



Contents lists available at ScienceDirect

Journal of Functional Analysis

journal homepage: www.elsevier.com/locate/jfa

Regular Article

Nonlinear aggregation-diffusion equations with Riesz potentials

Yanghong Huang^a, Edoardo Mainini^{b,*}, Juan Luis Vázquez^c,
Bruno Volzone^d^a Department of Mathematics, University of Manchester, Oxford Road, Manchester M13 9PL, United Kingdom^b Dipartimento di Ingegneria Meccanica, Energetica, Gestionale e dei Trasporti, Università degli studi di Genova, Via all'Opera Pia, 15 - 16145 Genova, Italy^c Departamento de Matemáticas, Universidad Autónoma de Madrid, 28049 Madrid, Spain^d Dipartimento di Matematica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133 Milano, Italy

ARTICLE INFO

Article history:

Received 3 June 2022

Accepted 7 April 2024

Available online 16 April 2024

Communicated by Guido De Philippis

MSC:

35K44

35R11

49K20

Keywords:

Aggregation-diffusion model

Gradient flow

Stationary states

Riesz potential

ABSTRACT

We consider an aggregation-diffusion model, where the diffusion is nonlinear of porous medium type and the aggregation is governed by the Riesz potential of order s . The addition of a quadratic diffusion term produces a more precise competition with the aggregation term for small s , as they have the same scaling if $s = 0$. We prove existence and uniqueness of stationary states and we characterize their asymptotic behavior as s goes to zero. Moreover, we prove existence of gradient flow solutions to the evolution problem by applying the JKO scheme.

© 2024 The Author(s). Published by Elsevier Inc. This is an open access article under the CC BY license (<http://creativecommons.org/licenses/by/4.0/>).

* Corresponding author.

E-mail addresses: yanghong.huang@manchester.ac.uk (Y. Huang), mainini@dime.unige.it (E. Mainini), juanluis.vazquez@uam.es (J.L. Vázquez), bruno.volzone@polimi.it (B. Volzone).

1. Introduction

We consider the Cauchy problem in the whole space \mathbb{R}^d , $d \geq 1$, for the aggregation-diffusion equation

$$\begin{cases} \partial_t \rho = \Delta \rho^m + \beta \Delta \rho^2 - \chi \nabla \cdot (\rho \nabla (K_s * \rho)), \\ \rho(0) = \rho^0, \end{cases} \quad (1.1)$$

where the initial datum ρ^0 is a nonnegative mass density in $L^1(\mathbb{R}^d) \cap L^m(\mathbb{R}^d)$, and the parameters satisfy $\chi > 0$, $\beta \geq 0$, $m > 2$. Here, K_s denotes the Riesz kernel of order $s \in (0, d/2)$, namely $K_s(x) = c_{d,s}|x|^{2s-d}$, being $c_{d,s}$ an explicit normalization constant defined as

$$c_{d,s} := \pi^{-d/2} 2^{-2s} \Gamma(\frac{d}{2} - s) / \Gamma(s) \quad (\text{as } s \downarrow 0 \text{ there holds } c_{d,s} \sim \pi^{-d/2} \Gamma(\frac{d}{2}) s). \quad (1.2)$$

The natural free energy associated with the nonlocal PDE (1.1) is given by

$$\mathcal{F}_s[\rho] = \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m(x) dx + \beta \int_{\mathbb{R}^d} \rho^2(x) dx - \frac{\chi}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho(x) \rho(y) dx dy. \quad (1.3)$$

We notice that (1.1) has the structure of a continuity equation $\partial_t \rho + \nabla \cdot (\rho \mathbf{u}) = 0$, where the velocity vector field is a gradient $\mathbf{u} = -\nabla \psi$ and the velocity potential

$$\psi := \frac{m}{m-1} \rho^{m-1} + 2\beta \rho - \chi K_s * \rho \quad (1.4)$$

is the functional derivative of $\mathcal{F}_s[\rho]$ with respect to ρ . For this reason, the evolution equation (1.1) is formally the gradient flow of functional \mathcal{F}_s with respect to the Wasserstein distance, having the structure described in [1, Chapter 11].

Our first objective is the analysis of stationary states of the dynamics, with most emphasis on their behavior as s becomes small. In fact, in our first result we show that for any given mass $M > 0$, \mathcal{F}_s has a unique minimizer over

$$\mathcal{Y}_M := \left\{ \rho \in L^1_+(\mathbb{R}^d) \cap L^m(\mathbb{R}^d) : \int_{\mathbb{R}^d} \rho(x) dx = M, \int_{\mathbb{R}^d} x \rho(x) dx = 0 \right\},$$

coinciding with the unique radial stationary state of the dynamics with mass M and center of mass at the origin. Properties of stationary states have been thoroughly investigated for $\beta = 0$ and for different ranges of m, s , which are usually classified as follows: by considering the homogeneity property of the terms of functional \mathcal{F}_s , diffusion and aggregation are in balance if m is equal to the critical exponent $m_c := 2 - 2s/d$, which

is the so called fair competition regime that is analyzed in [8,9]. The diffusion dominated regime $m > m_c$ was investigated in [12]. Uniqueness of stationary states has also been proved in the different regimes [10,15,17]. Since the diffusion exponents in (1.1) are greater than 2, here we are considering a diffusion dominated model. However, the competition between the additional quadratic diffusion term and the aggregation term becomes crucial in the small s regime, since the limiting critical exponent is exactly 2.

Let us introduce the precise notion of stationary state. We will check in Section 2 that the assumptions on ρ in the next definition entail $\rho \nabla \psi \in L^1_{loc}(\mathbb{R}^d)$, where ψ is given by (1.4).

Definition 1.1. Let $\rho \in W^{1,1}(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d) \cap C(\mathbb{R}^d)$ be a nonnegative density. Let $\Omega := \{x \in \mathbb{R}^d : \rho(x) > 0\}$ and let ψ be defined by (1.4). We say that ρ is a stationary state for the evolution equation in (1.1) if $\psi \in W^{1,\infty}(\Omega)$ and $\nabla \cdot (\rho \nabla \psi) = 0$ in $\mathcal{D}'(\mathbb{R}^d)$.

We stress that this definition differs from the one appearing in [8] and in later works. The definition that we propose is better suited to treat the small s regime. Indeed, as we explain in Section 2, minimizers of \mathcal{F}_s always satisfy the new definition. We have the following

Theorem 1.1. Let $\beta \geq 0$. If $\beta = 0$ and $s < 1/2$ assume in addition that $m < \frac{2-2s}{1-2s}$. For any mass $M > 0$, there exists a unique stationary state of mass M and center of mass 0. Such steady state is radially decreasing, compactly supported, Hölder on \mathbb{R}^d and C^∞ in the interior of its support. It coincides with the unique minimizer of the energy functional \mathcal{F}_s in the class \mathcal{Y}_M .

If $\beta = 0$ and $s < 1/2$, without the additional restriction $m < \frac{2-2s}{1-2s}$ we are not able to apply the radially result from [11] and [12], and for this reason we do not get the same conclusion, but we shall still obtain uniqueness in the class of radial stationary states.

We are mostly interested in the limiting behavior of stationary states as $s \rightarrow 0$. In this perspective, since $K_s \rightarrow \delta_0$, the limit functional is formally given by

$$\mathcal{F}_0[\rho] := \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m(x) dx + \left(\beta - \frac{\chi}{2} \right) \int_{\mathbb{R}^d} \rho^2(x) dx. \quad (1.5)$$

It is clear that the minimization problem $\min_{\mathcal{Y}_M} \mathcal{F}_0$ is strongly influenced by the sign of the coefficient $\beta - \chi/2$. Indeed, it has solutions if and only if $\beta < \chi/2$, and in such case we will check that there is a unique radially decreasing minimizer, given by the characteristic function of a ball. Our second main result is the following. It will be proven in Section 3, where some illustration of stationary states from numerical simulations will also be provided.

