_{\rm SE} \, (l_{\rm PE} - l_{\rm PEr}) \right] - 1 \right\}, \tag{5}$$

where  $f_{\text{SE}}$  and  $f_{\text{PE}}$  are the scaling factors,  $k_{\text{SE}}$  and  $k_{\text{PE}}$  are the spring constants, and  $l_{\text{SEr}}$  and  $l_{\text{PEr}}$  are the resting lengths for each element. Under mechanical equilibrium, the total force of the mechanical segment ( $F_{\text{segment}}$ ) is computed as  $F_{\text{segment}} = F_{\text{PE}} + F_{\text{SE}}$ ; length of the segment as  $l_{\text{PE}} = l_{\text{SE}} + SL$ ; and the series forces as  $F_{\text{SE}} = F_{\text{CE}}$ .

The three-element mechanical model equations were solved by using the numerical scheme described in Solovyova et al. (2002). Two new state variables were integrated into the model, i.e.  $l_1 = SL - SL_r$  and  $l_2 = l_{\rm PE} - l_{\rm PEr}$  and eqns (4) and (5), were modified accordingly. Moreover, an assumption on resting lengths was made,  $l_{\rm PEr} = SL_r$ , and, as a consequence,  $l_{\rm SEr} = 0$ . Based on this, the mechanical state of each element is defined by  $l_1$ ,  $l_2$ , v,  $\frac{dl_1}{dt}$ ,  $\frac{dl_2}{dt}$ ,  $\frac{dv}{dt}$ . Under iso-sarcometric condition for  $\lambda = 1$ , the lengths of contractile and parallel elements were held constant, i.e.  $\frac{dl_1}{dt}$ ,  $\frac{dl_2}{dt}$ ,  $\frac{dv}{dt} = 0$  whereas, for isometric simulations  $\frac{dl_2}{dt} = 0$ ,  $\frac{dl_1}{dt} = -v$ , and the sarcomere shortening velocity v was modelled as:

$$\frac{\mathrm{d}\nu}{\mathrm{d}t} = \frac{F_{\mathrm{CE}} - F_{\mathrm{SE}}}{Ir} \tag{6}$$



## Figure 3. Three element rheological scheme for modelling the mechanical activity of the cell in MBS2023 model

The scheme consists of a contractile element of length SL from RDQ2020 module, responsible for active force development and two passive elements, the series (SE), and the parallel element (PE) of length  $I_{SE}$  and  $I_{PE}$ , respectively.

where *Ir* is the inertial parameter, and all other constant parameter values are listed in Table A5, part C.

# CaMKII based phosphorylation of ionic currents and Ca<sup>2+</sup> handling

Ca<sup>2+</sup>/CMDN-dependent kinase II (CaMKII) is a multifunctional serine/threonine kinase expressed ubiquitously in the CMs (reviewed in Bers & Grandi, 2009). CaMKII plays an important role in regulating cardiac excitability and contractility. Ca2+ mediated modulation of ion channels is carried out either through direct Ca<sup>2+</sup>-CMDN binding or by the activation of CaMKII. CaMKII can phosphorylate numerous targets involved in ion channel gating,  ${\rm C}a^{2+}$  handling, and in regulatory proteins. In the literature, many studies have been performed on the biochemically detailed models of CaMKII signalling but the focus of these studies was on ventricular electrophysiology (Soltis & Saurceman 2010; Morotti et al., 2021). As the MEF integration differs among the atrial and ventricular CMs because of differences in ultrastructure and Ca<sup>2+</sup> homeostasis, therefore it is crucial to study the atrial model integrated with CaMKII modulated targets, based on atrial physiology and pathophysiology. For atrial CMs, Zhao et al. have developed an atrial cell model including CaMKII activation to explore the mechanism induced by atrial arrhythmia under oxidative stress (Zhao et al., 2020). Another model by Ni et al. integrates electrophysiology and Ca<sup>2+</sup> handling with a biochemically detailed systems model of upstream signalling pathways, i.e. protein kinase A (PKA) and CaMKII in human atria (Ni et al., 2023).

In human atrial CMs, CaMKII acts on diverse targets in ion channels and  $Ca^{2+}$  handling as listed in Table 1. The CaMKII model was adopted from O'Hara ventricular model formulation (O'Hara et al., 2011), where each target *T* is divided into two parts,

$$T = (1 - \phi_{\text{CAMK}}) T_{\text{NP}} + \phi_{\text{CAMK}} T_{\text{CaMK}},$$

where  $T_{\text{NP}}$  is the non-phosphorylated part of the target,  $T_{\text{CaMK}}$  is the CaMKII modulated part of the target, and the  $\phi_{\text{CAMK}}$  is the proportion affected by CaMKII and is given as:

$$\phi_{\text{CaMK}} = rac{1}{1 + rac{K_{\text{m,CaMK}}}{\text{CaMK}_{\text{active}}}}$$

where  $CaMK_{active}$  is the fraction of active CaMKII binding sites,  $K_{m,CaMK}$  is the half-saturation coefficient, which was the same as the O'Hara CaMKII model. CaMKII<sub>active</sub> is given as:

$$CaM K_{active} = CaMK_{bound} + CaMK_{trap}$$

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Target	Effect of CaMKII-dependent phosphorylation	Reference
Fast I <sub>Na</sub>	Slow inactivation gate slowed by 1.46 times	(Zhang et al., 2015)
I <sub>CaL</sub>	1. Current increased by 40%	(Christ et al., 2004)
	2. Voltage dependent inactivation gate slowed by 1.5	
/ <sub>t</sub>	1. Current reduced by 10%	(Tessier et al., 1999)
	2. Inactivation time slowed by 1.56 times	
I <sub>Kur</sub>	Current increased by 57%	
RyR	Opening probability increased and leak current increased by 2.5 times.	(Neef et al., 2010; Voigt et al., 2012)
SERCA	Pumping activity increased by 27% by a direct effect and reduction in PLB	(Koivumaki et al., 2009)
	inhibition, an indirect enhancement, of forward pumping rate by 70%.	

Table 1. Effect of CaMKII based phosphorylation of the cellular targets of human atrial cardiomyocytes

where  $CaMK_{bound}$  and  $CaMK_{trap}$  are the fractions of CaMKII binding sites bound to  $Ca^{2+}/CMDN$  and with trapped CMDN respectively and are given as:

$$CaM K_{bound} = CaMK_0 \frac{1 - CaMK_{trap}}{1 + \frac{K_{m,CaM}}{[Ca^{2+}]_{ss}}}.$$

Here CaMK<sub>0</sub> is the fraction of active CaMKII binding sites at equilibrium,  $K_{m,CaM}$  is the sensitivity factor of CaMKII bound to Ca<sup>2+</sup>/CMDN. Hence, CaMK<sub>bound</sub> quantifies the dependence on  $[Ca^{2+}]_{ss}$ . CaMK<sub>trap</sub> is a dynamic quantity and is given as:

$$\frac{dCaMK_{trap}}{dt} = \alpha_{CaMK} CaMK_{bound} (CaMK_{bound} + CaMK_{trap}) - \beta_{CaMK} CaMK_{trap},$$

where  $\alpha_{CaMK}$  and  $\beta_{CaMK}$  are the transition rates. All the values used for CaMKII model constants are listed in Table A5, section B.

The model presented is based upon an ordinary set of differential equations implemented in Matlab using a stiff ordinary differential equation solver method (ode15s). Under the basal conditions, the model was simulated at 1 Hz frequency and 800 beats to reach the steady state. The simplification made in the AP module made it more computationally efficient so that it was now defined by considering 30 state variables in comparison to 43 variables of the parent model. Consequently, this simplified description of intracellular Ca<sup>2+</sup> dynamics reduced the computational load of the AP module by 15% than the parent model for a 10-s simulation and at 1 Hz pacing takes about 1.7 s to run on a desktop PC (AMD Ryzen 7 CPU, 1.90 GHz, 16 GB of RAM). All the modifications introduced for the formulation of the model are shown in Appendix 1.

## Results

## AP, Ca<sup>2+</sup> transient and active force characteristics and biomarkers

The simulated steady state characteristics of the AP waveform, Ca<sup>2+</sup> transient in ss and bc compartments, and  $F_{\text{active}}$  are shown in Fig. 4. The modifications with respect to the parent model, KM2011, resulted in a type-3 human atria AP morphology, i.e. triangular AP with no dome (Dawodu et al., 1996) (Fig. 4A). The model shows a prominent plateau phase with respect to the parent model that was achieved by the slow inactivation of  $I_{CaL}$ . The role of the modified formulation of  $I_{CaL}$  (as explained in the Methods section) in producing the AP shape of MBS2023 is assessed in Fig. 5 where the modified  $I_{CaL}$ led to a prominent lengthening of the plateau, early and late repolarization phase of the AP (dashed line in red vs. dotted-dashed line in purple). The CaT produced by the model is presented in Fig. 4B. As expected, the amplitude is greater in the ss than in the bc compartment (0.678 vs. 0.22  $\mu$ M), with similar diastolic levels (0.179 vs. 0.187  $\mu$ M). A temporal heterogeneity is also evident in the dynamics of the transient; moving to the centre, CaT is slowed and delayed, as quantified by the differences in biomarker values between ss and bc:  $\delta ttp_{Ca}$  61 ms,  $\delta rt_{50Ca}$  100 ms, and  $\delta TT_{Ca}$  256 ms. The mean CaT (that is the weighted average between cytosolic compartments, to be compared with experimental measurements) was computed and its time course was substantially equal to that of the Ca<sup>2+</sup> transient in the bulk cytosol (not shown). The Ca<sup>2+</sup> handling behaviour for MBS2023 can be analysed using the  $[Ca^{2+}]_{SR}$  traces from ss and bc compartments in Fig. 4D. In the ss compartment, a prominently large Ca<sup>2+</sup> release can be seen from the SR into the ss compartment, with fast release flux dynamics (Fig. A4D vs. H, red continuous line). Conversely, in the bulk cytosol, the Ca<sup>2+</sup> release is smaller (Fig. A4G, red continuous line), slower, and slightly delayed, resulting in more  $Ca^{2+}$  SR load. The force developed ( $F_{active}$ ) in the bulk cytosol because of the coupling is shown in Fig. 4E.

A quantitative comparison of AP, CaT and  $F_{active}$  biomarkers with *in-vitro* human atrial data ranges is presented in Fig. 4 and Table 2. The simulated biomarkers are plotted using biomarker-specific markers and colour bars representing their *in-vitro* variability ranges. The dashed lines represent the *in-vitro* biomarker mean values, also reported in Table 2. Most of the biomarkers are within the range of the experimental ones. In particular, the APs simulated by MBS2023 (in blue) and KM2011 (in red) are compared in Fig. 4A, together with their APD<sub>50</sub> and APD<sub>90</sub> biomarkers. Our AP model has a more dominant plateau than KM2011 and is more in agreement with the mean APD<sub>50</sub> reported by the *in-vitro* data range. For late repolarization, APD<sub>90</sub> seems comparable for both models and lies on the mean value of the data. The MBS2023

CaT upstroke is faster compared to KM2011 (Fig. 4*C*), while its decay is slower (within the experimental range, Fig. 4*C*). The dynamics of  $F_{active}$  have been compared with the mean CaT of each model along with their ttp and rt50 biomarkers (Fig. 4*F*). The CaT arises earlier and decays more slowly than  $F_{active}$ , as also shown by the light and dark data range colours, respectively (Fig. 4*F*). The rising part of the CaT and  $F_{active}$  curves are compared using ttp, where a ttp<sub>Ca</sub> smaller than ttp<sub>Factive</sub> can be observed. The ttp<sub>Ca</sub> is slightly longer than the data range but ttp<sub>Factive</sub> is within the lower data bound (Fig. 4*F*). The comparison of the relaxation phase was carried out in Fig. 4*F* inset, where both curve peaks were aligned to highlight that rt<sub>50Ca</sub> is longer than rt<sub>50Factive</sub>. Moreover, rt<sub>50Factive</sub> lies quite close to the lower bound of the data range.

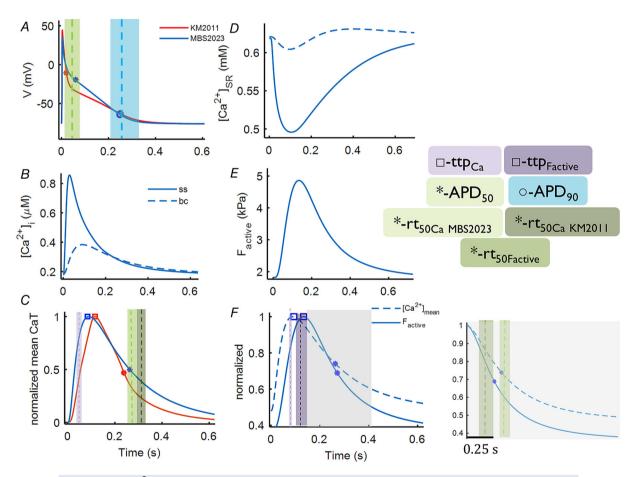


Figure 4. AP,  $Ca^{2+}$  handling and  $F_{active}$  characteristics obtained with the MBS2023 model compared to those obtained with the KM2011 model and experimental human atrial data A few selected biomarkers are also plotted using various marker labels like ttp<sub>Ca</sub> or ttp<sub>Factive</sub> with open squares,

APD<sub>50</sub> or rt<sub>50Ca</sub> or rt<sub>50Factive</sub> with asterisks, and APD<sub>90</sub> with open circles. The colour for each shaded area corresponds to the data range of a biomarker as also reported in Table 2. *A*, comparison of AP traces for MBS2023 (in blue) with KM2011 (in red) and with experimental APD<sub>50</sub> and APD<sub>90</sub>. *B*, Ca<sup>2+</sup> transients for subspace (continuous), and bulk cytosol (dashed line). *C*, comparison of normalized mean (i.e. weighted average between cytosolic compartments) Ca<sup>2+</sup> transient from MSB2022 (in blue) and KM2011 (in red) and with experimental ttp<sub>Ca</sub> and rt<sub>50Ca</sub>. *D*, Ca<sup>2+</sup> concentration inside the SR for both ss and bc subcompartments. *E*, *F*<sub>active</sub> produced by the EM coupling in MBS2023 model. *F*, comparison of normalized mean Ca<sup>2+</sup> transient and *F*<sub>active</sub> time course in the MBS2023 model (inset: alignment of the same curve with respect to their time to peaks for the sake of comparing the relaxation dynamics of the Ca<sup>2+</sup> transient (dashed line) and *F*<sub>active</sub> (continuous line)).