Theorem 1.2. For any $s \in (0, 1/2)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . If $0 \leq \beta < \chi/2$, there exists $\rho \in \mathcal{Y}_M$ such that $\rho_s \rightarrow \rho$ strongly in $L^m(\mathbb{R}^d)$ as $s \downarrow 0$, and

moreover ρ is the unique radially decreasing minimizer of the functional (1.5) over \mathcal{Y}_M . Else if $\beta \geq \chi/2$, we have $\lim_{s \downarrow 0} \mathcal{F}_s[\rho_s] = 0$ and $\rho_s \rightarrow 0$ uniformly on \mathbb{R}^d .

We next focus on the gradient flow structure of evolution problem (1.1). In this case, the initial datum ρ^0 is supposed to belong to $\mathcal{Y}_{M,2}$, being $\mathcal{Y}_{M,2}$ the set of all densities in \mathcal{Y}_M with finite second moment:

$$\mathcal{Y}_{M,2} := \left\{ \rho \in L_+^1(\mathbb{R}^d) \cap L^m(\mathbb{R}^d) : \int_{\mathbb{R}^d} \rho(x) dx = M, \int_{\mathbb{R}^d} x\rho(x) dx = 0, \int_{\mathbb{R}^d} |x|^2 \rho(x) dx < \infty \right\}.$$

The analysis of Keller-Segel models with Newtonian potential as Wasserstein gradient flows is found in [4,3,6,5]. More generally, there are many studies about gradient flow approach for interaction-driven evolutions. The case of the Riesz potential appears in [22], in the analysis of the porous medium equation with fractional pressure introduced in [7]. Problem (1.1) is formally preserving mass, positivity and center of mass, and a solution is naturally seen as a trajectory in the space \mathcal{Y}_M . A narrowly continuous curve $[0, +\infty) \ni t \mapsto \rho(t, \cdot) \in \mathcal{Y}_M$ (i.e., $t \mapsto \int_{\mathbb{R}^d} \varphi(x) \rho(t, x) dx$ is continuous for every continuous bounded function φ on \mathbb{R}^d), is a weak solution to (1.1) if $\rho(0) = \rho^0$ and for every $\varphi \in C_c^\infty(\mathbb{R}^d)$ and every $\eta \in C_c^\infty((0, +\infty))$

$$\begin{aligned} - \int_0^{+\infty} \int_{\mathbb{R}^d} \rho(t, x) \varphi(x) \eta'(t) dx dt &= \int_0^{+\infty} \int_{\mathbb{R}^d} \eta(t) \Delta \varphi(x) (\rho(t, x)^m + \beta \rho(t, x)^2) dx dt \\ &- \frac{(d-2s)c_{d,s}\chi}{2} \int_0^{+\infty} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \eta(t) \frac{(\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x-y)}{|x-y|^{d+2-2s}} \rho(t, x) \rho(t, y) dx dy dt. \end{aligned} \quad (1.6)$$

This notion of solution was introduced in [31] for the case of the interaction with the logarithmic potential, see also [4]. We shall construct weak solutions to problem (1.1) by applying the Jordan-Kinderlehrer-Otto [18] scheme. Therefore, denoting by W_2 the Wasserstein distance of order 2, for a discrete time step $\tau > 0$, we shall solve the recursive minimization problems

$$\rho_\tau^0 = \rho^0, \quad \rho_\tau^k \in \operatorname{argmin}_{\rho \in \mathcal{Y}_M} \left(\mathcal{F}_s[\rho] + \frac{1}{2\tau} W_2^2(\rho, \rho_\tau^{k-1}) \right), \quad k \in \mathbb{N},$$

and we shall prove that piecewise constant in time interpolations of minimizers do converge to a weak solution to (1.1) as $\tau \rightarrow 0$ along a suitable vanishing sequence $(\tau_n)_{n \in \mathbb{N}}$. A weak solution that is constructed in this way, that is, as a limit of the JKO scheme

applied to \mathcal{F}_s , will be called a gradient flow solution. We have the following existence result

Theorem 1.3. *Let $\beta \geq 0$ and $0 < s < \min\{1, d/2\}$. If $\beta = 0$, assume in addition that $d \geq 2$ and $1/2 \leq s < 1$. Let $\rho^0 \in \mathcal{Y}_{M,2}$. Then there exists a gradient flow solution to problem (1.1).*

Solutions are global-in-time, as expected in a diffusion dominated model. Further properties of solutions will be described in Section 4, along with some numerical analysis of the problem in Section 5. We shall also discuss an interesting feature: radial decreasing initial data do not necessarily preserve radial monotonicity during the dynamics.

We finally turn the attention to the behavior of solutions as $s \rightarrow 0$. The formal limiting equation reads

$$\partial_t \rho = \Delta \rho^m + (\beta - \chi/2) \Delta \rho^2, \quad (1.7)$$

and its behavior is again crucially depending on the sign of the coefficient $\beta - \chi/2$. Here, we limit ourselves to treat the case $\beta \geq \chi/2$, in which (1.7) is a standard degenerate parabolic equation. We have the following

Theorem 1.4. *Let $\beta \geq \chi/2$. Let $\rho^0 \in \mathcal{Y}_{M,2}$. Let $(s_n)_{n \in \mathbb{N}} \subset (0, 1/2)$ be a vanishing sequence, and for every $n \in \mathbb{N}$ let ρ_n be a gradient flow solution to (1.1) with $s = s_n$. Then the sequence $(\rho_n)_{n \in \mathbb{N}}$ admits strong $L^2_{loc}((0, +\infty); L^2(\mathbb{R}^d))$ limit points. If ρ is one of such limit points, then $[0, +\infty) \ni t \mapsto \rho(t, \cdot)$ is narrowly continuous with values in $\mathcal{Y}_{M,2}$, $\rho(0, \cdot) = \rho^0$ and ρ is a distributional solution to the nonlinear diffusion equation (1.7), i.e.,*

$$-\int_0^{+\infty} \int_{\mathbb{R}^d} \rho(t, x) \varphi(x) \eta'(t) dx dt = \int_0^{+\infty} \int_{\mathbb{R}^d} \eta(t) \Delta \varphi(x) (\rho(t, x)^m + (\beta - \chi/2) \rho(t, x)^2) dx dt$$

for every $\varphi \in C_c^\infty(\mathbb{R}^d)$ and every $\eta \in C_c^\infty((0, +\infty))$.

Plan of the paper

In Section 2 we discuss Definition 1.1 and prove uniqueness of stationary states along with some regularity properties. Asymptotic behavior of stationary states as s approaches 0 is investigated in Section 3. In Section 4 we prove the theorems about the evolution problem. In Section 5 we provide some numerical examples demonstrating other phenomena, and in Section 6 we provide a discussion on some open problems.

2. Stationary states and minimizers of the free energy

We start this section by discussing the definition of stationary states for the evolution equation in (1.1), see Definition 1.1. We begin with the following lemma about Riesz potentials of bounded continuous $W^{1,1}(\mathbb{R}^d)$ densities ρ , showing that indeed $\rho \nabla \psi \in L^1_{loc}(\mathbb{R}^d)$, where ψ is the velocity potential defined by (1.4). The lemma includes the equivalence of two formulations for the equation that governs stationary states in Definition 1.1.

Lemma 2.1. *Let $\rho \in W^{1,1}(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d) \cap C(\mathbb{R}^d)$ be a nonnegative function. Then $K_s * \rho \in W^{1,1}_{loc}(\mathbb{R}^d)$ and $\nabla(K_s * \rho) = K_s * \nabla \rho$. Moreover, $\psi \in W^{1,1}_{loc}(\mathbb{R}^d)$, where ψ is defined by (1.4). Finally,*

$$\nabla \cdot (\rho \nabla \psi) = 0 \quad \text{in } \mathcal{D}'(\mathbb{R}^d) \quad (2.1)$$

if and only if

$$\begin{aligned} & \int_{\mathbb{R}^d} \nabla(\rho^m + \beta \rho^2) \cdot \nabla \varphi \, dx \\ & + \frac{(d-2s)c_{d,s}\chi}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x-y) |x-y|^{2s-d-2} \rho(x) \rho(y) \, dx \, dy = 0 \end{aligned} \quad (2.2)$$

for any $\varphi \in C_c^\infty(\mathbb{R}^d)$.

We postpone the technical proof of the above lemma to the appendix. Here, we focus on its consequences in relation with the definition of stationary states and on further properties that follow from Definition 1.1.

Remark 2.2. By virtue of Lemma (2.1), we have the equivalence between the stationary version of (1.6) and (2.1). Therefore, Definition 1.1 agrees with the natural definition of steady solutions of the evolution equation in (1.1) in the formulation (1.6), i.e., time independent solutions.

Remark 2.3. Definition 1.1 provides a notion of stationary state which is weaker than the ones previously used in [8–10,12,11,17], see [8, Definition 2.1]. Indeed, the latter definition requires $\rho \nabla \psi = 0$ in $\mathcal{D}'(\mathbb{R}^d)$ along with more regularity properties of ρ . We point out that the $W^{1,1}(\mathbb{R}^d)$ regularity of the density ρ and the $W^{1,\infty}$ regularity of the velocity potential ψ in the interior of the support of ρ do always agree with the properties of the minimizers of the free energy functional \mathcal{F}_s that we shall discuss later on. On the other hand, such minimizers do not match the notion of stationary states from [8, Definition 2.1] if s is small.

Our aim is now to show that if ρ is a steady state, then we can obtain the natural zero-dissipation identity

$$\int_{\mathbb{R}^d} |\nabla \psi|^2 \rho \, dx = 0,$$

for ψ defined by (1.4).

Proposition 2.4. *Assume that ρ is a stationary state for the evolution equation in (1.1), according to Definition 1.1. Then $\int_{\mathbb{R}^d} |\nabla \psi|^2 \rho \, dx = 0$. In particular ψ is constant in each connected component of Ω .*

Proof. Since $\psi \in W_{loc}^{1,1}(\mathbb{R}^d)$ by Lemma 2.1, for any given $\eta \in C_c^\infty(\mathbb{R}^d)$ we have $\eta\psi \in W^{1,1}(\mathbb{R}^d)$, and then there exists a sequence $(\varphi_n)_{n \in \mathbb{N}} \subset C_c^\infty(\mathbb{R}^d)$ such that $\varphi_n \rightarrow \eta\psi$ in $W^{1,1}(\mathbb{R}^d)$ as $n \rightarrow \infty$. Since $\operatorname{div}(\rho \nabla \psi) = 0$ in $\mathcal{D}'(\mathbb{R}^d)$, we deduce that

$$\int_{\Omega} \nabla \varphi_n \cdot \nabla \psi \rho \, dx = 0$$

for every $n \in \mathbb{N}$. Since $\rho \nabla \psi \in L^\infty(\Omega)$, by taking the limit as $n \rightarrow \infty$ we find

$$\int_{\mathbb{R}^d} \nabla \psi \cdot \nabla (\eta\psi) \rho \, dx = 0 \quad (2.3)$$

for every $\eta \in C_c^\infty(\mathbb{R}^d)$. Let us now consider a radially decreasing function $\zeta \in C_c^\infty(\mathbb{R}^d)$ such that $\zeta(x) = 1$ if $|x| \leq 1$ and such that $0 \leq \zeta(x) \leq 1$ for every $x \in \mathbb{R}^d$. For every $k \in \mathbb{N}$, let $\eta_k(x) = \zeta(x/k)$, so that $\zeta_k \rightarrow 1$ pointwise and monotonically on \mathbb{R}^d and $\nabla \eta_k \rightarrow 0$ in $L^\infty(\mathbb{R}^d)$ as $k \rightarrow \infty$. With this choice of the test functions, from (2.3) we get

$$0 = \int_{\mathbb{R}^d} \psi \rho \nabla \psi \cdot \nabla \eta_k \, dx + \int_{\mathbb{R}^d} |\nabla \psi|^2 \eta_k \rho \, dx \quad (2.4)$$

for every $k \in \mathbb{N}$. The first term in the right hand side vanishes as $k \rightarrow +\infty$, since

$$\left| \int_{\mathbb{R}^d} \nabla \psi \cdot \nabla \eta_k \psi \rho \, dx \right| \leq \|\psi\|_{L^\infty(\mathbb{R}^d)} \|\nabla \psi\|_{L^\infty(\Omega)} \|\rho\|_{L^1(\mathbb{R}^d)} \|\nabla \eta_k\|_{L^\infty(\mathbb{R}^d)}$$

and since $\nabla \eta_k \rightarrow 0$ in $L^\infty(\mathbb{R}^d)$. Therefore from (2.4), by applying the monotone convergence theorem to the second term in right hand side, we obtain $\int_{\mathbb{R}^d} |\nabla \psi|^2 \rho \, dx = 0$. Since $\rho > 0$ in Ω , this implies $\nabla \psi = 0$ a.e. in Ω , thus ψ is constant in each connected component of Ω . \square

Stationary states are closely related to minimizers of the free-energy functional \mathcal{F}_s . Indeed, by adapting the arguments of [12], we shall prove that for every $M > 0$ a minimizer of functional \mathcal{F}_s over \mathcal{Y}_M does exist, that it is necessarily radially decreasing and it satisfies a suitable Euler-Lagrange equation. Thanks to these properties, a minimizer of \mathcal{F}_s over \mathcal{Y}_M will be immediately seen to be a stationary state of mass M . In the analysis of minimizers we shall make use of the following notation

$$\begin{aligned}\mathcal{H}_m[\rho] &:= \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m(x) dx + \beta \int_{\mathbb{R}^d} \rho^2(x) dx, \\ \mathcal{W}_s[\rho] &:= -\frac{\chi}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho(x) \rho(y) dx dy,\end{aligned}\tag{2.5}$$

so that $\mathcal{F}_s[\rho] = \mathcal{H}_m[\rho] + \mathcal{W}_s[\rho]$. We notice that since $m > 2$, both $\mathcal{H}_m[\rho]$ and $\mathcal{W}_s[\rho]$ are finite if $\rho \in L^1(\mathbb{R}^d) \cap L^m(\mathbb{R}^d)$, as a consequence of the Hardy-Littlewood-Sobolev inequality (see [20, Theorem 4.3]), that reads (notice that due to the condition $m > 2$ we always have $1 < 2d/(d+2s) < m$)

$$\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho(x) \rho(y) dx dy \leq H_{d,s} \|\rho\|_{L^{\frac{2d}{d+2s}}(\mathbb{R}^d)}^2,\tag{2.6}$$

where the optimal constant $H_{d,s}$ is given by

$$H_{d,s} = \pi^{\frac{d-2s}{2}} \frac{\Gamma(s)}{\Gamma(s+d/2)} \left(\frac{\Gamma(d/2)}{\Gamma(d)} \right)^{-\frac{2s}{d}}.$$

As a direct consequence, we check that for any $\rho \in \mathcal{Y}_M$ we have, for some constant \bar{C} only depending on χ, m, s, d ,

$$\frac{1}{2(m-1)} \|\rho\|_{L^m(\mathbb{R}^d)}^m \leq \mathcal{F}_s[\rho] + \bar{C}.\tag{2.7}$$