Action poten	tial		Ca <sup>2+</sup> transien	t		Active force		
Biomarker	Mean value range [LB, UB]	Model output	Biomarker	Mean value range [LB, UB]	Model output	Biomarker	Mean value Range [LB, UB]	Model output
APD <sub>90</sub> (ms)	263.05 [202, 332]	248	Са <sub>атр</sub> (µМ)	0.268 [0.18, 0.4]	0.221	F <sub>active</sub> (kPa)	5.92 [3.14, 9.5]	4.7
APD <sub>50</sub> (ms)	50.02 [25, 94.14]	60.2	Ca <sub>dias</sub> (µM)	0.223 [0.2, 0.25]	0.187	ttp <sub>Factive</sub> (ms)	104.98 [79.5, 161]	135.7
APD <sub>30</sub> (ms)	7.725 [5, 13.9]	13.4	ttp <sub>Ca</sub> (ms)	52.5 [49.4, 55.6]	93.7	rt <sub>50Factive</sub> (ms)	80.92 [60.2, 118.6]	140
V <sub>amp</sub> (mV)	103.266 [83, 130]	111.7	rt <sub>50Ca</sub> (ms)	177.5 [168.5, 186.5]	176	rt <sub>90Factive</sub> (ms)	200 [153, 235.9]	406
V <sub>diast</sub> (mV)	-75.1 [-68, -75.1]	-75.7	TT <sub>Ca</sub> (ms)	539.1 [508.1, 570.1]	690	TT <sub>Factive</sub> (ms)	433.3 [413.1, 453.5]	663
d <i>V</i> /dt <sub>max</sub> (mV/ms)	179.23 [159, 231.9]	177	au (s <sup>-1</sup> )	335 [220, 450]	288			

Table 2. Comparison of AP,  $Ca^{2+}$  transient, and  $F_{active}$  biomarkers obtained with the model and the values (lower and upper bounds along with their mean values) from *in-vitro* human atrial data extracted from references in Table A2 and A4

# Effect of CaMKII inhibition on MBS2023 model characteristics

All the major currents for the MBS2023 model at the basal frequency of 1 Hz are shown in Fig. A3 (red continuous line). The effect of CaMKII inhibition (blue dashed line) on the AP is compared with the basal condition of the model where shortening of the mid-repolarization phase is observed, as shown in Fig. A3A. Two of the

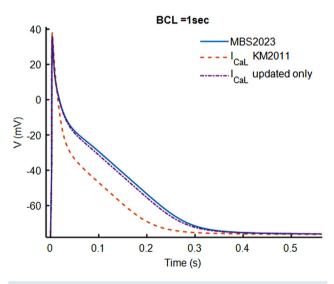


Figure 5. Role of membrane current  $I_{CaL}$  in producing action potential for MBS2023 model (blue continuous line) at 1 Hz frequency

Steady-state action potential (AP) for basal model (blue continuous line) compared with AP using KM2011 formulation of  $I_{CaL}$  (red dashed line) and AP using updated  $I_{CaL}$  current only (in purple dotted-dashed line). [Colour figure can be viewed at wileyonlinelibrary.com]

AP biomarkers, i.e.  $dV/dt_{max}$  and  $V_{amp}$  are compared with the experimental data (Leberk et al., 2018) where a qualitative equivalence can be observed. Among the ionic currents, I<sub>Na</sub> shows a increase (Fig. A3B), and the IK current shows slightly fast inactivation (Fig. A3F and G). The  $Ca^{2+}$ -transient and all the incoming and outgoing fluxes are shown in Fig. A4 (red continuous line) along with the effect of CaMKII inhibition (blue dashed line). The CaT in ss (Fig. A4A) rises instantly and decay is also fast in comparison to CaT in bc (Fig. A4E). This is evident with fast release and uptake flux dynamics of the ss compartment (Fig. A4D and C). The CaT propagates towards the bc compartment via diffusion  $(J_{\text{diff}})$  (Fig. A4B). In the bc compartment, the release flux is quite small whereas the uptake is very large in magnitude (Fig. A4H and G) in comparison to ss. CaMKII inhibition has reduced the  $J_{up}$  flux and has reduced  $F_{active}$  (Fig. A4F), as was observed by Lebek et al. (2018).

### Mechano-calcium feedback effects

The model characteristics were analysed by keeping SL fixed, in an iso-sarcometric condition as usual in *in-vitro* experiments (referred to in Table A2). The consequence of coupling on AP, CaT and  $F_{active}$  for the novel MBS2023 model was analysed by comparing the model output under basal conditions, i.e. bidirectional coupling (continuous line), with noMCF condition (dashed lines), i.e. by removing the dynamic effect of the Ca<sup>2+</sup> bound to troponin (d[Ca<sup>2+</sup>]<sub>TRPN</sub>/dt as shown in eqn (1)). The MBS2023 AP waveform is compared with AP under the noMCF condition in Fig. 6A. The AP curve and its kinetics are only slightly affected by MCF and this is demonstrated

by the early repolarization phase. Accordingly,  $I_{CaL}$ (Fig 6B) shows more inactivation during the plateau of the curve under the noMCF condition. Similarly, under the noMCF condition, we observed an enhanced systolic level of  $Ca^{2+}$  in the bulk cytosol (Fig. 6D) due to the elevated amount of  $Ca^{2+}$  bound to TRPN (Fig. 6F), which is not sequestered from the cytoplasm. Moreover, in comparison to the basal state of the model, the CaT peak arrives earlier, and the decay is faster when the MCF effect is switched off. In particular, under the basal state the  $rt_{50Ca}$  for the model was within the data range; however, it went completely out of range under the noMCF condition. Although the model in the basal state produces a  $TT_{Ca}$  that is longer than the experimental range, under noMCF the TT<sub>Ca</sub> is far too short, showing a large difference from the experimental range (reported in Table 2). Because of raised Ca<sup>2+</sup> levels in the bc under the noMCF condition, the exchanger I<sub>NaCa</sub> shows enhanced inward current followed by a reduction in the latter phase of time, as shown in Fig. 6C. Similarly, the  $F_{\text{active}}$  curve with fast relaxation dynamics was also observed under the noMCF condition, as shown in Fig. 6E. Overall, the kinetics of CaT is fully dependent on the dynamics of MCF, i.e. the concentration of  $Ca^{2+}$ bound to TRPN.

### Quantifying SR Ca<sup>2+</sup> activity

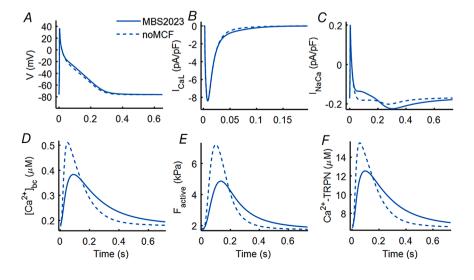
SR Ca<sup>2+</sup> content is a major modulator of CaT. Therefore, we simulated the caffeine effect to quantify the SERCA activity by measuring SR Ca<sup>2+</sup> content and its pumping rate. In experiments, atrial myocytes are preconditioned with an  $I_{CaL}$  activation voltage clamp protocol for 1 min. Caffeine (10 mM) causes a large rise in CaT, as a result of SR Ca<sup>2+</sup> release. The slow decay of CaT is mainly through the  $I_{NaCa}$  exchanger, and the integral of the current is a measure of SR Ca<sup>2+</sup> content. Following the experiment,  $Ca^{2+}$  current and transient were simulated (Fig. 7A) following the voltage clamp protocol presented in Grandi et al. (2011) and Voigt et al. (2014) at 0.5 Hz, as shown in Fig. 7A (inset). In detail, the protocol starts with a 100 ms ramp from -80 mV (holding potential) to -40 mV to inactivate  $I_{Na}$ , followed by a step pulse to +10 mV for 100 ms (Grandi et al., 2011; Voigt et al., 2014). The resultant  $CaT_{bc}$  decay phase was fitted on a mono-exponential equation (Voigt et al., 2014 Supplementary information).

The voltage clamp protocol resulted in  $CaT_{svs} = 0.41 \ \mu M$  (Fig. 7C, middle),  $Ca_{diast} = 0.179 \ \mu M$ (Fig. 7C, left), and decay time constant of 288 ms (Fig. 7D), corresponding to a decay rate  $k_{sys} = 3.47 \text{ s}^{-1}$ . In comparison to the experiments in Fig. 7C, the CaT amplitude matches well with data from Fakuade et al. (2021) and Voigt et al. (2012, 2014), and it is slightly smaller than the value reported by Heijman et al. (2020). The decay rate for CaT (Fig. 7D) is in between the reported data, i.e. smaller than Fakuade et al. and Heijman et al. but greater than Voigt et al.

The caffeine induced Ca<sup>2+</sup> transient (cCaT) experiment was replicated by switching the membrane voltage to holding potential (-80 mV) and running the simulation in the RyRs open state, i.e. both the activation, inactivation and Ca<sup>2+</sup> SR-dependent activation gates were set to 1, to allow for the maximum release flux to the cytosol. The cCaT is shown in Fig. 7*B* (bottom panel) where the cCaT has a slow decay phase with high amplitude (2.6 times greater than CaT, Fig. 7*E*). The cCaT amplitude is quite equivalent to Fakuade et al.'s data and is smaller than Voigt et al.'s data. The resulting inward current flowing through  $I_{NaCa}$  is shown in Fig. 7*B* (top panel), while the integral of  $I_{NaCa}$ , representing the measure of SR-Ca<sup>2+</sup> content, is shown in Fig. 7*F*. The model has a SR Ca<sup>2+</sup> load equivalent to Voigt et al.'s data but is smaller than the data bar plots

#### Figure 6. Effect of mechano-calcium feedback on MBS2023

Consequences of coupling on MBS2023 AP (continuous line) in comparison to noMCF (dashed line) condition (*A*), ionic currents  $I_{CaL}$  and  $I_{NaCa}$  plots with respect to noMCF condition (*B* and C), Ca<sup>2+</sup> transient in bulk cytosol with respect to noMCF condition (*D*), comparison of  $F_{active}$  time course with noMCF condition (*E*), and concentration of Ca<sup>2+</sup> bound to TRPN with and without MCF (*F*). [Colour figure can be viewed at wileyonlinelibrary.com]



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from Heijman and Fakuade et al. The SERCA pumping rate ( $k_{\text{SERCA}}$ ) is estimated as:

$$k_{\text{SERCA}} = k_{\text{sys}} - k_{\text{caff}}$$

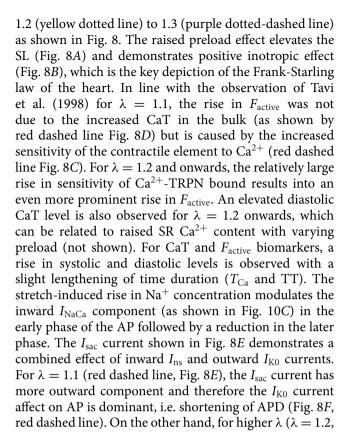
where  $k_{\text{caff}}$  is given as:

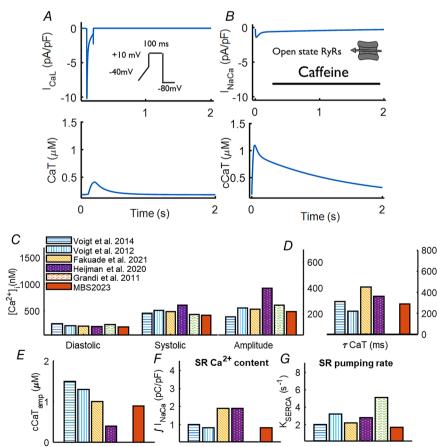
$$k_{\rm caff} = k_{\rm NaCa} + k_{\rm PCa},$$

where  $k_{PCa}$  is a small contribution of sarcolemma Ca<sup>+</sup> current. Using monoexponential equation with cCaT<sub>amP</sub> = 0.9  $\mu$ M, and cCaT<sub>dias</sub> = 0.163  $\mu$ M we obtained  $k_{caff} = 1.8 \text{ s}^{-1}$  yielding  $k_{SERCA} = 1.9 \text{ s}^{-1}$  (Fig. 7*G*). Thus, the comparison shows that the model has a SR Ca<sup>2+</sup> content in between the experimental data ranges whereas the SERCA pumping time is slower than most of the experimental values (Fig. 7*F*) but not too far from Voigt et al. and Fakuade et al.

### Effect of stretch variation on AP and contractility

Using the MBS2023 model, we simulated the effect of stretch variations on AP, CaT and  $F_{\text{active}}$  characteristics under isometric conditions, where for a given SL the preload effect was varied using  $\lambda$ . We simulated the  $I_{\text{sac}}$  current, which is the combination of  $I_{\text{ns}}$  and  $I_{\text{K0}}$ , and varied  $\lambda$  from 1 (blue continuous line), 1.1 (red dashed line),





#### Figure 7. Validation of SR Ca<sup>2+</sup> content and pumping rate in the MBS2023 model

A, the voltage clamp protocol (inset) resulting in I<sub>CaL</sub> current (top) and I<sub>CaL</sub> triggered Ca<sup>2+</sup> transient (CaT) (bottom). B, the caffeine induced inward INACa current (top) with CaT (cCaT) (bottom). The experimental data from bar plots extracted from Voigt et al. (2014) (blue horizontal lines), Voigt et al. (2012) (blue vertical lines), Fakuade et al. (2021) (yellow diagonal lines), Heijman et al. (2020) (purple dots), Grandi et al. (2011) (green curvy lines), and MBS2023 output (filled red bars). C shows the characteristics of I<sub>CaL</sub> triggered CaT diastolic (left), systolic (middle), and amplitude (right). D shows the decay time constant for the CaT. E, the caffeine induced CaT amplitude. F and G, SR Ca<sup>2+</sup> content is the integral of INaCa current under caffeine application and the  $k_{\text{SERCA}}$  is the measure of SERCA pump rate and is compared for the simulated output with the experimental data in panels F and G, respectively. [Colour figure can be viewed at wileyonlinelibrary.com]

This conduction slowd

yellow dotted line, and  $\lambda = 1.3$ , purple dotted-dashed line) the effect of  $I_{\rm ns}$  predominates, as can be seen by the inward current and slow kinetics of  $I_{\rm sac}$  current (Fig. 8*E*). As expected, for large  $\lambda$  ( $\lambda = 1.2, 1.3$ ), the depolarization of RMP is more pronounced.

The reversal potential of  $I_{ns}$  lies near 0 mV (as shown in Fig. A7, red triangles) whereas the  $I_{K0}$  current has a potential quite close to the RMP. This can also be depicted from the time course for each current shown in Fig. A9E. Consequently,  $I_{ns}$  depolarizes the RMP of the AP with lengthening of late repolarization  $(APD_{90})$ and no effect on the mid-repolarization phase  $(APD_{50})$ (Fig. A9F on left). In contrast,  $I_{K0}$  has a slight shortening effect on the late repolarization phase of the AP (Fig. A9F on right). The combined effect of  $I_{sac}$  on AP is shown in Fig. 8F, where the shortening effect is predominant for smaller preload ( $\lambda = 1.1$ ), and the depolarized RMP is observed for higher  $\lambda$ . Consequently, the depolarized RMP causes a 30% slowing of  $dV/dt_{max}$  at  $\lambda = 1.2$  and will be discussed in detail in the Discussion section. The depolarization of RMP gets strengthened by increased conductance of  $g_{nsmax}$ , as shown in Fig. A10F where the basal state of the model (in yellow) at  $\lambda = 1.2$  demonstrates another 41% slowing of  $dV/dt_{max}$  for a 50% rise in  $g_{nsmax}$ 

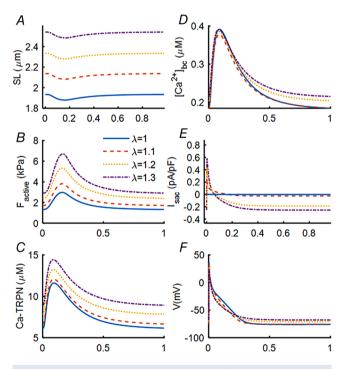


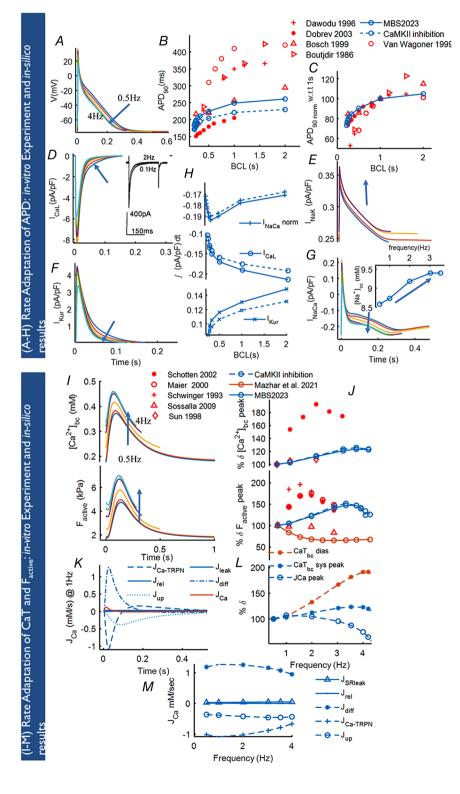
Figure 8. MBS2023 model characteristics for stretch variations  $\lambda$  from 1 (continuous blue line), 1.1 (red dashed line), 1.2 (yellow dotted line), and 1.3 (purple dotted-dashed line) *A*, the raised level of contractile element length (SL). *B*, active contractile force ( $F_{active}$ ). *C*, the concentration of Ca<sup>2+</sup> bound to troponin (TRPN). *D*, Ca<sup>2+</sup> transient in bulk cytoplasm (bc). *E*, stretch activated current  $I_{sac}$ , which is the sum of  $I_{ns}$  and  $I_{KO}$ , *F*, the action potential (AP). [Colour figure can be viewed at wileyonlinelibrary.com]

(in red). This conduction slowdown is usually found to be a consequence of acute atrial dilatation that can be considered vulnerable to AF (Ravelli et al., 2011).

#### Rate adaptation of AP and contractility

An increase in pacing rate causes a physiological shortening of APD of atrial myocytes as quantified in various experimental studies (referred to in Table A4). We simulated the rate dependent behaviour of our model by running the model for a physiological range of BCLs (2, 1, 0.5, 0.33 and 0.283 s). AP curve rate dependence is shown in Fig. 9A. In addition, the APD<sub>90</sub> shortening trend with BCL for a sparse set of *in-vitro* data is shown in Fig. 9B (continuous line). The relative change in APD<sub>90</sub> for BCLs [2, 0.3] is -20% for the MBS2023 model and that gets increased to -28.8% under CaMKII inhibition (dashed line). Hence, the monotonically increasing activity of CaMKII<sub>active</sub> for the physiological range of frequency (0.5-3 Hz) plays its part in steep APD<sub>90</sub> rate dependent shortening. Given the sparsity of the experimental data, we compared the data with the APD<sub>90</sub> output normalized with respect to BCL = 1 s, as shown in Fig. 9C where the model output is quite within the range of the data spread.

The mechanism behind the APD<sub>90</sub> shortening with rate is explained with the help of the ionic currents, which show substantial rate dependence, as shown in Fig. 9D-H: the inactivation of  $I_{CaL}$  reduces the plateau of the current (Fig. 9D) and has a time dependent recovery property that results in significant rate dependence and hence contributes to APD<sub>90</sub> shortening which is pronounced only for longer BCLs (2-0.5 s) (Li et al., 1997). A parallel phenomenon of rate dependent Na<sup>+</sup> accumulation in the cytosol occurs in ventricular CMs (Harrison et al., 1995; Verdonck et al., 2004) and in atrial guinea-pig (Wang et al., 1993). The model shows a similar Na<sup>+</sup> accumulation (Fig. 9G, inset) which causes the enhanced pumping activity of Na<sup>+</sup>/K<sup>+</sup> pump with rate. This increased pumping results in an increased outward current  $I_{\text{NaK}}$ (Fig. 9*E*) that links to APD<sub>90</sub> shortening. Since  $I_{Kur}$  block can contribute to the positive inotropic effect, the rate dependent trend of  $I_{Kur}$  integral is analysed in Fig. 9H (lower panel), which showed a rate dependent reduction mainly because of the later decay phase of the current (Fig. 9F). Thus, reducing outward  $I_{Kur}$  is contributing to the accumulation of  $Ca^{2+}$  in the cytosol (Fig. 9J). During fast pacing, a combination of both  $Na^+$  and  $Ca^{2+}$  (Fig. 9G and 91) increases in the cytosol contributes to making the  $I_{\text{NaCa}}$  exchanger work more in reverse mode. In the model output, this is shown by the integral of I<sub>NaCa</sub> normalized to BCL with the rate (Fig. 9H, top) decreasing or becoming more positive for shorter BCLs (after 330 ms) (Fig. 9G). The CaMKII inhibition effect on the integral of currents (Fig. 9*H*, dashed lines) reduces the  $I_{CaL}$  and  $I_{Kur}$  currents for longer BCLs. For the  $I_{NaCa}$  integral, CaMKII inhibition shows a slight reduction in the overall current. Hence, a combined increase in outward  $I_{NaK}$  and  $I_{NaCa}$ , and a reduction in inward  $I_{CaL}$  currents is the major contributor of APD<sub>90</sub> rate dependent shortening for the physiological range of BCLs. The rate dependent behaviour of CaT in bulk cytosol and the  $F_{\text{active}}$  for our model is shown in Fig. 9*I* -*M*. The positive inotropic effect with a parallel rise in CaT with rate is observed in the *in-silico* model, as shown in Fig. 9*I*. Likewise, CaT time course shows a rise in both systolic and diastolic levels. The biphasic



#### Figure 9. Rate adaptation response of MBS2023 and its dependence on CaMKII activity

A–H, rate dependence of AP and the cellular mechanism behind it. A, in-silico results of AP time course for each rate. B, APD<sub>90</sub> shortening trend for the model (continuous line) with BCL compared to experimental data listed in Table A4 and CaMKII inhibition (dashed line). C, APD<sub>90</sub> rate dependence normalized with respect to BCL = 1 s. D,  $I_{CaL}$  rate dependence in experiments (Li et al., 1997) (inset) compared with in-silico results. E-H, the mechanism of APD<sub>90</sub> shortening. E, increasing I<sub>NaK</sub> traces with rate. F, rate dependency of IKur current. G, increasing inward I<sub>NaCa</sub> current traces with inset showing rate dependent Na<sup>+</sup> accumulation in the cytosol. H, rate dependent trend of the integral of I<sub>NaCa</sub> (top) normalized with respect to the BCL, I<sub>CaL</sub> (middle), and I<sub>Kur</sub> (bottom) is shown. I-M, rate dependence of Factive and CaT and its cellular mechanism. I, J, in-silico results of CaT and Factive rate dependence as percentage change of CaT peak (top) and Factive peak values produced by the model in comparison to experimental data (references in Table A3) and CaMKII inhibition effect (dashed lines). K-M, the mechanism of CaT rate dependence. K, the comparison of magnitude and dynamics of inward fluxes to bc, diffusion from ss to bc compartment (J<sub>diff</sub>) (dotted-dashed line), the SR release  $(J_{rel})$  and leak  $(J_{SRleak})$ ; and outward fluxes, SR uptake  $(J_{up})$  (dotted line), J<sub>Ca-TRPN</sub> (dashed line) at 1 Hz frequency. The sum of all fluxes is  $J_{Ca}$  in red. L, rate dependence trend of CaT systolic (asterisks), diastolic (cross) and J<sub>Ca</sub> (open circles). M, rate dependency of peak values of each flux using various markers, J<sub>diff</sub> (asterisk), Jup (open circles), JCa-TRPN (crosses), J<sub>SRleak</sub> (triangles), J<sub>rel</sub> (dotted).

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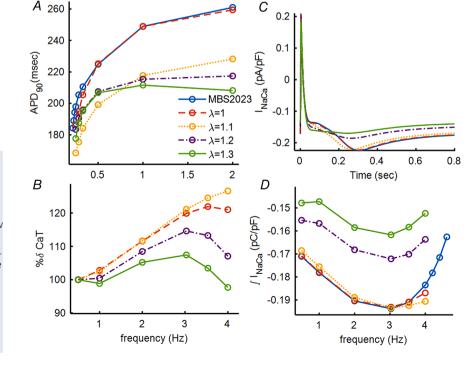
percentage change of CaT peak in the bc compartment is compared with limited available data (Fig. 9J). Similarly, the percentage change of  $F_{\text{active}}$  is compared with human experimental data (references in Table A2) in Fig. 9J (bottom panel). The rate dependent shortening of  $F_{\text{active}}$ biomarkers normalized with respect to 1 Hz is shown in Fig. A2. Among the biomarkers, the late relaxation phase, i.e. rt<sub>90</sub> shows a prominent rate dependency and model output is quite in line with the experimental data, as shown in Fig. A2. In contrast, the rt<sub>50</sub> and ttp show a flat rate dependency. The cellular mechanism behind the rate dependence of CaT is explained in Fig. 9K-M). Following this, all the inward and outward fluxes in the bc compartment are compared (Fig. 9K) under steady state conditions and 1 Hz frequency. Hence, a difference in dynamics and amplitude of the inward fluxes, i.e. diffusion from ss to bc  $(J_{diff})$ , SR release in the bc  $(J_{rel})$ , and the SR leak ( $J_{SRleak}$ ); and outward fluxes, i.e. SR uptake ( $J_{up}$ ) from the bulk,  $Ca^{2+}$  bound to TRPN ( $J_{Ca-TRPN}$ ) can be observed.  $J_{diff}$  shows a large peak comparable with  $J_{Ca-TRPN}$ and followed by  $J_{up}$  whereas, minimum contribution is observed by  $J_{rel}$  and  $J_{SRleak}$ . In addition, flux  $J_{Ca-TRPN}$  is the fastest followed by  $J_{\text{diff}}$  and  $J_{\text{up}}$ . The sum of all the fluxes is shown by  $J_{Ca}$  (in red) in Fig. 9K.  $J_{Ca}$  has fast dynamics and a depressed systolic level. The rate dependence of CaT systolic and diastolic values can be appreciated in Fig. 9I and is quantified, in terms of percentage in Fig. 9L. The progressive increase in diastolic  $[Ca^{2+}]_{bc}$  is the main determinant for a consequent increase in the systolic  $[Ca^{2+}]_{bc}$ ; however, a concurrent progressive decrease with

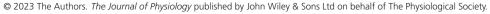
rate of the net Ca<sup>2+</sup> influx into the bulk cytoplasm,  $J_{Ca}$  (peak values) takes place, resulting in the overall biphasic rate dependence of the CaT systolic peaks. A further investigation behind this decreasing rate dependence of  $J_{Ca}$  is carried out by plotting the peak values of each flux with frequency as shown in Fig. 9*M*. The  $J_{diff}$  peak (asterisk) increases slightly at lower rates (up to 1 Hz) and then reduces progressively with the rate. On the other hand, the outward flux,  $J_{up}$  (open circle) shows a slight increase with rate till 4 Hz then a minute decrease, whereas a prominent decrease in  $J_{Ca-TPRN}$  (plus) is observed with rate. Hence, the declining trend of diffusion with a parallel reduced amount of Ca<sup>2+</sup> bound to TRPN at higher rate (after 3.5 Hz) gives rise to the biphasic decline of CaT in the bulk cytosol.

### Effect of mechano-electric feedback on rate adaptation.

We have analysed the effect of stretch variation via  $I_{sac}$  current, under isometric condition, on rate adaptation of APD and CaT where the rate varies from BCL = 2 to 0.283 s (Fig. 10). APD<sub>90</sub> rate dependent shortening for varying  $\lambda$  from1 to 1.3 is shown in Fig. 10*A*. For  $\lambda = 1$ , the rate adaptation of APD<sub>90</sub> (red dashed line) is quite overlapping with the basal state of the model (MBS2023, blue continuous line, Fig. 10*A*). When the stretch is increased, for  $\lambda = 1.1$  (yellow dotted line, Fig. 10*A*), the APD<sub>90</sub> demonstrates an overall shortening with a no sharp change in rate adaptation trend. With another rise in  $\lambda = 1.2$ , the APD<sub>90</sub> rate dependence is flattened first for longer BCLs (from 2 to 1 s) and the

**Figure 10. Stretch variation response on rate adaptation of APD and CaT** The model was compared under four different simulation settings: control condition as MBS2023 (blue continuous line),  $\lambda = 1$  (red dashed line),  $\lambda = 1.1$  (yellow dotted line),  $\lambda = 1.2$  (purple dotted-dashed line), and  $\lambda = 1.3$  (in green continuous line). *A*, the rate adaptation of APD<sub>90</sub>; *B*, the rate dependence of percentage change of CaT systolic peaks; C, the time course of  $I_{NaCa}$ current; *D*, integral of  $I_{NaCa}$  current trend with rate normalized with respect to each BCL. [Colour figure can be viewed at wileyonlinelibrary.com]





trend gets reversed for higher stretch levels ( $\lambda = 1.3$ ). In addition, the slope of APD<sub>90</sub> rate dependence for shorter BCLs (0.5 onwards) was reduced equally for both  $\lambda = 1.2$ and 1.3 (Fig. 10A). This can be a consequence of the exponentially rising  $I_{ns}$  current with  $\lambda$ . Hence, the overall depression of APD<sub>90</sub> magnitude with rate is potentially due to the rise in the  $I_{K0}$  contribution to AP for lower  $\lambda$  $(\lambda = 1.1 \text{ in Fig. 10A})$  and the change in slope of APD is because of the increasing role of  $I_{\rm ns}$  with rising  $\lambda$ . Apart from this, the mechanical stretch increases the inward flow of Na<sup>+</sup> ions in the myocytes and, hence, can affect the operating modes of the Na<sup>+</sup>/Ca<sup>2+</sup> exchanger ( $I_{NaCa}$ ). Accordingly, we analysed the time course of the  $I_{NaCa}$ current for a basal rate of 1 Hz (Fig. 10C) and the rate dependence of integral of I<sub>NaCa</sub> current normalized with respect to its BCL (Fig. 10D) for all the conditions. Under basal conditions, the rate dependent Na<sup>+</sup> accumulation in the cytosol increases the inward flow of current up to BCL = 0.33 s (3 Hz in Fig. 10D) and then there is a shift to outward mode. For varying  $\lambda$ , the integral of  $I_{NaCa}$  trend shows a decrease in overall magnitude of the current with no striking change in the rate adaptation trend.