Indeed, by the sharp Hardy-Littlewood-Sobolev inequality (2.6), letting

$$S_{d,s} := c_{d,s} H_{d,s} = (4\pi)^{-s} \frac{\Gamma(-s+d/2)}{\Gamma(s+d/2)} \left(\frac{\Gamma(d/2)}{\Gamma(d)} \right)^{-2s/d},\tag{2.8}$$

there holds

$$|\mathcal{W}_s[\rho]| = \frac{\chi}{2} c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho(x) \rho(y) dx dy \leq \frac{\chi}{2} S_{d,s} \|\rho\|_{L^{\frac{2d}{d+2s}}(\mathbb{R}^d)}^2.\tag{2.9}$$

By taking advantage of the interpolation inequality

$$\|\rho\|_{L^{\frac{2d}{d+2s}}(\mathbb{R}^d)}^2 \leq M^{2-2\theta} \|\rho\|_{L^m(\mathbb{R}^d)}^{2\theta}, \quad \theta := \frac{(d-2s)m}{2d(m-1)}$$

and of Young inequality (taking conjugate exponents p, p' with $p = 2\theta$), from (2.9) we deduce

$$\begin{aligned} |\mathcal{W}_s[\rho]| &= \frac{\chi}{2} c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho(x) \rho(y) dx dy \leq \frac{\chi}{2} S_{d,s} M^{2-2\theta} \|\rho\|_{L^m(\mathbb{R}^d)}^{2\theta} \\ &\leq \bar{C} + \frac{1}{2(m-1)} \|\rho\|_{L^m(\mathbb{R}^d)}^m, \end{aligned} \quad (2.10)$$

where $\bar{C} = \bar{C}(\chi, m, s, d)$ is defined by

$$\bar{C} = \frac{m-2\theta}{m} \left(\frac{m}{4\theta(m-1)} \right)^{-\frac{2\theta}{m-2\theta}} \left(\frac{\chi}{2} S_{d,s} M^{2-2\theta} \right)^{\frac{m}{m-2\theta}}, \quad \theta = \frac{(d-2s)m}{2d(m-1)}. \quad (2.11)$$

Then we have

$$\frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m dx \leq \mathcal{F}_s[\rho] + |\mathcal{W}_s[\rho]| \leq \mathcal{F}_s[\rho] + \bar{C} + \frac{1}{2(m-1)} \|\rho\|_{L^m(\mathbb{R}^d)}^m,$$

thus (2.7) holds.

We have the following result about existence of minimizers, which employs the classical concentration compactness theorem of Lions [21].

Lemma 2.5. *The functional \mathcal{F}_s admits a minimizer over \mathcal{Y}_M . If $\rho_s \in \operatorname{argmin}_{\mathcal{Y}_M} \mathcal{F}_s$, then ρ_s is radially decreasing, compactly supported and it satisfies*

$$\frac{m}{m-1} \rho_s^{m-1} + 2\beta \rho_s = (\chi K_s * \rho_s - \mathcal{C}_s)_+ \quad \text{in } \mathbb{R}^d, \quad (2.12)$$

where

$$0 < \mathcal{C}_s := -\frac{2}{M} \mathcal{F}_s[\rho_s] - \frac{1}{M} \frac{m-2}{m-1} \int_{\mathbb{R}^d} \rho_s^m(x) dx. \quad (2.13)$$

In particular, we have

$$\mathcal{F}_s[\rho_s] = -\frac{1}{d-2s} \left(\frac{dm-2d+2s}{m-1} \|\rho_s\|_m^m + 2s\beta \|\rho_s\|_2^2 \right) < 0 \quad (2.14)$$

$$\mathcal{C}_s = \frac{1}{M(d-2s)} \left(\frac{dm+2sm-2d}{m-1} \|\rho_s\|_m^m + 4\beta s \|\rho_s\|_2^2 \right) \quad (2.15)$$

For the proof of the above lemma we follow the concentration compactness argument as applied in Appendix A.1 of [19]. Indeed, the proof is based on [21, Theorem II.1,

Corollary II.1], that are next recalled, denoting by $\mathcal{M}^p(\mathbb{R}^d)$ the Marcinkiewicz space or weak L^p space.

Theorem 2.6. [21, Theorem II.1] *Suppose $K \in \mathcal{M}^p(\mathbb{R}^d)$, $1 < p < \infty$, and consider the problem*

$$I_M = \inf_{\rho \in \mathcal{Y}_{q,M}} \left\{ \int_{\mathbb{R}^d} j(\rho) dx - \frac{\chi}{2} \iint_{\mathbb{R}^d \times \mathbb{R}^d} K(x-y) \rho(x) \rho(y) dx dy \right\}.$$

Here,

$$\mathcal{Y}_{q,M} = \left\{ \rho \in L^q(\mathbb{R}^d) \cap L^1(\mathbb{R}^d), \rho \geq 0 \text{ a.e.}, \int_{\mathbb{R}^d} \rho(x) dx = M \right\}, \quad q = \frac{p+1}{p},$$

and the nonlinearity $j: \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is a strictly convex nonnegative function such that

$$\lim_{t \rightarrow 0^+} \frac{j(t)}{t} = 0, \quad \lim_{t \rightarrow +\infty} \frac{j(t)}{t^q} = +\infty.$$

Then there exists a minimizer of problem (I_M) if the following holds:

$$I_{M_0} < I_M + I_{M_0-M} \quad \text{for all } M \in (0, M_0). \quad (2.16)$$

Proposition 2.7. [21, Corollary II.1] *Suppose there exists some $\lambda \in (0, N)$ such that*

$$K(tx) \geq t^{-\lambda} K(x)$$

for all $t \geq 1$. Then (2.16) holds if and only if

$$I_M < 0 \quad \text{for all } M > 0. \quad (2.17)$$

Proof of Lemma 2.5. The proof is similar to [12, Theorem 5], but we sketch the main lines of the arguments for the sake of completeness. Let $p = \frac{d}{d-2s}$ and $q = \frac{p+1}{p}$. We first notice that our potential $K_s(x) = c_{d,s}|x|^{2s-d}$ is in $\mathcal{M}^p(\mathbb{R}^d)$ and it is clear that K_s verifies the homogeneity assumption in Proposition 2.7 for $\lambda = d - 2s$. Moreover the nonlinearity

$$j(t) = \frac{1}{m-1} t^m + \beta t^2$$

verifies the properties of Theorem 2.6, since $m > 2$ implies $m > q$. Then we just have to show that there exists some density $\rho \in \mathcal{Y}_{q,M}$ such that $\mathcal{F}_s[\rho] < 0$. We fix $R > 0$ and define

$$\rho_*(x) := \frac{dM}{\sigma_d R^d} \mathbb{1}_{B_R}(x),$$

where B_R denotes the ball centered at zero and of radius $R > 0$, and where $\sigma_d = 2\pi^{d/2}/\Gamma(d/2)$ is the surface area of the d -dimensional unit ball. Then

$$\begin{aligned} \mathcal{H}_m[\rho_*] &= \frac{1}{m-1} \int_{\mathbb{R}^d} \rho_*^m dx + \beta \int_{\mathbb{R}^d} \rho^2(x) dx = \frac{(dM)^m \sigma_d^{1-m}}{d(m-1)} R^{d(1-m)} + \beta \frac{(dM)^2 \sigma_d^{-1}}{d} R^{-d}, \\ \mathcal{W}_s[\rho_*] &= -\frac{\chi}{2} \iint_{\mathbb{R}^d \times \mathbb{R}^d} K_s(x-y) \rho_*(x) \rho_*(y) dx dy \\ &= -\chi c_{d,s} \frac{(dM)^2}{2\sigma_d^2 R^{2d}} \iint_{\mathbb{R}^d \times \mathbb{R}^d} |x-y|^{2s-d} \mathbb{1}_{B_R}(x) \mathbb{1}_{B_R}(y) dx dy \\ &\leq -\chi c_{d,s} \frac{(dM)^2}{2\sigma_d^2 R^{2d}} (2R)^{2s-d} \frac{\sigma_d^2}{d^2} R^{2d} = -\chi c_{d,s} 2^{2s-d-1} M^2 R^{2s-d} < 0. \end{aligned}$$

Therefore we find

$$\mathcal{F}_s[\rho_*] \leq \frac{M^m d^{m-1} \sigma_d^{1-m}}{(m-1)} R^{d(1-m)} + \beta d M^2 \sigma_d^{-1} R^{-d} - \chi c_{d,s} 2^{2s-d-1} M^2 R^{2s-d}.$$

Since $m > 2$ we have $d(1-m) < -d < 2s-d$, then we can choose $R > 0$ large enough such that $\mathcal{F}_s[\rho_*] < 0$, and hence condition (2.17) is verified. Then Proposition 2.7 and Theorem 2.6 implies that there exists a minimizer ρ_s of \mathcal{F}_s in $\mathcal{Y}_{q,M}$.