The rise in CaT systolic level can be related to the enhanced sensitivity of Ca2+-TRPN binding, as was already shown in Fig. 8C. The rate dependent biphasic behaviour of percentage change in CaT<sub>peak</sub> is shown in Fig. 10B (right) where under isometric mode of the model, it shows a rise in  $[Ca^{2+}]_i$  up to 3.5 Hz and then a decline. The rate dependency of percentage change in CaT loses its biphasic trend for  $\lambda = 1.1$ , i.e. it always increases; however, with a rise in  $\lambda$  (1.2 and 1.3), the biphasic relation is restored at higher frequencies with a trade-off of some decrease in current at lower frequency (at 1 Hz). This trend can be related to the area of *I*<sub>NaCa</sub> current normalized with respect to the BCL, as shown in Fig. 10D. The integral of  $I_{NaCa}$  at higher frequencies (from 3 Hz and onwards) showing a transition from the inward to the outward mode of the I<sub>NaCa</sub> current at higher frequencies corresponds to the biphasic relation of CaT. The outward mode of  $I_{NaCa}$ allows the influx of Ca<sup>2+</sup> that via diffusion enters the bc and enhances the SR Ca<sup>2+</sup> storage, which can be shown by the elevated diastolic level of CaT at  $\lambda = 1.2$  and onwards (as shown in Fig. 8D).

## Effect of [Ca<sup>2+</sup>]<sub>o</sub> variations

The evaluation of the Ca<sup>2+</sup> handling mechanism underlying the intracellular Ca<sup>2+</sup> homeostasis of the model was carried out by simulating the elevation of the extracellular Ca<sup>2+</sup> concentration  $[Ca^{2+}]_0$  from 0.9 to 3.2 mM. In line with experiments (Brixius et al., 1997) reported in Fig. 11*A* (left), the model reproduced the elevation of CaT and *F*<sub>active</sub> systolic and diastolic levels, as shown in Fig. 11B (right top and bottom panels, respectively). The percentage change trend of systolic (triangles) and diastolic (open circles) for CaT (in red) and Factive (in blue) is shown in Fig. 11B (bottom). The diastolic rise is equivalent for both CaT and  $F_{\text{active}}$  whereas, the systolic level elevation is more pronounced than CaT for all [Ca<sup>2+</sup>]<sub>o</sub> values. The comparison of the percentage rise in systolic peaks of CaT and  $F_{\text{active}}$  with experimental data is quantified in panel C (left panel). Consistent with experimental evidence, the model shows a significant percentage rise in  $F_{\text{active}}$  and CaT throughout the varying  $[Ca^{2+}]_{0}$  and this increased more at 3.2 mM, where a 172% rise in  $F_{\text{active}}$  and 140% in CaT was achieved. Similarly, we compared the percentage change of kinetics of F<sub>active</sub> using biomarkers ttp<sub>Factive</sub>, rt<sub>50Factive</sub> and TT<sub>Factive</sub> for varying  $[Ca^{2+}]_0$  in Fig. 11C (right panel): there is no substantial difference in the percentage changes of biomarkers for  $F_{\text{active}}$ , as can be seen in the experimental values.

The diastolic rise of CaT with elevated  $[Ca^{2+}]_o$  can be speculated to be the result of more  $Ca^{2+}$  influx via  $Ca^{2+}$ channels in the subspace and less efflux either via SERCA activity or the  $I_{NaCa}$  exchanger current. Since the model lacks T-tubules, inside the bulk cytoplasm, the rise in  $Ca^{2+}$  influx is mainly due to a 175% increase in diffusion flux ( $J_{diff}$ ) and slightly to a release flux ( $J_{relbc}$ ), whereas, the extrusion of  $Ca^{2+}$  is dependent on SERCA uptake. Under  $[Ca^{2+}]_o$  elevation, the model showed a 123% rise in SR  $Ca^{2+}$  content (not shown), thus also resulting in more  $Ca^{2+}$  bound to TRPN. Hence, these changes of intracellular CaT with varying  $[Ca^{2+}]_o$  were accompanied by a rise in diastolic and systolic values of  $F_{active}$ .

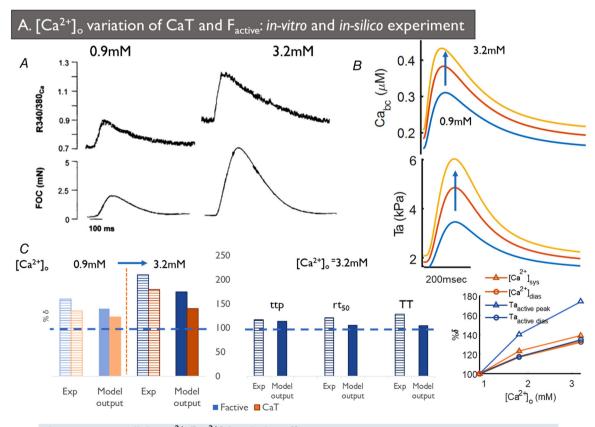
## Enhanced contractility by IKur blocking

The block of  $I_{Kur}$  is expected to prolong the AP plateau that will eventually lead to an increase in contractility. Accordingly, we analysed the blocking effect of  $I_{Kur}$ by setting the current conductance  $g_{Kur}$  to 100%, 80% and 20% of its regular value. In accordance with the experiments (Wettwer et al., 2004), the in-silico model reproduced a prolongation of the plateau phase, as shown in Fig. 12A. Moreover, a change of AP morphology from type 3 to type 4 (Dawodu et al., 1996) for human atria is also observed at a higher percentage (80%) of  $I_{Kur}$  block. The AP biomarker's response to  $I_{Kur}$  blocking is shown in Fig. 12 (right). AP biomarkers included APD at 20% and 30% of repolarization (APD<sub>20</sub>, APD<sub>30</sub>) and the plateau potential (PP) calculated as a mean voltage amplitude within a time window of 20-80 ms after the upstroke. In line with experiments (Wettwer et al., 2004), APD<sub>30</sub> is prolonged from 8 to 105 ms and the plateau potential rises from -16 to +10 mV with the increasing block percentage of  $I_{Kur}$ . In contrast, we observed a lengthening of APD<sub>90</sub> in the model instead of a shortening as reported in the experiments (Wettwer et al., 2004).

The model also demonstrated the secondary effect of  $I_{Kur}$  blockage on AP. As a result, an increased amplitude of  $I_{CaL}$  current was observed at more positive AP potentials (Fig. 12B). Consequently, a large  $I_{CaL}$  activated more outward repolarizing currents like  $I_{\rm Kr}$  (Fig. 12C) and  $I_{\rm Ks}$ (not shown). The reduced activation of Na<sup>+</sup> channels results in a decreasing steady state Na<sup>+</sup> concentration in the cytosol. Similarly, we analysed the prolongation of the APD<sub>90</sub> effect on I<sub>NaCa</sub> current, as shown in Fig. 12D. The decrease in inward  $I_{NaCa}$  current during the early repolarization phase of AP (most dominant in the yellow line) is compensated by the increased current during the later AP phase. In other words, the integral of  $I_{NaCa}$  current (inset) does not show a large change even for a high percentage of IKur block, i.e. 80%. The positive inotropic response of our model is in line with human experimental results (Schotten et al., 2007) (Fig. 12E). The concentration dependent effect of IKur block (Kv1.5) by AVE0118 is modelled using the formulation by Decher et al. (2006). Based on this relation, our model showed good agreement (open circles) with the elevation of contractility caused by the various degrees of  $I_{\text{Kur}}$  block (Fig. 12*F*), as reported in experiments (open squares).

# Simulation of impaired SERCA activity as a mechanism for post-operative AF

The simulation of PoAF associated pro-arrhythmogenic abnormalities can be analysed in the form of beat-to-beat oscillations in CaT and AP, i.e. alternans. The PoAF arrhythmogenic substrate was simulated by the depression of SERCA activity, expression and pumping rate (Fakuade et al., 2021). Following the experimental protocols, we simulated the model from 0.5 to 5 Hz of frequency and analysed the mechanism behind the incidence of CaT and AP alternans under control conditions and the shift in alternans threshold in the PoAF condition. Following the experimental results, the PoAF condition was simulated in the model by reducing the SERCA activity, i.e. the pumping rate to 70% and the expression levels, to 80%.

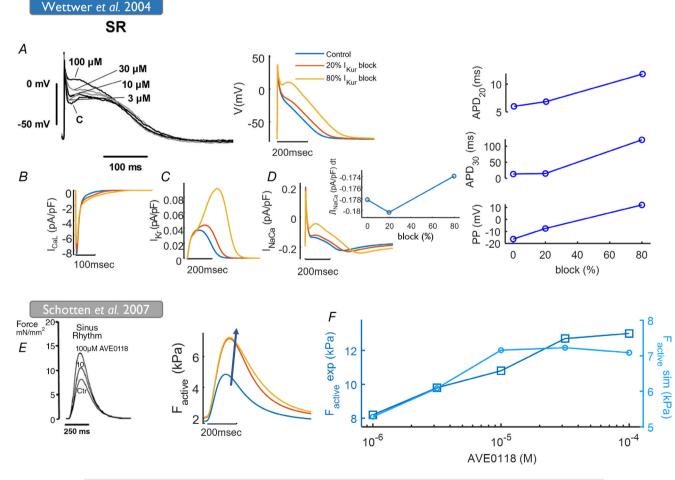


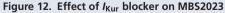
## Figure 11. Extracellular $Ca^{2+}$ ([ $Ca^{2+}$ ]<sub>o</sub>) variation effect on MBS2023

*A*, experimental data (Brixius et al., 1997) for CaT (fura-2 ratio) (top) and  $F_{active}$  elevation (bottom) with  $[Ca^{2+}]_o$  compared with *in-silico* results of CaT (top) and  $F_{active}$  (bottom) time course. *B*,  $F_{active}$  and CaT biomarker response to  $[Ca^{2+}]_o$  variations reproduced with MBS2023 model. *C*, left percentage change of systolic peaks of  $F_{active}$  (red bars) and CaT<sub>bc</sub> (blue bars) compared with the experimental data for varying  $[Ca^{2+}]_o$  from 0.9mM to 3.2mM. *C*, right, percentage change in the kinetics of  $F_{active}$  using biomarkers ttp, rt<sub>50</sub>, and TT with same experimental data.

The onset of AP and CaT alternans at high frequency is observed by the biomarker bifurcation diagram, where the divergence point is the key indicator of alternans initiation, as shown in Fig. 13A and B. The biomarkers for AP are APD<sub>50</sub> and APD<sub>90</sub> (Fig. 13A), and for CaT the amplitude of CaT peaks at consecutive beats in bc (Fig. 13B, left) and ss (Fig. 13B, right) compartments. Hence, under control conditions (100% SERCA activity) (in blue), alternans both in AP and CaT appeared at a frequency of 4.6 Hz. In experiments (Fakuade et al., 2021), the threshold for AP and CaT alternans was shifted down to a lower frequency in the PoAF group. This alternans threshold shifting was demonstrated qualitatively by the model under the PoAF condition (in red). Hence, a higher susceptibility to AP and CaT alternans was found under PoAF, as the alternans threshold was shifted to 4 Hz. Furthermore, under PoAF, a reduced Ca<sup>2+</sup> level is seen in the ss compartment (Fig. 13B, right) with respect to control (red vs. blue) with no change in bc CaT levels (Fig. 13*B*, left). The consecutive three beats, in terms of AP (first panel), CaT in ss (continuous line) and bc (dashed line) (second panel) and  $[Ca^{2+}]_{SR}$  time course (third panel) evoked at 4.4 Hz under control conditions are shown in Fig. 13*C* and *D*. A similar plot of AP, CaT and  $[Ca^{2+}]_{SR}$  under PoAF (Fig. 13*C* and *E*, in red) illustrates the beat-to-beat alternations in CaT systolic and diastolic levels both in ss and bc compartments (second panel) and in  $[Ca^{2+}]_{SR}$  release and end diastolic value (third panel).

The mechanism for the onset of CaT alternans is investigated in Fig. 13F by plotting the three consecutive beats of RyR steady state inactivation gate for the ss compartment at 4.4 Hz for control (top panel, blue traces) and PoAF conditions (bottom panel, red traces). Under control conditions (Fig. 13F, top panel), the peak availability of RyR gate is about 14%, it is uniform all along the consecutive beats, and this gives a stable CaT in the subsequent beats (in light shading). In contrast, in PoAF (bottom panel), the model demonstrates



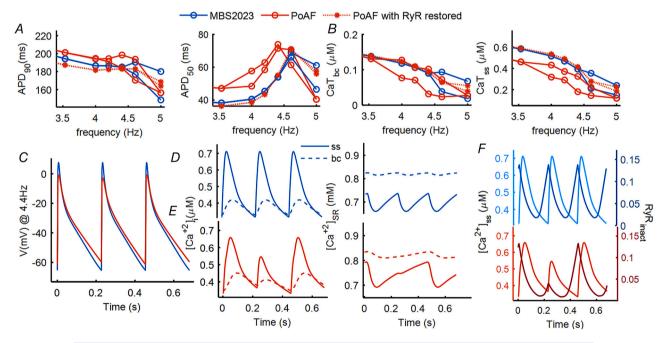


*A*, *in-silico* results of  $I_{Kur}$  blocking on AP time course compared with *in-vitro* experiments (Wettwer et al., 2004). The AP biomarkers were shown, plateau potential (PP), APD<sub>30</sub>, and APD<sub>20</sub> trend. Secondary effect of  $I_{Kur}$  blocking on  $I_{CaL}$  current (*B*),  $I_{Kr}$  current (*C*), and  $I_{NaCa}$  (*D*) with the integral (in the inset). *E* and *F*, the drug concentration dependent inotropic response of the model compared with experimental data (Schotten et al., 2007).

beat-to-beat alternations in RyR inactivation and recovery from inactivation time. The recovery from inactivation is not uniform from beat to beat, and as a consequence the availability of RyR gate alternates between 4% and 14%, resulting in short and long CaTs, respectively. Hence, the refractoriness of RyRs in the ss compartment is one of the mechanism of alternans incidence in the model under PoAF conditions.

Dissecting parameters involved in alternans development. With the help of a dynamic alternans pacing protocol, we have investigated the role of contraction, Ca<sup>2+</sup> handling, and the CaMKII based phosphorylation effect on the model that can be appreciated by APD<sub>90</sub>-ANM and CaT-ANM scores, as shown in Fig. 14, where the ANM threshold is set to 5% and is shown by the red dashed line. Under basal conditions, MBS model has a propensity to alternans incidence at a BCL of 200 ms where Ca alternans are stronger than APD, as shown by the ANM score (7.34% vs. 5.9% in red bars). Among the contraction related parameters, the slowing of transition rates  $K_{\text{off}}$  and  $K_{\text{basic}}$  did not play a significant role in strengthening the alternans magnitude. On the other hand, the reduced sensitivity of thin filament, i.e. increasing  $k_{d-TRPN}$  results in enhancement of alternans more for CaT than APD (ANM scores: 12% vs. 9%) and the alternans completely vanished on decreasing  $k_{d-TRPN}$ . The parameter that has the greatest impact on CaT alternans magnitude is the coupling effect, i.e. MCF feedback as shown in Fig. 14. On reducing the Ca-TRPN binding (by 0.5 times), i.e. a weak MCF effect, the model demonstrated alternans with greater sensitivity to CaT oscillations than APD (ANM score: 13% vs. 10.01%). This major rise in ANM score is a consequence of the enhanced level of free CaT in the cytosol that is achieved because of the reduced concentration of Ca-TRPN bound in the model. Ca-TRPN has a dampening effect on the beat-to-beat oscillations that is even enhanced by increasing the Ca-TRPN binding affinity, as was seen already by reducing the  $k_{d-TRPN}$ parameter.