Now, the Hardy-Littlewood-Sobolev inequality (2.6) implies that $\rho_s \in L^m(\mathbb{R}^d)$. Moreover, the Riesz rearrangement inequality [20, Theorem 3.7] yields that ρ_s is radially decreasing (then $K_s * \rho_s$ is also radially decreasing and vanishing at infinity).

The fact that ρ_s satisfies (2.12) can be obtained by computing the first variation of \mathcal{F}_s at ρ_s , as done in [14, Theorem 3.1]. Then (2.13) follows by taking into account the equation (2.12) satisfied by ρ_s , multiplying it by ρ_s and integrating over \mathbb{R}^d . Since \mathcal{C}_s in (2.12) is positive, and since both ρ_s and $K_s * \rho_s$ are radially decreasing and vanishing at infinity, from (2.12) we deduce that ρ_s is compactly supported. In particular $x\rho_s \in L^1(\mathbb{R}^d)$ and $\rho_s \in \mathcal{Y}_M$.

Eventually we prove (2.14) and (2.15). Using the mass invariant dilation $\rho_{s,\lambda}(x) = \lambda^d \rho_s(\lambda x)$ we easily find

$$h(\lambda) := \mathcal{F}_s[\rho_{s,\lambda}] = \frac{1}{m-1} \lambda^{d(m-1)} \|\rho_s\|_m^m + \beta \lambda^d \|\rho_s\|_2^2 + \lambda^{d-2s} \mathcal{W}_s[\rho_s].$$

By the minimality of ρ_s , the function $h(\lambda)$ has its unique minimizer at $\lambda = 1$, which implies by differentiation

$$\|\rho_s\|_m^m + \beta \|\rho_s\|_2^2 = -\frac{d-2s}{d} \mathcal{W}_s[\rho_s].$$

Then we have (2.14) hence substituting in the expression of \mathcal{C}_s we find (2.15). \square

Concerning the regularity of the minimizers, we have the following result

Lemma 2.8. *Let $\beta > 0$. Any minimizer of \mathcal{F}_s over \mathcal{Y}_M is bounded and Lipschitz continuous in \mathbb{R}^d and it is C^∞ in the interior of its support.*

Proof. The boundedness of the minimizers follows by [12, Theorem 7] up to minor modifications, nevertheless in the case $s \in (0, 1)$ we propose here a more direct proof based on [15, Proposition 2.8]. Indeed, let $\rho_0 \in \mathcal{Y}_M$ be any minimizer of \mathcal{F}_s , which is radially decreasing and compactly supported by Lemma 2.5. Then (2.12) implies that the Riesz potential $v_0 := K_s * \rho_0$ is a radially decreasing, vanishing at infinity solution to the fractional PDE

$$(-\Delta)^s v = \mathbf{g}(v), \quad (2.18)$$

where the nonlinearity \mathbf{g} is defined as $\mathbf{g}(t) = f^{-1}((t - C_s)_+)$, being f^{-1} the inverse function of the convex nonlinearity $f(t) = \frac{m}{m-1} t^{m-1} + 2\beta t$, $t \geq 0$. Since $\beta > 0$, f^{-1} is Lipschitz continuous therefore so is \mathbf{g} . Now, since $\rho_0 \in L^m(\mathbb{R}^d) \cap L^1(\mathbb{R}^d)$ with $m > 2$, we have $\rho_0 \in L^{2d/(d+2s)}(\mathbb{R}^d)$ thus the Hardy-Littlewood-Sobolev inequality implies that $v_0 \in L^{2_s^*}(\mathbb{R}^d)$, being $2_s^* = 2d/(d - 2s)$. Since $\mathbf{g}(v_0) \leq (2\beta)^{-1}v_0$, we find in particular $\mathbf{g} \circ v \in L^{2d/(d+2s)}(\mathbb{R}^d)$: indeed, if $q = 2d/(d + 2s)$, by Hölder inequality and the radial monotonicity there is a constant $C \geq 0$ such that

$$\int_{\mathbb{R}^d} \mathbf{g}(v_0)^q dx = \int_{\{v_0 > C_s\}} \mathbf{g}(v_0)^q dx \leq C \int_{\{v_0 > C_s\}} v_0^q \leq C \left(\int_{\mathbb{R}^d} v_0^{2_s^*} dx \right)^{\frac{d-2s}{d+2s}}.$$

Then we can proceed exactly as in [15, Proposition 2.8] in order to conclude that $v_0 \in L^\infty(\mathbb{R}^d)$, which implies in turns by (2.12) that $\rho_0 \in L^\infty(\mathbb{R}^d)$.

Now we turn to the regularity properties of ρ_0 and we refer mainly to [12, Theorem 8]. Let us first consider the easiest case $s > 1/2$. We have in this case $u_0 \in W^{1,\infty}(\mathbb{R}^d)$ from [12, Lemma 1], thus by (2.18) we have

$$\rho_0 = \mathbf{g}(u_0), \quad (2.19)$$

therefore (since \mathbf{g} is Lipschitz) $\rho_0 \in W^{1,\infty}(\mathbb{R}^d)$. If $s \in (0, 1/2)$, by the first part of the proof of [12, Theorem 8] it follows that $u_0 \in C^{0,\gamma}$ for any $\gamma < 2s$, then equation (2.19) gives $\rho_0 \in C^{0,\gamma}$ for any $\gamma < 2s$. By the Hölder regularity of the Riesz potential (see again [12, Eq. 3.24]) we have $u_0 \in C^{0,\gamma}$ for any $\gamma < 4s$. Bootstrapping, we finally find $u_0, \rho_0 \in W^{1,\infty}(\mathbb{R}^d)$. The case $s = 1/2$ can be treated analogously, see again the proof of [12, Theorem 8]. Finally, [12, Theorem 10] gives the smoothness of ρ_0 inside its support. \square

Remark 2.9. The case $\beta = 0$ of Lemma 2.8 is treated in [12, Theorem 8, Remark 2]. As shown therein, if $\beta = 0$ any minimizer in \mathcal{Y}_M of \mathcal{F}_s is bounded, smooth inside its support, and enjoys suitable Hölder regularity on \mathbb{R}^d depending on m, s .

It is now easy to check that minimizers of functional \mathcal{F}_s over \mathcal{Y}_M , whose existence is ensured by Lemma 2.5, are stationary states.

Proposition 2.10. *Let $\beta \geq 0$. If ρ_s minimizes \mathcal{F}_s over \mathcal{Y}_M , then ρ_s is a stationary state according to Definition 1.1.*

Proof. By Lemma 2.5, Lemma 2.8 and Remark 2.9, ρ_s is continuous, compactly supported, radially decreasing and smooth on B_R , where R is the radius of its support. Therefore its radial profile is absolutely continuous in $[0, R]$, and belongs to the weighted Sobolev space $W^{1,1}((0, R), r^{d-1} dr)$. This implies $\rho_s \in W^{1,1}(B_R)$, see [16, Theorem 2.3]. Since ρ_s is also vanishing on $\mathbb{R}^d \setminus B_R$, we conclude that it belongs to $W^{1,1}(\mathbb{R}^d)$. By Lemma 2.1, we get $\psi \in W_{loc}^{1,1}(\mathbb{R}^d)$, where ψ is defined by (1.4), thus $\rho \nabla \psi \in L^1(\mathbb{R}^d)$. But the validity of (2.12) implies that ψ is constant on B_R . In particular $\rho \nabla \psi = 0$ a.e. in \mathbb{R}^d and (2.1) follows. \square

The complete characterization of stationary states (according to Definition 1.1) is finally given by the next theorem.