Among the Ca<sup>2+</sup> handling parameters, the slow inactivation of RyR release gates has slightly reversed the trend, i.e. this parameter is more vulnerable to APD alternans than CaT with ANM score of 9.2% *vs.* 8.78%. The enhanced SERCA activity is favourable for alternans vulnerability, as demonstrated by enhanced SERCA expression and pumping rate. The enhanced CaMKII effect protects the MBS2023 model against the



**Figure 13. Effect of SERCA activity variation as a measure of PoAF condition on MBS2023 model** The SERCA activity was reduced from 100% (control in blue) to the PoAF condition induced by depressed SERCA, i.e. 80% reduced SERCA expression and 70% reduced pumping rate (PoAF in red), and under same PoAF conditions with RyR inactivation kinetics restored as was in KM2011 (in red asterisks). *A*, bifurcation diagram of AP biomarker (APD<sub>50</sub>) (left panel), and APD<sub>90</sub> (right panel). *B*, difference of CaT amplitude on the consecutive beats of CaT in bc (left panel) and ss (right panel) compartments at different frequencies. *C*, consecutive APs at 4.4 Hz under control (blue continuous line) and PoAF (red dashed line) conditions. *D* and *E*, consecutive CaTs in both compartments (left panels) and Ca<sup>2+</sup> concentration in SR,  $[Ca^{2+}]_{SR}$  (right panels) at 4.4 Hz in control (*D*) and PoAF (*E*) conditions. *F*, consecutive beats of RyR steady-state inactivation gate in comparison with CaT from ss compartment at 4.4 Hz under control and PoAF conditions.

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incidence of alternans. The CaMKII expression reduction induces very strong alternans in APD with ANM 12.16% that is even greater than the CaT ANM magnitude (11.05%). Overall, the presence of a coupling MCF effect reduces the alternans induced vulnerability in the model, where the sensitivity of the thin filament is the key modulator of alternans development. In Ca<sup>2+</sup> handling, enhanced SERCA activity and RyR slow inactivation times can play a role in increasing the alternans incidence probability in our model. The role of each parameter discloses some insight into mechanisms behind the alternans vulnerability, as will be explained further in the Discussion section.

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Stretch strengthens the alternans susceptibility. Under isometric conditions, the MBS2023 model demonstrates increased vulnerability to alternans incidence when paced for higher rates, as shown by the bifurcation diagram in Fig. 15. The diagram demonstrates biomarkers for AP, APD<sub>90</sub> (top left), and APD<sub>50</sub> (bottom left) and for CaT amplitude both in bc and ss compartments. The stretch condition at  $\lambda = 1.2$  under  $I_{sac}$  (red triangles) protects the model from both AP and CaT alternans development as was seen under control conditions (in blue). In contrast, the exclusion of outward  $I_{K0}$  from  $I_{sac}$  current resulted in an enhanced vulnerability to alternans development, as shown by the yellow cross in Fig. 10. This shifts the alternans threshold from 4.6 to 3.5 Hz and strengthens the magnitude of oscillations (shown by increased opening of the branches) for both AP and CaT biomarkers. The exclusion of  $I_{\rm K0}$  removes the shortening effect of APD and, hence, susceptibility to alternans incidence.

The mechanism behind the stretch induced alternans incidence at BCL 0.283 s (3.5 Hz) can be seen in Fig. 16 under  $\lambda = 1.2$  with  $I_{sac}$  (continuous blue line) and with  $I_{\rm ns}$  only (red dashed line). In the basal state of the model, one of the mechanisms of alternans incidence was the time constant of the RyR inactivation gate which determines the availability of the RyRs. Here, the alternans appeared at 3.5 Hz (instead of 4.6 Hz as in the basal model), hence resulting into more availability (18%) of RyRs for  $\lambda = 1$  (Fig. 16G). On increasing the preload  $\lambda = 1.2$ , the recovery from inactivation is not uniform from beat to beat, and consequently the availability of RyR gate alternates between 24% and 11%, resulting in short and long CaTs, respectively (Fig. 16E). Other than this, the stretch induced rise in alternans magnitude can be associated with Ca<sup>2+</sup>-TRPN concentration (Fig. 16D). As shown before,  $I_{ns}$  current increases the binding affinity of  $Ca^{2+}$ -TRPN which gets amplified by the  $I_{ns}$  only condition

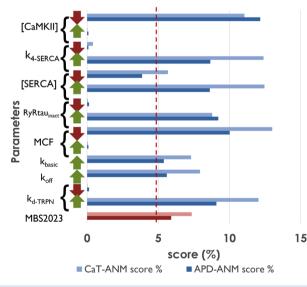


Figure 14. Effect of parameter variations on alternans normalized magnitude (ANM) for CaT<sub>peak</sub> (CaT-ANM) (in light shading) and APD<sub>90</sub> (APD-ANM) (in dark shading) All the parameter values were varied by increasing 2-folds and reducing 0.5-fold. Parameters related to Ca<sup>2+</sup> handling: RyR inactivation time (RyR<sub>tauinact</sub>), SERCA activity, i.e. expression level ([SERCA]), pumping rate ( $k_{4-SERCA}$ ); the feedback: mechano-calcium feedback (MCF) induced by Ca<sup>2+</sup>-TRPN binding; Contraction: thin filament transition rate ( $K_{off}$ ), sensitivity ( $k_{d-TRPN}$ ), and cross-bridge cycling (XBcy) rate ( $k_{basic}$ ); CaMKII: expression ([CaMKII]). [Colour figure can be viewed at wileyonlinelibrary.com]

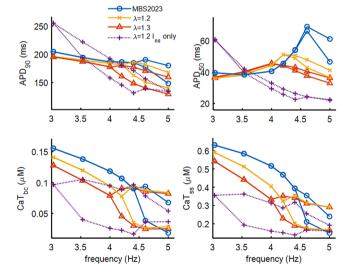


Figure 15. Bifurcation diagram illustrating the comparison of the basal condition of the model MBS2023 (in blue), with conditions under  $\lambda = 1.2 I_{sac}$  current (yellow crosses),  $\lambda = 1.3$  (red triangles), and  $\lambda = 1.2$  for  $I_{ns}$  only (purple crosses) The divergence of the bifurcation branch is a key indicator of alternans incidence in AP biomarkers, APD<sub>90</sub> (top left), APD<sub>50</sub> (bottom left), and CaT amplitude in bulk cytosol (bc) (top right) and the subspace (ss) compartments (bottom right). The alternans magnitude is increasing with increased level of stretch. [Colour figure can be viewed at wileyonlinelibrary.com]

(Fig. 16*B* and *C*). For the  $\lambda = 1.2$  I<sub>ns</sub> only condition, for the long beat (Fig. 16*B*, red dashed line), the raised diastolic level of Ca<sup>2+</sup>-TRPN concentration at the end of the BCL indicates that not all the Ca<sup>2+</sup> is released from the TRPN buffer, showing that the sarcomere was not fully relaxed during the diastole. This increased activation of TRPN buffer during the subsequent short beat increases the fraction of bound Ca<sup>2+</sup>-TRPN, which strengthens the cooperative activation of neighbouring regulatory units resulting in a stronger binding affinity. Thus, an increased Ca<sup>2+</sup>-TRPN fraction will release a small CaT, resulting in a relative beat-to-beat oscillatory mechanism.

### Discussion

We have developed a novel mathematical model that integrates the MER loop with feedforward, the EM coupling and the feedback pathways, via the MCF and MEF. The model couples one AP module, the modified KM2011 (Koivumäki et al., 2011), with a contractility module, RDQ2020 (Regazzoni et al., 2020). Despite the simplification induced in the AP module, the model is capable of preserving the characteristics of centripetal  $Ca^{2+}$  diffusion from the membrane to the bulk.

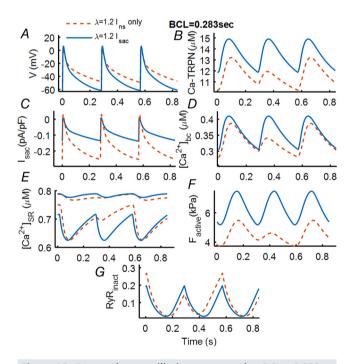


Figure 16. Beat-to-beat oscillations appeared at BCL = 0.238 s under no preload  $\lambda$  = 1 (continuous blue lines) and with preload  $\lambda$  = 1.2 (dashed red lines) conditions

*A*, action potential (AP); *B*, non-selective cation stretch activated current ( $I_{ns}$ ); *C*, SR Ca<sup>2+</sup> content; *D*, amount of Ca<sup>2+</sup> bound to troponin (TRPN); *E*, Ca<sup>2+</sup> transient in subspace ss (large amplitude) *vs*. in bulk cytosol (bc) (small amplitude); *F*, active force ( $F_{active}$ ); *G*, the steady state inactivation gating variable of RyR<sub>ss</sub>. [Colour figure can be viewed at wileyonlinelibrary.com]

# AP and CaT characteristics of the model and the *in-vitro* data

The modifications introduced into the KM2011 model resulted in a type-3 AP morphology of human atria, which has a more prominent plateau phase than the parent model. In other words, we achieved an APD<sub>50</sub> that is more in agreement with the mean human *in-vitro* data; nevertheless, a high level of AP shape heterogeneity has been reported in human atria (Colman et al., 2013; Feng et al., 1998) as reflected in our data collection in Table 2. Of note, the APD<sub>50</sub> obtained from the model is in close agreement with the value for type-3 AP morphology reported by Dawodu (Dawodu et al., 1996). Likewise, the model was able to reproduce faster systolic  $Ca^{2+}$  rise and longer Ca<sup>2+</sup> decay (in the bc compartment) than those of the contraction force, hence having a longer time course than the twitch time (Fig. 4F (inset)), in agreement with the observation made by Brixius (Brixius et al., 1997). Moreover, a quantitative analysis of this phenomenon based on the biomarkers ttp and  $rt_{50}$  for  $F_{active}$  and and  $rt_{50}$ for CaT<sub>bc</sub> resulted in agreement with their experimental data ranges (Table 2). In particular, the CaT produced by the model has slower dynamics than the parent model (Fig. 4C) and than the defined experimental data range. On dissecting the decay phases, the model produced a CaT characterized by an early decay phase  $(rt_{50Ca})$  within the physiological data range (Table 2 and Fig. 4) followed by a slower later decay phase (TT<sub>Ca</sub> in Table 2). The slow dynamics of decay in the later phase usually depend more on the kinetics of Ca<sup>2+</sup> removal from the cytosol than the buffering effect (Smith et al., 2019). In our model, the efflux of Ca<sup>2+</sup> from the bc compartment is carried out by SR Ca<sup>2+</sup> uptake since the exchanger and sarcolemmal pump are localized in the ss compartment. The uptake flux  $(J_{up})$  dynamics in bc compartment was quite slow and small in magnitude (Fig. 8L). This was also evident when comparing the SR pumping rate ( $K_{\text{SERCA}}$ ) with experimental data, as shown in Fig. 7G. This demonstrates a need for for further investigation of the  $Ca^{2+}$  efflux mechanism in the bc of the developed model.

### Atrial stretch enhances the vulnerability to arrhythmias

Experimental evidence indicates that  $I_{ns}$  current plays a role in promoting arrhythmias in intact human atrium (Bode et al., 2000), in rat atrial tissue (Tavi et al., 1996), and in human atrial myocytes (Kamkin et al., 2003). Consistently, at basal frequency of 1 Hz, the rise in stretch (for  $I_{ns}$  only) lengthens the AP with a depolarized RMP in our model (Fig. A9, left panels). In the rat atria, a similar lengthening of APD was demonstrated under the moderate level of stretch applied during the AP diastole (Tavi et al., 1998). In contrast, a

shortening of APD was also reported in guinea-pig atria, although the mechanical stimulus was applied during the mid-repolarization phase (Nazir & Lab, 1996). In experiments, the electrophysiological response to stretch is dependent on timing of the stretch with respect to the AP dynamics. In our model the reversal potential of the  $I_{\rm ns}$  current is near 0 mV and hence the model showed a lengthening effect. The depolarization of RMP resulted in decreased membrane upstroke velocity  $(dV/dt_{max})$  that was slowed down by 30% for  $\lambda = 1.2$ . The sensitivity analysis of  $g_{ns}$  current shows an additional 41% slowdown of  $dV/dt_{max}$  for a 50% rise in  $g_{ns}$  (red line, Fig. A10). The reduction of cell excitability is associated with the macroscopic mechanism of reduced conduction velocity (CV) observed in acutely dilated human atrium (Ravelli et al., 2011). Accordingly, to have a better connection between micro-scale stretch ratio variation effect and macro-scale cell excitability, we simulated a 1D-discrete tissue level experiment for varying  $\lambda$  as shown in Fig. A6. For a 20% stretch rise ( $\lambda = 1.2$ ), the 30% percentage change in  $dV/dt_{max}$  in an isolated cell (red asterisk) can be translated into 10% slowing of CV in the 1D-tissue (blue asterisk, Fig. A6). These results qualitatively agree with the findings of Ravelli et al., where for a 23% increase in atrial volume a 17% slowing of CV was found (Ravelli et al., 2011).

At basal rate, the varying length elevated the systolic and diastolic levels of the CaT with slight lengthening of the total transient duration. The systolic peak of CaT is highly dependent on the amount of Ca<sup>2+</sup>-TRPN bound. Increasing length induces small increments in the CaT amplitude and more substantial ones in the amount of Ca<sup>2+</sup> bound to TRPN. Moreover, the role of individual  $I_{\rm sac}$  currents means that  $I_{\rm ns}$  current formulation is more sensitive to changing  $\lambda$  than  $I_{K0}$  current (Fig. A9D). In other words, given no rise in CaT peak for  $I_{K0}$  current (right panels) for  $\lambda = 1.3$  in purple dotted-dashed line, the rise in  $F_{\text{active}}$  (Fig. A9B) is because of increased sensitivity of the thin filament. In contrast, in the presence of  $I_{ns}$ current (left panels), the increased CaT peak (for  $\lambda = 1.2$ , yellow dotted line) combines with an increased affinity for bound Ca-TRPN (Fig. A9C) to elevate  $F_{\text{active}}$  (Fig. A9B). This result is consistent with experiments by Tavi et al. in rat atrial myocytes (Tavi et al., 1998) highlighting Ca<sup>2+</sup> systolic rise and Ca<sup>2+</sup> buffering by TRPN as the major modulators of the contractile force. The enhanced sensitivity of Ca<sup>2+</sup>-TRPN binding can be explained by the length-dependent increase of contraction (Lookin et al., 2021). The study relates the extent of the length dependent activation with both peak systolic CaT and its decaying kinetics. At higher rates (3.5 Hz and above), the stretched sarcomeres reduced the alternans incidence threshold under control conditions from 4.6 to 3.5 Hz. These findings are quite similar to those from Tsai et al. (2011), who showed increased alternans vulnerability at BCL 250–67 ms for a longer SL. Therefore, atrial stretch is a important modulator of alternans development and can be considered as a potential therapeutic target to reduce the propensity for arrhythmia in patients with AF.

In our contraction module, as in experiments, the  $EC_{50}$ value of force-pCa curve is a linear inverse function of increasing SL while varying the SL within the physiological range. Hence, the length dependent decrease in  $EC_{50}$  (or increased sensitivity of  $Ca^{2+}$  binding to TRPN) results in increased bound  $\dot{Ca}^{2+}$  in the cytosol in spite of quite similar levels of free Ca<sup>2+</sup>. This length dependent rise in Ca<sup>2+</sup> bound to TRPN is further amplified once translated into active force (Fig. 8). The model was also able to predict the secondary effects of MEF. First, the decrease of I<sub>CaL</sub> in both peak and plateau phases (not shown); the increasing calcium dependent inactivation (CDI) effect, because of increasing  $J_{relss}$  peak (not shown), is the main modulator of  $I_{CaL}$  inactivation with stretch. On the other hand, Na<sup>+</sup> accumulation due to stretch increased the inward  $I_{NaCa}$  current at basal rate (Fig. 10C). Under increasing stretch, the combination of increased inward  $I_{\rm sac}$ , reduced  $I_{\rm CaL}$  and increased  $I_{\rm NaCa}$  inward component resulted in the lengthening of APD<sub>90</sub>.

# Rate dependent modulation of MBS2023 and comparison with *in-silico* studies

Computational modelling and simulation are powerful tools for accelerating the mechanistic understanding of ionic mechanisms of the heart and eventually the whole heart. For human atrial myocytes, the CRN (Courtemanche et al. 1998) and Nygren (Nygren et al., 1998) models were the pioneer in-silico models, with a detailed description of ionic membrane currents and different descriptions of CaT dynamics, as reviewed by Cherry (Cherry et al., 2008). However, neither model was able to reproduce slow Ca<sup>2+</sup> wave-like diffusion towards the centre of the cell, and experimental phenomena like a correct rate dependence of CaT, as shown in our previous work using the CRN model (Mazhar et al., 2021) and in Fig. 9J (red line). The rate dependent trend of  $APD_{90}$ and CaT for the CRN model can be seen in Fig. A8 (blue dashed lines). Later, KM2011 was proposed as the model with a detailed Ca<sup>2+</sup> wave 'fire-diffuse-fire-propagation' mechanism. However, the model showed quite different results for Ca<sup>2+</sup> rate dependence when compared with the dome shaped experimental trend (Fig. A8, blue dotted lines). One explanation could be that the  $Ca^{2+}$  homeostasis validity of the model was confined to a basal frequency of 1 Hz only. Among the existing human atrial computational models, the one by Grandi et al. (2011) demonstrated a positive CaT rate dependency (Fig. A8, blue dotted-dashed line). However, the rate dependence trend for APD<sub>90</sub> was compromised at lower BCL (also

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reviewed in Dössel et al., 2012). In contrast, MBS2023 adapted the Ca<sup>2+</sup> handling formulations from KM2011 and calibrated them for a set of physiological pacing rates, i.e. 0.5 to 3 Hz. Moreover, the inclusion of the MCF effect resulted in a more physiological shape of CaT and therefore of the  $F_{\text{active}}$  dynamics, as quantified by ttp and rt<sub>50</sub> biomarkers.

Our model reproduces the APD<sub>90</sub> rate dependence as shown by the APs within type-3 AP morphology (Fig. 9B). In line with experimental findings (Dawodu et al., 1996), for longer BCLs (2 to 1 s), our model shows a less steep APD<sub>90</sub> rate dependence, whereas for shorter BCLs (0.5 s onwards) the trend gets steeper, with decreasing BCL increasing rate. We analysed the ionic rate dependent mechanisms underlying APD rate dependency. We observed that the rate dependence of  $I_{CaL}$  inactivation results in a decreasing inward current with BCL (Fig. 9D). The effect of rate dependent reduction in  $I_{CaL}$  current on  $APD_{90}$  was quantified by blocking the current by 50% as shown in Fig. 17 (blue dashed line on the left). Block of  $I_{CaL}$  results in a flat APD<sub>90</sub> rate dependency for BCL [2,1] similar to what was observed by Wagoner et al. by the application of 10  $\mu$ M nifedipine (Van Wagoner et al., 1999), as shown in Fig. 17 (left in green markers). This was further confirmed by the observation made by Li et al. on human atrial myocytes, where the reduction of  $I_{CaL}$  was dominant for higher BCLs (0.5 to 2 s) (Li et al., 1997). Another test was to clamp  $[Na^+]_{ss}$  to its steady state value and quantify the impact on rate adaptation of  $APD_{90}$  and  $CaT_{bc}$  as shown in Fig. 17 (dotted line).  $[Na^+]_{ss}$ clamp was similar to I<sub>CaL</sub> block; however, its effect was dominant from BCL 2 to 0.5 s. Thus, a combination of 50%  $I_{CaL}$  block and  $[Na^+]_{ss}$ -clamp (dotted-dashed line) was simulated that reversed the shortening of APD<sub>90</sub> for BCLs 2 to 0.33 s to give a lengthening effect. Since  $I_{Kur}$ demonstrates rate dependency for the physiological range (Fig. 9*F*), we therefore made a test by blocking 50% of the current (dashed red line). Block of  $I_{Kur}$  can enhance the plateau of the AP, hence resulting in more activation of  $I_{CaL}$ 

Figure 17. Effect of blocking each

MBS2023 original rate dependence

(continuous line) is compared with ICaL

(Van Wagoner et al., 1999) depicts the

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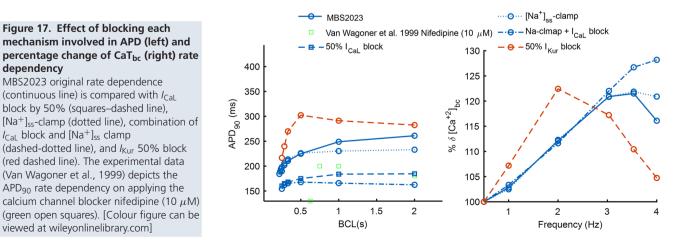
block by 50% (squares-dashed line),

I<sub>Cal</sub> block and [Na<sup>+</sup>]<sub>ss</sub> clamp

dependency

current, which eventually leads to a reverse rate dependent APD<sub>90</sub> trend, i.e. lengthening effect with rate. This was also observed in experiments blocking I<sub>Kur</sub> with Acatein  $(10 \,\mu\text{M})$  (Li et al., 2008).  $I_{\text{Kur}}$  block gives a similar response to the 50%  $I_{CaL}$  block +  $[Na^+]_{ss}$ -clamp experiment though the latter reduced the overall APD<sub>90</sub> magnitude whereas the former enhanced it considerably. In other words, the magnitude and rate dependency of the current integrals  $I_{CaL}$  and  $I_{NaCa}$  (Fig. 9H) were reduced for combined 50%  $I_{CaL}$  block +  $[Na^+]_{ss}$ -clamp simulation test, and were enhanced for the 50%  $I_{\rm Kur}$  block test. For the 50%  $I_{\rm Kur}$ block test, the increased activation of I<sub>CaL</sub> results in a large I<sub>NaCa</sub> current with its reverse operating mode switched at BCL 0.5 s, which is the major player behind the rate adaptation of the APD trend under IKur block. Overall, for longer BCLs, I<sub>CaL</sub> reduction was a dominant mechanism of APD<sub>90</sub> shortening, whereas for shorter BCLs a combined increase in outward  $I_{\text{NaK}}$  pump and reverse mode of the I<sub>NaCa</sub> exchanger currents resulted in a steep APD<sub>90</sub> rate dependence.

As a result, our model shows an increase in CaT and  $F_{\text{active}}$  peaks with rate (Fig. 9J) up to 3.5 Hz and then a small decline afterwards. Experimental data for this rate dependence are quite sparse, making a close fit between the experimental data and the model results meaningless. Rather our focus was to investigate the mechanisms that give rise to biphasic rate dependence of the CaT. Therefore, we performed a detailed analysis of the inward and outward Ca<sup>2+</sup> fluxes to/from the bc compartment and their rate dependence and quantified the contribution of each flux at each increasing rate. The major source of Ca<sup>2+</sup> accumulation with rate is the prominent rise of diastolic  $Ca^{2+}$  (see Fig. 9M) because less time is available for its complete removal from the bulk. The  $Ca^{2+}$  efflux is mainly dependent on  $J_{up}$ , whereas transient sequestration of Ca<sup>2+</sup> by  $J_{Ca-TRPN}$  during the systolic phase strongly modulates the CaT peak, as is also shown by simulations in which the MCF is absent (Fig. 6). An increasing  $Ca^{2+}$  content in SR with the rate (from 0.615 to 0.821 mM for frequencies



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of 1 to 4.5 Hz) is shown by the model and it qualitatively matches experimental evidence (Maier et al., 2000) from human atrial myocytes. An increasing  $[Ca^{2+}]_{SR}$  might produce an increasing  $J_{rel}$ ; however, the model produced a J<sub>relbc</sub> with a very small amplitude that always decreased with rate. The main reason lies in the refractoriness of RyR gate that resulted in unavailability of fully activated RyR channels with increasing frequency. On the other hand,  $J_{\rm relss}$  has faster dynamics then  $J_{\rm relbc}$  and demonstrates a biphasic behaviour with increasing rate (not shown) that impacts on  $J_{\text{diff}}$  rate dependency. Consequently,  $J_{\text{diff}}$  peak increases slightly at lower rates (up to 2 Hz) and then reduces progressively with rate. On the other hand, after 2 Hz, a reduction in outward flux  $J_{Ca-TPRN}$  (asterisks in Fig. 9M) was observed with a very slight contribution of  $J_{up}$  (open circle in Fig. 9*M*). Hence, this balancing of Ca<sup>2+</sup> fluxes was observed at each frequency; therefore, inward flux was dominant till 3.5 Hz after which the outward flux dominated.

### Inotropic effect of IKur blocking

 $I_{Kur}$  block enhances the atrial AP plateau, which can be therapeutically beneficial for treating atrial contractile dysfunction (Schotten et al., 2007), thus reducing the AF burden in persistent AF (Ford et al., 2016) and chronic AF (Wettwer et al., 2004) groups. In accordance with the experiments, the model can reproduce and explain the positive inotropic response on  $I_{Kur}$  block (Fig. 10). The elevated plateau potential and duration predicts secondary effects of IKur block on membrane currents like increases in  $I_{CaL}$  and  $I_K$  ( $I_{Ks}$  and  $I_{Kr}$ ) currents. On the other hand, the model also produces a contradictory result, i.e. a lengthening of APD<sub>90</sub> instead of a shortening as reported in human atrial myocytes (Wettwer et al., 2004). The reason for this discrepancy lies in the low maximal conductance value of IKr current present in our AP module, which was adopted from the parent model. Indeed, KM2011 inherited that conductance value from the Nygren model, which was shown to have a very small  $I_{\rm Kr}$  with respect to CRN (Cherry et al., 2008). A detailed calibration of IKr formulation for human atrial myocytes will be the objective for future developments of our model. Therefore, the overall effect of  $I_{Kur}$  block was a prolongation of APD<sub>90</sub>. In addition, the type-3 AP morphology has no substantial effect on APD<sub>90</sub> with *I*<sub>Kur</sub> block; it still produces an increase in contractility, as observed by Schotten in canine myocytes (Schotten et al., 2007). Based on these observations, the model predicted that a strong positive inotropic effect is not dependent on the late repolarization phase of the AP, whereas the AP plateau phase is crucial for this phenomenon.

# Depressed SERCA activity can affect the alternans threshold

Reduced atrial contractile function is a predictor of PoAF (Fakuade et al., 2021). Based on this fact, it was observed experimentally that SERCA activity reduction was the only mechanism of impaired Ca<sup>2+</sup> handling in a PoAF patient group. Based on this observation, we simulated the SERCA activity when reduced to 75% and analysed the set of mechanisms underlying the onset of alternans in the model (Fig. 11). The model produced alternans at 4 Hz under PoAF but not under control conditions. In our model, under PoAF conditions, due to SERCA depression, the removal of calcium from the cytoplasm and refilling of SR between two successive releases is slowed down (Fig. 11D and E). As a consequence, the diastolic value of Ca<sup>2+</sup> in ss is relatively high, keeping it in a range for which recovery from inactivation of RyRs is very small. In this condition, the amount of available RyRs at the beginning of two successive releases can change a lot (more than doubled, from 0.024 to 0.058, in Fig. 11E). In other words, the 'gain' of the feedback loop into the intracellular calcium handling system is very high, and this gives rise to oscillations/alternans. Ca<sup>2+</sup> dependent channels and transporters control the translation of beat-to-beat alternans in CaT into AP shape. Therefore, CaT is the driving mechanism of AP alternans, as demonstrated by comparing the APD<sub>50</sub> as a measure of AP alternans. Hence, under PoAF, reduced SERCA activity results in the slowing of  $J_{up}$ , which can promote alternans even at modest SR loads and small release, consistent with the results found in rabbit atrial myocytes (Shkryl et al., 2012). Consequently, the slow uptake results in less  $Ca^{2+} J_{rel}$  flux and the small release with slow recovery kinetics contributed to moving the alternans threshold down to 4 Hz under PoAF conditions, which matches the findings of human atrial myocytes (Llach et al., 2011).

We have also observed, under MEF conditions, that alternans incidence was demonstrated at 3.5 Hz, as shown in Fig. 15 (yellow crosses). Similarly, the dynamic alternans protocol simulation results show that parameters like CaMKII overexpression, reduced SERCA pumping rate, fast RyRtau<sub>inact</sub>, and the increased Ca<sup>2+</sup>-TRPN sensitivity factor protect the model from alternans vulnerability, whereas the weak MCF condition can enhance it. In principle, CaMKII inhibition should suppress abnormal increases in activity of the RyR inactivation gate, providing a possible strategy to protect against alternans incidence. Nevertheless, the suppression of RyR activity can result in an increased propensity for Ca<sup>2+</sup> alternans because of the prolongation of RyR refractoriness, as was seen experimentally in isolated mice ventricular myocytes (Saderdini et al., 2021). Therefore, the development of alternans incidence in the model highlights the role of possible chains of mechanisms involved. In addition, the PoAF condition simulated in the model is based on one mechanism, i.e. depressed SERCA activity and we speculate that other mechanisms are also involved, as was observed by Van Wagoner et al. (1999), who showed that a  $g_{CaL}$  percentage raised to 50% was the main mechanism in the patients who developed PoAF condition. This point was also discussed in the work of Fakaude et al. and no clear reason for this discrepancy was found. Thus, we have observed alternans threshold shifting in the model that is a qualitative representation of Fakuade et al.'s experiments.