Theorem 2.11. *Let $\beta > 0$ (resp. $\beta = 0$). For any mass $M > 0$, there exists a unique stationary state (resp. a unique radial stationary state) of mass M and center of mass 0. Such steady state is radially decreasing, compactly supported, Lipschitz on \mathbb{R}^d (resp. Hölder on \mathbb{R}^d) and C^∞ in the interior of its support. Moreover, it coincides (up to translation) with the unique minimizer of the energy functional \mathcal{F}_s in the class \mathcal{Y}_M .*

Proof. Let $\beta > 0$. Let ρ be a stationary state according to Definition 1.1. Let $v := K_s * \rho$. Besides $v \in W_{loc}^{1,1}(\mathbb{R}^d)$, which is proven in Lemma 2.1, by using Sobolev embeddings it is possible to prove that exists $\gamma \in (0, 1)$ such that $v \in C^{0,\gamma}(\mathbb{R}^d)$ (see again [12, Theorem 8]).

Let as usual $\Omega = \{\rho > 0\}$. Since ρ is continuous, Ω is open and thus $\Omega = \cup_{n=1}^\infty \Theta_n$, where the Θ_n 's are the countably many open connected components of Ω . Let us introduce the continuous functions

$$\rho_n(x) := \begin{cases} \rho(x) & \text{if } x \in \Theta_n \\ 0 & \text{if } x \in \mathbb{R}^d \setminus \Theta_n. \end{cases} \quad (2.20)$$

From Proposition 2.4, for each $n \in \mathbb{N}$ there is a constant $Q_n \in \mathbb{R}$ such that there holds

$$f(\rho_n(x)) = f(\rho(x)) = \chi v(x) - Q_n \quad \text{for every } x \in \Theta_n,$$

where $f(t) := \frac{m}{m-1}t^{m-1} + 2\beta t$, $t \geq 0$. The constant Q_n is nonnegative, since ρ and v are nonnegative continuous and ρ vanishes on $\partial\Theta_n$. Therefore we have

$$\rho_n(x) = \mathfrak{g}(\chi v(x) - Q_n) \quad \text{for every } x \in \overline{\Theta_n},$$

where $\mathfrak{g} : [0, +\infty) \rightarrow [0, +\infty)$ is the inverse function of f , and we notice that \mathfrak{g} is Lipschitz on $[0, +\infty)$ since $\beta > 0$. Hence, we have $\rho_n \in C^{0,\gamma}(\overline{\Theta_n})$, and (2.20) implies $\rho_n \in C^{0,\gamma}(\mathbb{R}^d)$. We stress that the Hölder constant of ρ_n , that we denote by c , is independent of n , since it only depends on the Hölder constant of v and the Lipschitz constant of f . Now, for every two distinct points $x \in \mathbb{R}^d$ and $y \in \mathbb{R}^d$, there exist $n \in \mathbb{N}$ and $m \in \mathbb{N}$ such that $\rho(x) = \rho_n(x)$ and $\rho(y) = \rho_m(y)$ and

$$|\rho(x) - \rho(y)| = |\rho_n(x) - \rho_m(y)| \leq |\rho_n(x) - \rho_n(y)| + |\rho_m(x) - \rho_m(y)| \leq 2c|x - y|^\gamma,$$

so that $\rho \in C^{0,\gamma}(\mathbb{R}^d)$. Since $v = K_s * \rho$, the Hölder estimates from [29] entail $v \in C^{0,\gamma+2s-\varepsilon}(\mathbb{R}^d)$ for every arbitrarily small ε . Then we may bootstrap this argument and in a finite number of steps we get $v \in C^{0,1}(\mathbb{R}^d)$ and $\rho \in C^{0,1}(\mathbb{R}^d)$. Since $m > 2$ and ρ is bounded, we conclude that ρ^{m-1} is Lipschitz on \mathbb{R}^d as well. We recall that by Proposition (2.4) the velocity potential ψ defined by (1.4) is constant on each connected component of Ω .

Now, an easy modification of [12, Theorem 3], which crucially exploits the Lipschitz regularity of ρ^{m-1} , shows the radially of the steady state ρ : indeed, the actual entropy

$$\mathcal{H}_m[\rho_*] = \frac{1}{m-1} \int_{\mathbb{R}^d} \rho_*^m dx + \beta \int_{\mathbb{R}^d} \rho_*^2(x) dx$$

decreases under the modified continuous Steiner symmetrization introduced in the proof therein.

All the regularity properties of the steady states can be argued by the proof of Lemma 2.8. The uniqueness follows by the results of [17], namely by the Remark under the statement of [17, Remark 1.2], because the nonlinearity $\Phi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$, defined by $\Phi(\rho) = \frac{1}{m-1}\rho^m + \beta\rho^2$ entering in the nonlinear diffusion is a strictly increasing smooth convex function. Finally, the identification (up to translation) with the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M follows from Proposition 2.10.

Finally, if $\beta = 0$, we only consider the class of radial stationary states, and the uniqueness in this class follows from the results in [17]. Again the unique radial stationary state of mass M and center of mass 0 is the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M , see Proposition 2.10, and the other properties follow from Lemma 2.5 and [12, Theorem 8]. \square

Remark 2.12. If $\beta = 0$ the general uniqueness result of Theorem 2.11 (without assuming radially) is still an open problem, mainly due to the fact that the inverse function of $f(t) = t^{m-1}$ is $C^{1/(m-1)}$ for $m > 2$ but not Lipschitz, preventing to obtain that ρ^{m-1}

is Lipschitz for a stationary state ρ , which is crucial for applying the radially result of [11]. However, we still have radially of every stationary states, thus uniqueness, as long as we assume in addition that $m < m^*$, because in this case ρ^{m-1} can be proven to be Lipschitz, by reasoning on connected components as done in the previous proof and by applying the bootstrap argument from [12, Theorem 8]. Here, $m^* := \frac{2-2s}{1-2s}$ if $s < 1/2$ and $m^* := +\infty$ otherwise (so that there is no restriction if $d \geq 2$ and $s \geq 1/2$).

Proof of Theorem 1.1. The result follows from Theorem 2.11 and Remark 2.12. \square

3. Asymptotic behavior of stationary states as $s \rightarrow 0$

We investigate the asymptotic behavior of ρ_s as s approaches 0, where ρ_s is, for each small s , the unique stationary state of given mass M and center of mass 0 for the equation in (1.1), provided by Theorem 2.11. Thanks to the identifications with minimizers of the free-energy, this will be done by showing that functionals \mathcal{F}_s Γ -converge to the limit energy functional \mathcal{F}_0 defined by

$$\mathcal{F}_0[\rho] := \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m(x) dx - \frac{\chi-2\beta}{2} \int_{\mathbb{R}^d} \rho^2(x) dx, \quad (3.1)$$

whose minimization is governed by the following proposition.

Proposition 3.1 (Minimization of the limit functional \mathcal{F}_0). Suppose that $0 \leq \beta < \chi/2$. Then functional \mathcal{F}_0 admits a unique radially decreasing minimizer over \mathcal{Y}_M , given by

$$\rho_0(x) := \left(\frac{\chi-2\beta}{2} \right)^{1/(m-2)} \mathbb{1}_{B_{R_0}}(x), \quad \text{where } R_0 = \left(\frac{dM}{\sigma_d} \right)^{1/d} \left(\frac{\chi-2\beta}{2} \right)^{-\frac{1}{d(m-2)}}. \quad (3.2)$$

Else if $\beta \geq \chi/2$, functional \mathcal{F}_0 does not admit a minimizer over \mathcal{Y}_M and $\inf_{\mathcal{Y}_M} \mathcal{F}_0 = 0$.

Proof. Through the proof, we let for simplicity $\gamma := -\beta + \chi/2$. For every $\rho \in \mathcal{Y}_M$ and every $\lambda > 0$, letting $\rho_\lambda(x) := \lambda^d \rho(\lambda x)$, $x \in \mathbb{R}^d$, direct computations show that

$$\begin{aligned} \mathcal{F}_0[\rho_\lambda] &= \frac{1}{m-1} \lambda^{d(m-1)} \int_{\mathbb{R}^d} \rho^m(x) dx - \gamma \lambda^d \int_{\mathbb{R}^d} \rho^2(x) dx, \\ \frac{d}{d\lambda} \mathcal{F}_0[\rho_\lambda] &= d\lambda^{d-1} \left(\lambda^{d(m-2)} \int_{\mathbb{R}^d} \rho^m(x) dx - \gamma \int_{\mathbb{R}^d} \rho^2(x) dx \right). \end{aligned} \quad (3.3)$$

Assume that $0 \leq \beta < \chi/2$. In this case, by (3.3) the map $(0, +\infty) \ni \lambda \mapsto \mathcal{F}_0(\rho_\lambda)$ is uniquely minimized at

$$\lambda = \lambda_* := \left(\gamma \int_{\mathbb{R}^d} \rho^2(x) dx \right)^{\frac{1}{d(m-2)}} \left(\int_{\mathbb{R}^d} \rho^m(x) dx \right)^{-\frac{1}{d(m-2)}}$$

with value

$$\mathcal{F}_0[\rho_{\lambda_*}] = \frac{2-m}{m-1} \gamma^{\frac{m-1}{m-2}} \left(\int_{\mathbb{R}^d} \rho^2(x) dx \right)^{\frac{m-1}{m-2}} \left(\int_{\mathbb{R}^d} \rho^m(x) dx \right)^{-\frac{1}{m-2}}.$$

But writing

$$2 = \frac{m-2}{m-1} + \frac{m}{m-1},$$

Hölder inequality with exponents $p = (m-1)/(m-2)$, $p' = m-1$ yields

$$\left(\int_{\mathbb{R}^d} \rho^2(x) dx \right)^{m-1} \leq M^{m-2} \int_{\mathbb{R}^d} \rho^m(x) dx,$$

thus for every $\rho \in \mathcal{Y}_M$ there holds

$$\mathcal{F}_0[\rho] \geq \frac{2-m}{m-1} \gamma^{\frac{m-1}{m-2}} \left(\int_{\mathbb{R}^d} \rho^2(x) dx \right)^{\frac{m-1}{m-2}} \left(\int_{\mathbb{R}^d} \rho^m(x) dx \right)^{-\frac{1}{m-2}} \geq M \frac{2-m}{m-1} \gamma^{\frac{m-1}{m-2}}. \quad (3.4)$$

Therefore, if there is a density ρ achieving the constant at the right-hand side of (3.4), then ρ is a minimizer of \mathcal{F}_0 . However, the above Hölder inequality is an equality if and only if ρ is a multiple of a characteristic function, i.e., $\rho(x) = t \mathbb{1}_{\Omega}(x)$ for some measurable subset Ω of \mathbb{R}^d , and the condition $\int_{\mathbb{R}^d} \rho(x) dx = M$ implies $|\Omega| > 0$ and $t = M|\Omega|^{-1}$. In particular, the second inequality in (3.4) is an equality if and only if $\rho(x) = M|\Omega|^{-1} \mathbb{1}_{\Omega}(x)$. On the other hand, the first inequality in (3.4) is an equality if and only if $\rho = \rho_{\lambda_*}$, i.e., $\lambda_* = 1$, which means $\gamma \int_{\mathbb{R}^d} \rho^2(x) dx = \int_{\mathbb{R}^d} \rho^m(x) dx$, and this condition, in case $\rho(x) = M|\Omega|^{-1} \mathbb{1}_{\Omega}(x)$, entails $|\Omega| = M\gamma^{-\frac{1}{m-2}}$. We conclude that both inequalities in (3.4) are equalities if and only if $\rho(x) = M|\Omega|^{-1} \mathbb{1}_{\Omega}(x)$ for some measurable set $\Omega \subset \mathbb{R}^d$ such that $|\Omega| = M\gamma^{-\frac{1}{m-2}}$, implying that this family of functions coincides with $\operatorname{argmin}_{\mathcal{Y}_M} \mathcal{F}_0$ up to translations. In this family there is a unique radially decreasing function ρ_0 , obtained by letting $\Omega = B_R$ and by choosing R such that $|\Omega| = M\gamma^{-\frac{1}{m-2}}$. Hence, ρ_0 is given by (3.2).

The statement concerning the case $\beta \geq \chi/2$ (i.e., $\gamma \leq 0$) follows by letting $\lambda \rightarrow 0$ in (3.3). \square

The steady states ρ_s in one dimension for different values of s are shown in Fig. 1, with $m = 3$, $\chi = 1$, obtained by iterating the governing equation (2.12), i.e.,

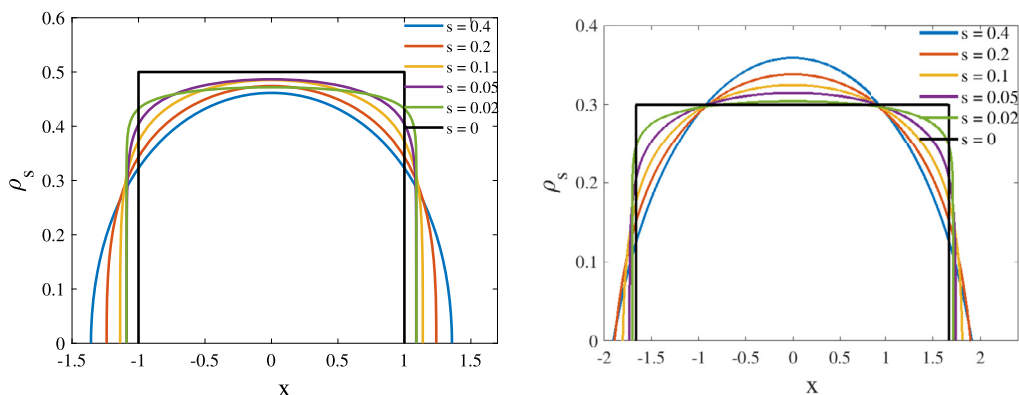


Fig. 1. The steady states for different $s > 0$ with $m = 3$ and $\chi = 1$ (Left figure: $\beta = 0$ and Right figure: $\beta = 0.2$). The expected limiting steady state with $s = 0$, which is a characteristic function with height $\left(\frac{\chi - 2\beta}{2}\right)^{1/(m-2)}$ is also plotted for reference.

$$\frac{m}{m-1} (\tilde{\rho}_s^{(new)})^{m-1} + 2\beta \tilde{\rho}_s^{(new)} = (\chi K_x * \rho_s^{(old)} - C_s^{(old)})_+,$$

with $C_s^{(old)}$ given by Eq. (2.13), followed by a spatial scaling $\rho_s^{(new)}(x) = \tilde{\rho}_s^{(new)}(\lambda^{(new)}x)$ such that the total mass is exactly M . The convergence towards the limit ρ_0 in Proposition 3.1 as s goes to zero is illustrated, for both $\beta = 0$ (left figure) and $\beta > 0$ (right figure). The regularizing effect with $\beta > 0$ is obvious, especially near the boundary of the support.

The next step is the investigation of the behavior of \mathcal{F}_s on characteristic functions.