#### Limitations

We have presented a fully integrated human atrial model that highlights novel insights into various physiological mechanisms. For model parameterization and evaluation (Fig. 6), we utilized extensive human-specific independent atrial datasets to generate human-specific results and conclusions that are highly translatable. Nonetheless, lack of data availability or large variability in available datasets introduces some uncertainty. As a relevant example, the experimental data available on CaT biomarkers in humans show a high level of variability due to the differences in dyes and temperature used. In addition, the mechanical conditions under which the CMs are tested also differ substantially, and the available data for atrial human CMs are also quite limited. While the model captures the CaT characteristics, the systolic and diastolic levels and rt<sub>50Ca</sub> well, it produces a slow TT<sub>Ca</sub> and therefore a slow twitch time. A possible reason behind this could be the delayed efflux of  $Ca^{2+}$  from the bulk, which is solely based on  $J_{uv}$ . However, the small amount of data on Ca<sup>2+</sup> kinetics in human atrial myocytes limits the robustness of the comparison. The rate dependent trend for CaT and  $F_{\text{active}}$  fits qualitatively with the dome shaped experimental data shown. Potential parameters and mechanisms were highlighted; however, we were not able to obtain the right rate dependency for all the characteristics of the model since we have to deal with so many parameter calibrations at a time. In future, further optimization and validation of the presented biomarkers will be performed as new data becomes available.

It is known that T-tubules also exist in atrial myocytes (Richards et al., 2011) and the transverse-axial tubule system can play an important role in modulating intracellular calcium handling and, therefore, electromechanical coupling. The inclusion in the model of these aspects is beyond the scope of the present work and will be addressed in future developments focusing on pathological conditions.

#### Conclusions

In this work, we have presented a new electro-mechanical model for human atrial cardiomyocytes along with mechano-electric regulation. The model was developed, calibrated and evaluated against a wide range of *in-vitro* data. We found that our physiologically accurate description of  $Ca^{2+}$  handling can reproduce many experimental phenomena and help us to gain insights into the pathophysiological mechanisms. This computationally efficient and coupled model opens new pathways for future multi-scale modelling and investigation of human atrial electromechanics.

#### Appendix 1: Modified model equations

The model equations are based on KM2011 and RDQ2020 models. Here we list all those that have been modified or added.

Volume of cytosolic region

$$V_{\rm bc_{Total}} = \sum_{i = 1 \ to \ 4} V_{\rm bc_i}$$

where  $V_{bc1}$ - $V_{bc4}$  refers to the volume of four bc compartments in KM2011 model and  $V_{bc_{Total}}$  is the total volume of the bc compartment in MBS2023. *Volume of SR compartment* 

$$V_{\rm SRss} = 2.25\% (V_{\rm ss} + V_{\rm bc4}),$$

$$V_{\text{SRbc}} = 2.25\% (\text{sum} (V_{\text{bc1}} : V_{\text{bc3}}))$$

*Intracellular free*  $Ca^{2+}$  *concentration* 

$$\frac{\mathrm{dCai}_{\mathrm{bc}}}{\mathrm{d}t} = \mathrm{beta}_{\mathrm{nj}} \frac{J\mathrm{Ca}}{V_{\mathrm{bc}_{\mathrm{Total}}}}.$$

Addition of dynamic troponin effect,  $\frac{d[Ca^{2+}]_{TRPN}}{dt}$ , in  $[Ca^{2+}]_i$ :

$$\frac{\mathrm{dCai}_{\mathrm{bc}}}{\mathrm{d}t} = \mathrm{beta}_{\mathrm{bc}} \frac{J\mathrm{Ca}}{V_{\mathrm{bc}_{\mathrm{Total}}}} - \frac{\mathrm{d}[\mathrm{Ca}^{2+}]_{\mathrm{TRPN}}}{\mathrm{d}t}.$$

*Free*  $Ca^{2+}$  *in SR* 

$$\begin{aligned} \frac{\mathrm{dCa}_{\mathrm{SR}_{\mathrm{ss}}}}{\mathrm{d}t} &= \mathrm{beta}_{\mathrm{SR}_{\mathrm{ss}}} \left( \mathrm{DCa}_{\mathrm{SR}} \left( \frac{-\mathrm{Ca}_{\mathrm{SR}_{\mathrm{ss}}} + \mathrm{Ca}_{\mathrm{SR}_{\mathrm{bc}}}}{\mathrm{d}x^2} \right. \right. \\ &+ \frac{\mathrm{Ca}_{\mathrm{SR}_{\mathrm{ss}}} - \mathrm{Ca}_{\mathrm{SR}_{\mathrm{bc}}}}{2\left(4\mathrm{d}x^2\right)} \right) + \frac{J_{\mathrm{SRCa}_{\mathrm{ss}}}}{V_{\mathrm{SR}_{\mathrm{ss}}}} \right), \end{aligned}$$

$$\begin{aligned} \frac{\mathrm{dCa}_{\mathrm{SR}_{\mathrm{bc}}}}{\mathrm{d}t} &= \mathrm{beta}_{\mathrm{SR}_{\mathrm{bc}}} \left( \mathrm{DCa}_{\mathrm{SR}} \left( \frac{\mathrm{Ca}_{\mathrm{SR}_{\mathrm{ss}}} - \mathrm{Ca}_{\mathrm{SR}_{\mathrm{bc}}}}{\mathrm{d}x^2} \right. \\ &+ \frac{\mathrm{Ca}_{\mathrm{SR}_{\mathrm{ss}}} - \mathrm{Ca}_{\mathrm{SR}_{\mathrm{bc}}}}{2\left(3\mathrm{d}x^2\right)} \right) + \frac{J_{\mathrm{SRCa}_{\mathrm{bc}}}}{V_{\mathrm{SR}_{\mathrm{bc}}}} \right). \end{aligned}$$

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 $I_{\rm Na}$  current modified as shown,

$$E_{\rm Na} = \frac{RT}{F} \log\left(\frac{{\rm Na}_{\rm o}}{{\rm Na}_{\rm ss}}\right),$$

$$m_{\infty} = \frac{1}{1 + \exp((V + 36.3) / - 7.8)},$$

$$h_{\infty} = \frac{1}{1 + \exp\left(\frac{V+67}{5.6}\right)}$$

$$m_{\text{tau}} = 0.00001 + 0.00013 \exp\left(-\left(\frac{V+48}{15}\right)^2\right) + \frac{0.000045}{1 + \exp((V+42)/-5)},$$

$$h 1_{\text{tau}} = 0.00007 + \frac{0.034}{1 + \exp((V + 41)/5.5)} + \exp\left(-\frac{V + 41}{14}\right) + \frac{0.0002}{1 + \exp(-(V + 79)/14)}$$

$$h 2_{\text{tau}} = 0.0007 + \frac{0.15}{1 + \exp((V + 41)/5.5)} + \exp\left(-\frac{V + 41}{14}\right) + \frac{0.002}{1 + \exp(-(V + 79)/14)}$$

$$I_{\rm Na} = g_{\rm Na} \ mh1h2 \left( V - E_{\rm Na} \right).$$

## **Appendix 2: Figures**

*L-type* Ca<sup>2+</sup> current reformulation

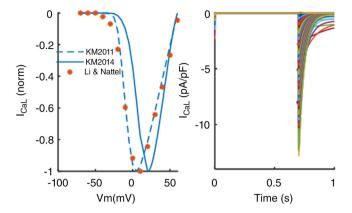
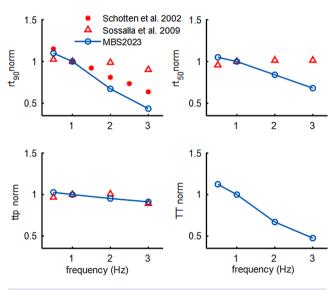


Figure A1. Voltage clamp protocol reproduced from Li & Nattel (1997)

Left panel shows *I-V* plot for two different formulations, i.e. the modified KM2011 used by MBS2023 (dashed line) and the KM2014 (continuous line) in comparison to the experimental data points. [Colour figure can be viewed at wileyonlinelibrary.com]

Rate dependent characteristics of contractility biomarkers



**Figure A2. Rate dependent characteristics of**  $F_{active}$  **biomarkers** The experimental data reported are from Table A2 and are normalized with respect to 1 Hz frequency. Relaxation kinetics at 90% and 50% of the peak ( $rt_{90}$  and  $rt_{50}$ ) are shown in top left and right panels, time to peak (ttp) in the bottom left panel and the twitch time TT is shown on the bottom right panel. [Colour figure can be viewed at wileyonlinelibrary.com]

# MBS2023 model ionic currents, $Ca^{2+}$ handling and CaMKII inhibition

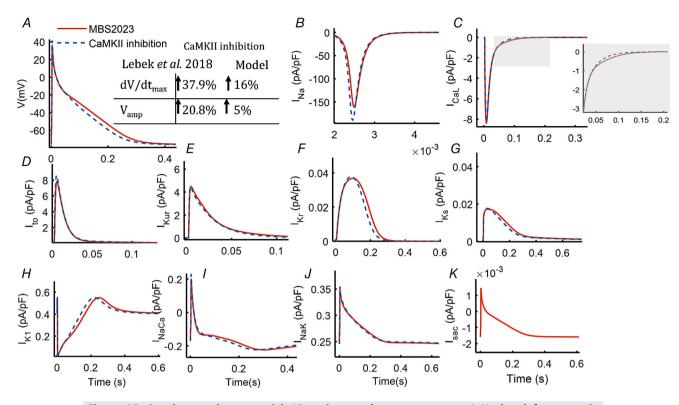


Figure A3. Steady-state human atrial AP, major membrane currents at 1 Hz basal frequency in comparison with CaMKII inhibition (blue dashed line) in MBS2023 (red continuous line) A shows the AP morphology for MBS2023, table in the inset illustrates relative percentage change in  $V_{amp}$  and  $dV/dt_{max}$  from Lebek et al. (2018) in comparison to the model output. B and C show the fast  $I_{Na}$  and the  $I_{CaL}$  current with magnified plateau region of  $I_{CaL}$  in the inset. D and E show the phase 1 repolarizing currents  $I_{to}$  and  $I_{Kur}$ . F, G and H show  $I_K$  currents for phase 2 repolarization, i.e.  $I_{Kr}$ ,  $I_{Ks}$ , and  $I_{K1}$ . I shows the stretch activated current  $I_{sac}$ . J and K represent the exchanger  $I_{NaCa}$  and the pump currents  $I_{NaK}$ . [Colour figure can be viewed at wileyonlinelibrary.com]

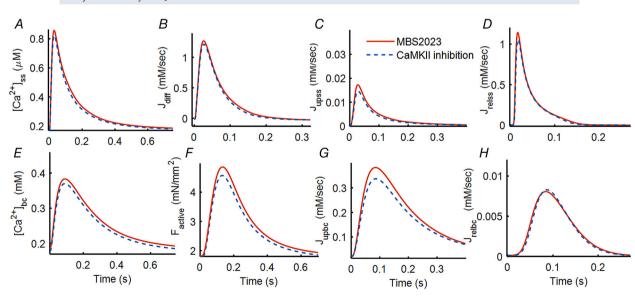
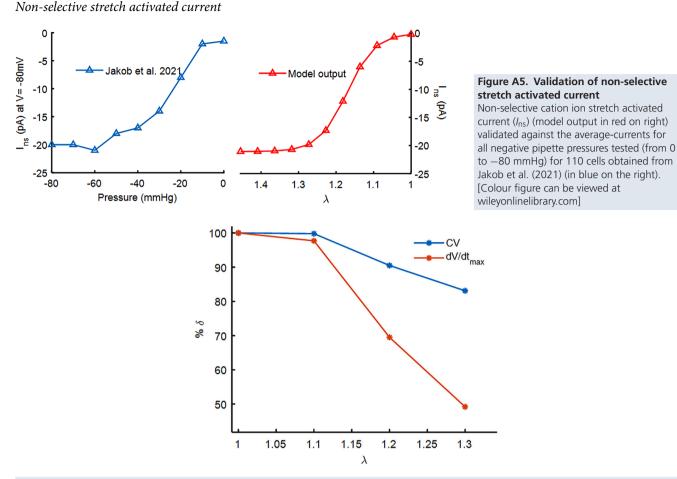


Figure A4. Ca<sup>2+</sup> handling in MBS2023 model (red continuous line) in comparison with CaMKII inhibition (blue dashed line)

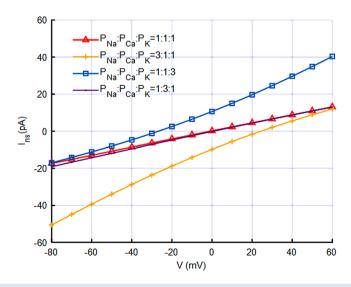
*A*, Ca<sup>2+</sup> transient in the ss compartment. *B*, diffusion from ss to bc ( $J_{diff}$ ). *C*, uptake flux via SERCA pump ( $J_{upss}$ ). *D*, release flux ( $J_{relss}$ ). *E*, slow and smaller Ca<sup>2+</sup>-transient in the bc. *F*,  $F_{active}$  produced. *G*, uptake flux ( $J_{upbc}$ ). *H*, small release flux ( $J_{relbc}$ ). [Colour figure can be viewed at wileyonlinelibrary.com]

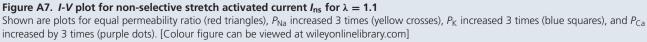
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#### Figure A6. Shortening effect of stretch ratio $\lambda$ on cell level parameters

The maximal shortening velocity ( $dV/dt_{max}$ ) is shown in red and the tissue level parameter, the conduction velocity (CV) is shown in blue. The 1D-discrete tissue level simulation was carried out by connecting 40 cells, with gap junctional resistance ( $R_{gap}$ ) 1.5 M $\Omega$ , for 10 beats using stimulation current of 5.06 nA. The CV was calculated based on the propagation time between cell number 8 and cell number 16. [Colour figure can be viewed at wileyonlinelibrary.com]

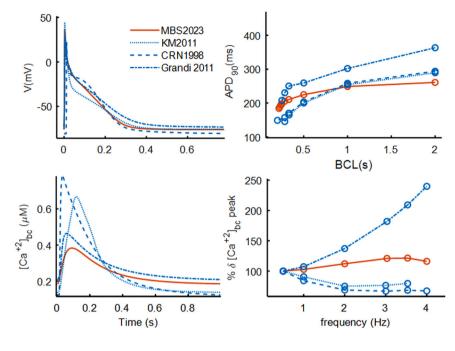




## Comparison with in-silico models

#### Figure A8. Comparison of human atrial computational models characteristics with MBS2023 for basic cycle lengths ranging from 2 to 0.25 s

The Courtemanche (CRN) 1998 model is shown in blue dashed lines, the Grandi, 2011 model in blue dotted-dashed lines, and the Koivumaki (KM) 2011 model in blue dotted lines. The MBS2023 model is shown in red continuous lines. Action potential (AP) morphology is shown in the top left panel, AP duration at 90% of repolarization (APD<sub>90</sub>) rate adaptation in the top right panel, Ca<sup>2+</sup> transient (CaT) time course in the bottom left panel, and rate dependency of percentage change in CaT ( $\%\delta$ CaT) in the bottom right panel. KM2011 demonstrated an irregular waveform at 4 Hz, hence its simulation value is not shown in the top and bottom right panels. [Colour figure can be viewed at wileyonlinelibrary.com]



Effect of non-selective  $(I_{ns})$  and potassium-selective  $(I_{K0})$  cation currents on MBS2023

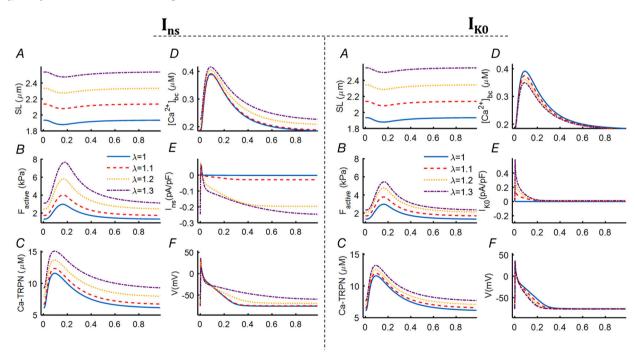
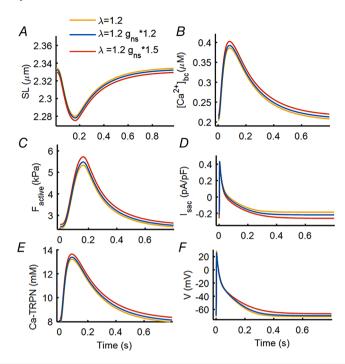


Figure A9. MBS2023 model characteristics for stretch variations using non-selective and selective cations current,  $I_{ns}$  (left panels) and  $I_{K0}$  (right panels), respectively

Stretch is varied by increasing preload effect, i.e.  $\lambda$  from 1 (blue continuous line), 1.1 (red dashed lines), 1.2 (yellow dashed line), and 1.3 (purple dotted-dashed line). *A*, the increasing length of contractile element ( $I_{CE}$ ). *B*, active contractile force ( $F_{active}$ ). *C*, the concentration of Ca<sup>2+</sup> bound to troponin (TRPN). *D*, Ca<sup>2+</sup> transient in bulk cytoplasm (bc). *E*,  $I_{ns}$  (left) and  $I_{KO}$  (right) variation with stretch. *F*, the action potential (AP). [Colour figure can be viewed at wileyonlinelibrary.com]

## Effect of increased conductance of $I_{ns}$ current on MBS2023



### Figure A10. Increased conductance $g_{ns}$ of $I_{ns}$ current for $\lambda = 1.2$ response of model MBS2023

Basal state of the model is in yellow, 20% increase in  $g_{ns}$  is in blue, and 50% increase in  $g_{ns}$  is in red. A, length of contractile element ( $I_{CE}$ ). B,  $F_{\text{active}}$ . C, concentration of Ca<sup>2+</sup> bound to TRPN. D, Ca<sup>2+</sup> transient in the bulk cytosol. E, increasing  $I_{\text{sac}}$  current. F, the action potential. [Colour figure can be viewed at wileyonlinelibrary.com]

## **Appendix 3: Tables data references and** model parameters

Reference	Mean (pF)	SD	Ν
(Voigt et al., 2012)	114.8	33.2	26
(Wettwer et al., 2004)	101	35	29
(Gulais et al., 2004)	99.4	9.5	16
(Dobrev et al., 2001)	92	65.7	30
(Voigt et al., 2014)	113.6	36	35
(Voigt et al., 2010)	94		
(Dobrev et al., 2002)	84	32.9	68
(Fermini et al., 1992)	73.7	24.3	10
(Fakuade et al., 2021)	95.97	51.1	78
(Wang et al., 1993)	66.1	19.2	34
(Kamkin et al., 2003)	72.4	6.4	14
(Schotten et al., 2002)	72	3	167
(Li et al., 1997)	67	35.4	35
(Van Wagoner et al., 1999)	67.6	26.8	86
(Madsen et al., 2004)	58	7.8	16
(Giles & Imaizumi 1988)	54.3	5.87	5
(Nygren et al., 1998)	51.9	3.5	52
(Wang et al., 1999)	29.6	1.8	

Reference	Tissue preparation	Biomarker	Frequency dependent Biomarkers
(Schotten et al., 2002)	Muscle strips obtained from right atrial appendages from patients of mitral valve surgery at 1 Hz, 37°C (n = 31)	F <sub>max</sub> , rt <sub>90Factive</sub>	F <sub>max</sub> , rt <sub>90Factive</sub>
(Schwinger et al., 1993)	Right atrial trabeculae from patients who underwent aortocoronary bypass operations at 1 Hz, 37°C ( $n = 9$ )	F <sub>max</sub>	F <sub>max</sub>
(Sossalla et al., 2010)	Thin right atrial trabeculae were micro-dissected ( $n = 79$ )	F <sub>max</sub> , ttp <sub>Factive</sub> , rt <sub>50Factive</sub> , rt <sub>90Factive</sub>	F <sub>max</sub> , ttp <sub>Factive</sub> , rt <sub>50Factive</sub> , rt <sub>90Factive</sub>
(L S. Maier et al., 2000)	Muscle strips from right atrial trabeculae obtained from patients undergoing aortocoronary bypass operation at $37^{\circ}$ C ( $n = 15$ )	$F_{max}, \Pi_{Factive}, ttp_{Factive}, \\ rt_{50Factive}, rt_{90Factive}, \\ CaT_{max}$	F <sub>max</sub> , CaT <sub>max</sub>
(Flesch et al., 1997)	Isolated electrically driven human right atrial trabecula from non-failing hearts at 1 Hz, $37^{\circ}$ C ( $n = 15$ )	ttp <sub>Factive</sub> , rt <sub>50Factive</sub>	
(Brixius et al., 1997)	Right atrial tissue from patients having aortocoronary bypass surgery ( $n = 19$ )	ttp <sub>Factive</sub> , rt <sub>50Factive</sub> , CaT <sub>max</sub>	
(Schotten et al., 2007)	Muscle strip preparations from atrial trabeculae obtained from right atrial appendages at 1 Hz, 37°C ( $n = 14$ )	F <sub>max</sub>	
(Brixius et al., 1999)	Auricular trabeculae were selected from right atrial tissue. Using Fura-2 ratio method for Ca transient and force in muscle strips at 37°C, 1 Hz	F <sub>max</sub> , F <sub>min</sub> , ttp <sub>Factive</sub> , rt <sub>50Factive</sub> , CaT <sub>max</sub> , CaT <sub>min</sub> , ttp <sub>CaT</sub> , rt <sub>50CaT</sub>	F <sub>max</sub> , F <sub>min</sub> , ttp <sub>Factive</sub> , rt <sub>50Factive</sub> , CaT <sub>max</sub> , CaT <sub>min</sub> , ttp <sub>CaT</sub> , rt <sub>50CaT</sub>
(Voigt et al., 2012)	Human atrial myocytes isolated from right atrial appendage at 37°C for 40 myocytes extracted from 20 patients	CaT <sub>min,</sub> CaT <sub>sys</sub> , CaT <sub>amp</sub> , τ	
(Voigt et al., 2014)	Human atrial appendages from 73 patients were isolated at 37°C	$CaT_{min}$ , $CaT_{sys}$ , $CaT_{amp}$ , $\tau$	
(Heijman et al., 2020)	71 out of 149 RA-appendages were subject to cardiomyocytes isolation at 37°C	$CaT_{min}$ , $CaT_{sys}$ , $CaT_{amp}$ , $\tau$	
(Grandi et al., 2011)	Atrial myocytes extracted from right atrial appendages were isolated using enzymatic dissociation at 37°C	CaT <sub>min,</sub> CaT <sub>sys</sub> , CaT <sub>amp</sub>	
(Fakuade et al., 2021)	Myocytes ( $n = 78$ ) were isolated from 38 patients extracted from right atrial appendages	CaT <sub>min,</sub> CaT <sub>sys</sub> , CaT <sub>amp</sub> , τ	

#### Table A2. Human atrial experimental data used for calibration of contraction parameters in the Electro-mechanical model

 $F_{max}$ : peak tension;  $F_{min}$ : diastolic tension; ttp<sub>Factive</sub>: time to  $F_{max}$ ; TT<sub>Factive</sub>: twitch time; rt<sub>50Factive</sub>, rt<sub>90Factive</sub>: relaxation time at 50% and 90% of  $F_{max}$ ; CaT<sub>max</sub>: calcium transient peak; CaT<sub>min</sub>: calcium transient diastolic; CaT<sub>amp</sub>: calcium transient amplitude; rt<sub>50CaT</sub>: relaxation time at 50% of CaT<sub>max</sub>; ttp<sub>CaT</sub>: time to CaT<sub>max</sub>;  $\tau$ : decay time.

Parameters	Calibrated data	RDQ2020 calibration	
Troponin sensitivity ( $k_{d0}$ ) (mM)	1.4e-3	0.3e-3	
Troponin dissociation rate $(k_{off})$ (s <sup>-1</sup> )	197.6	100	
Cooperativity factor (y) (-)	13.43	12	
$r_0$ (s <sup>-1</sup> )	55.5	134.31	
α (-)	22.2	25.184	
$\mu_{f_{-}}^{0}$ (s <sup>-1</sup> )	13.5	32.653	
$\mu_{f_p}^0$ (s <sup>-1</sup> ) $\mu_{f_p}^1$ (s <sup>-1</sup> )	0.304	0.778	
a <sub>XB</sub> (kPa)	52.89e3	22.894e3	

Table A3. Contraction module calibration based on biomarkers from human atrial experimental data

## Table A4. Human atrial in-vitro data used for evaluation of action potential biomarkers

Reference	Experimental protocols	Biomarkers	Frequency dependent biomarkers
(Bosch et al., 1999)	Whole cell voltage clamp technique was used to record AP from rod shaped cells of the RA appendages at 36°C	APD <sub>90</sub>	APD <sub>90</sub>
(Van Wagoner et al., 1999)	Whole cell patch clamp technique was used to record AP from isolated atrial myocytes at 35°C	APD <sub>90</sub> , APD <sub>50</sub>	APD <sub>90</sub> , APD <sub>50</sub>
Dobrev et al., 2002	Atrial myocytes were recorded using patch electrodes at 37°C	APD <sub>90</sub> , APD <sub>50</sub> , RMP	APD <sub>90</sub>
(Dawodu et al., 1996)	RA trabeculae were simulated with intracellular microelectrode technique to record APs at 31°C (Type 3 morphology data only were considered)	APD <sub>90</sub> , APD <sub>50</sub> , APD <sub>30</sub> , RMP, d <i>V</i> /dt <sub>max</sub> , APA	RMP, d <i>V</i> /dt <sub>max</sub>
(Lagrutta et al., 2006)	Microelectrode technique was used to measure action potential duration from human atrial myocytes at 35°C.	APD <sub>90</sub> , APD <sub>50</sub> , RMP, d <i>V</i> /dt <sub>max</sub> , APA	
(Wang et al., 1990)	Microelectrode technique was used to measure AP characteristics from atrial strips obtained from right atrial appendage at 36°C	RMP, d <i>V</i> /dt <sub>max</sub> , APA	
(Wettwer et al., 2004)	Microelectrode technique was used in right atrial trabeculae	RMP	
(Dobrev et al., 2001)	Microelectrode technique used to measure the APs from right atrial trabeculae at 37°C	APD <sub>50</sub> , RMP	
(Workman et al., 2003)	Whole cell patch clamp technique was used to record the APs at 35–37°C	APA	
(Ford et al., 2016)	Sharp microelectrode was used to measure action potentials from small pieces of RAA at 36°C and basal frequency of 1Hz	APD <sub>90</sub> , APD <sub>50</sub> , APD <sub>20</sub> , PLT <sub>20</sub> , RMP, dV/dt <sub>max</sub> , APA, ERP	APD <sub>90</sub> , APD <sub>50</sub> , APD <sub>20</sub> , PLT <sub>20</sub> , RMP, dV/dt <sub>max</sub> , APA, ERP

Table A5. Modified parameter values in MBS2023				
A. Ionic current related parameters				
Parameter	MBS2023	KM2011	Motivation	
g <sub>Na</sub> (nS)	350	$P_{\rm Na} = 0.00182 \ {\rm nL} \ {\rm s}^{-1}$	Reformulated according to new experimental data as described in (Skibsbye et al., 2016)	
g <sub>CaL</sub> (nS)	14.5	25.3125	To avoid irregular reactivation of I <sub>CaL</sub> current at the lower frequency of 0.5 Hz	
g <sub>Kur</sub> (nS)	2.35	2.447	To achieve an APD <sub>50</sub> more inline with the experimental bound as shown in Fig. 4 <i>A</i>	
g <sub>K1</sub> (nS)	3.44	3.44	-	
g <sub>to</sub> (nS)	8.175	8.175	-	
$g_{ m NaCa}$ (nS)	0.0083	0.0084 (pA/(mmol/L) <sup>4</sup> )	As suggested by Skibsbye et al., 2016	
$g_{Cab}$ (nS)	0.085	0.0952		
k <sub>Ca</sub> (mM)	0.6e-3	1e-3		
B. CaMKII related par	rameters			
CaMK <sub>0</sub> (-)	0.05			
K <sub>mCaM</sub> (mM)	0.0015			
K <sub>mCaMK</sub> (-)	0.15			
$\alpha_{CaMK}$ (s <sup>-1</sup> )	50			
$\beta_{CaMK}$ (s <sup>-1</sup> )	0.68			
C. Stretch activated c				
K <sub>ns</sub> (-)	100			
α <sub>ns</sub> (-)	30			
g <sub>ns</sub> (nS)	0.461e-6			
f <sub>CE</sub> (kPa)	100			
f <sub>PE</sub> (kPa)	0.006			
f <sub>SE</sub> (kPa)	2.8			
$k_{\rm SE} (\mu {\rm m}^{-1})$	14.6			
$k_{\rm PE} (\mu { m m}^{-1})$	14.6			

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## **Additional information**

## Data availability statement

The original contributions presented in the study are included in the article; further inquiries can be directed to the corresponding author.

## **Competing interests**

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

## Author contributions

S.S. was responsible for the study conception. F.M. and S.S. were responsible for the study design. FM was responsible for simulations. F.M., C.B. and S.S. were responsible for data analysis. F.M., C.B., M.P. and S.S. were responsible for data interpretation. F.M., C.B. and S.S. were responsible for writing the manuscript. F.M., C.B., F.R., M.P., L.D., A.Q., C.C. and S.S. were responsible for the final approval of the manuscript submitted for publication. All authors agree to be accountable for all aspects of the work. All persons designated as authors qualify for authorship, and all those who qualify for authorship are listed.

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#### **Keywords**

 $Ca^{2+}$ -handling, cardiomechanics, computational modelling, contractility, coupling, human atrial cells, rate adaptation, stretch activated currents

## **Supporting information**

Additional supporting information can be found online in the Supporting Information section at the end of the HTML view of the article. Supporting information files available:

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