Lemma 3.2. *Let $s \in (0, 1)$ ($s < 1/2$ if $d = 1$). For any $R > 0$, let $\mathcal{Y}_M \ni \rho_R := \frac{dM}{\sigma_d R^d} \mathbb{1}_{B_R}$, where B_R is the ball of radius R , centered at the origin. Assume that $0 \leq \beta < \chi/2$. Then there exists a unique positive number R_s such that*

$$\mathcal{F}_s[\rho_{R_s}] = \min_{R>0} \mathcal{F}_s[\rho_R].$$

In particular, for $\beta = 0$ its value is

$$R_s = \left(\frac{2\sqrt{\pi} M^{m-2} \Gamma(s+1) \Gamma(\frac{d}{2} + s + 1)}{\chi \omega_d^{m-2} \Gamma(s + \frac{1}{2}) \Gamma(\frac{d}{2} - s + 1)} \right)^{\frac{1}{2s+d(m-2)}}.$$

Moreover, the map $(0, 1/2) \ni s \mapsto \mathcal{F}_s[\rho_{R_s}]$ is continuous, it has negative value for any $s \in (0, 1/2)$ and there holds

$$\lim_{s \downarrow 0} \mathcal{F}_s[\rho_{R_s}] = -M \frac{m-2}{m-1} \left(\frac{\chi - 2\beta}{2} \right)^{\frac{m-1}{m-2}}. \quad (3.5)$$

In particular, the value (3.5) is exactly $\mathcal{F}_0(\rho_{R_0})$, where R_0 is given by (3.2).

Proof. By the proof of Lemma 2.5 we argue that for large R we have $\mathcal{F}_s[\rho_R] < 0$. Moreover, by the proof of [15, Lemma 6.1] one has

$$\mathcal{F}_s[\rho_R] = C_1 R^{-\alpha} + C_2 R^{-\gamma} - C_3 R^{-\delta}$$

where

$$\alpha = d(m-1), \gamma = d, \delta = d-2s, \alpha > \gamma > \delta$$

and

$$C_1 = \frac{(dM)^m \sigma_d^{1-m}}{d(m-1)}, \quad C_2 = \beta \frac{(dM)^2 \sigma_d^{-1}}{d}, \quad C_3 = \frac{\chi d^2 M^2 \Gamma(s + \frac{1}{2}) \Gamma(\frac{d}{2} - s)}{4\sqrt{\pi} \sigma_d \Gamma(s+1) \Gamma(\frac{d}{2} + s + 1)}.$$

Since $m > 2$ and $\beta < \chi/2$, it is then readily seen that for all $s \in [0, 1/2)$, the map $(0, +\infty) \ni R \mapsto \mathcal{F}_s[\rho_R]$ admits a unique minimizer R_s , which is the unique solution to the algebraic equation

$$\mathfrak{h}(s, R) := -\alpha C_1 - \gamma C_2 R^{\alpha-\gamma} + \delta C_3 R^{\alpha-\delta} = 0. \quad (3.6)$$

Due to the structure of the coefficients, observe that the map $\mathfrak{h}(\cdot, R)$ is continuous for $s \in [0, 1/2)$. We claim that the map $s \mapsto R_s$ is continuous up to $s = 0$ in the interval $(0, 1/2)$. To prove the claim, for any $s_0 \in [0, 1/2)$, let (s_0, R_{s_0}) be the unique solution to (3.6). Then we have

$$\frac{\partial \mathfrak{h}}{\partial R}(s_0, R_0) = \frac{1}{R_0} ((\alpha - \delta) \delta C_3 R_0^{\alpha-\delta} - \gamma C_2 (\alpha - \gamma) R_0^{\alpha-\gamma}).$$

Thus, using that $\mathfrak{h}(s_0, R_0) = 0$,

$$\frac{\partial \mathfrak{h}}{\partial R}(s_0, R_0) = \frac{1}{R_0} ((\gamma - \delta) \gamma C_2 R_0^{\alpha-\gamma} + \alpha (\alpha - \delta) C_1) > 0.$$

Then the Implicit Function Theorem assures that the map $s \rightarrow R_s$ is continuous in a neighborhood of any point s_0 and the claim follows. This implies in particular that the map $s \rightarrow \mathcal{F}_s(\rho_{R_s})$ is continuous at $s = 0$ and we have

$$\lim_{s \downarrow 0} \mathcal{F}_s[\rho_{R_s}] = \mathcal{F}_0(\rho_{R_0}),$$

then the only issue is to compute the value R_0 , which is obtained by equation (3.6) letting $s \rightarrow 0$. Such computation shows that the value R_0 is the one in (3.2). Then, inserting ρ_{R_0} in the expression (3.1) of \mathcal{F}_0 we have

$$\mathcal{F}_0[\rho_{R_0}] = -M \frac{m-2}{m-1} \left(\frac{\chi - 2\beta}{2} \right)^{\frac{m-1}{m-2}},$$

as desired. \square

Remark 3.3. Since the limit value as $s \rightarrow 0$ of \mathcal{F}_s in ρ_R is given by

$$\mathcal{F}_0[\rho_R] = \frac{(dM)^m}{d(m-1)} \sigma_d^{1-m} R^{-d(m-1)} + \frac{dM^2}{\sigma_d} \left(\beta - \frac{\chi}{2} \right) R^{-d}$$

and $m > 2$, we observe that $R \mapsto \mathcal{F}_0[\rho_R]$ does not admit a minimum for $\beta \geq \chi/2$.

Next we investigate some asymptotic properties of minimizers as $s \downarrow 0$.

Lemma 3.4. Fix any $s_0 \in (0, 1/2)$. For any $s \in (0, s_0)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . Then $\sup_{s \in (0, s_0)} \|\rho_s\|_{L^\infty(\mathbb{R}^d)} < +\infty$.

Proof. Since ρ_s is continuous and radially decreasing by Lemma 2.5 and Lemma 2.8, we have $\|\rho_s\|_{L^\infty(\mathbb{R}^d)} = \rho_s(0)$. By (2.12) and (2.13) we have, letting B_1 denote the unit ball centered at the origin,

$$\begin{aligned} \frac{m}{m-1} \rho_s(0)^{m-1} + 2\beta \rho_s(0) &= \chi c_{d,s}(|\cdot|^{2s-d} * \rho_s)(0) - \mathcal{C}_s \leq \chi c_{d,s} \int_{\mathbb{R}^d} |y|^{2s-d} \rho_s(y) dy \\ &\leq \chi c_{d,s} \rho_s(0) \int_{B_1} |y|^{2s-d} dy + \chi c_{d,s} \int_{\mathbb{R}^d \setminus B_1} \rho_s(y) dy \\ &\leq \chi c_{d,s} \rho_s(0) \sigma_d \int_0^1 r^{2s-1} dr + \chi c_{d,s} M \\ &\leq \frac{\chi c_{d,s} \sigma_d}{2s} \rho_s(0) + \chi c_{d,s} M. \end{aligned}$$

Therefore $\rho_s(0)^{m-1} \leq a \rho_s(0) + b$ for any $s \in (0, s_0)$, where

$$a := \frac{m-1}{m} \sup_{s \in (0, s_0)} \frac{\chi c_{d,s} \sigma_d}{2s} < +\infty \quad \text{and} \quad b := \frac{m-1}{m} \sup_{s \in (0, s_0)} \chi c_{d,s} M < +\infty.$$

Notice that $a < +\infty$ follows from $\lim_{s \downarrow 0} c_{d,s}/s = \pi^{-d/2} \Gamma(d/2)$, see (1.2). Since $m > 2$, we conclude that $\rho_s(0) \leq \bar{x}$ for any $s \in (0, s_0)$, where \bar{x} is the unique positive number such that $\bar{x}^{m-1} = a\bar{x} + b$. \square

Lemma 3.5. Let $0 \leq \beta < \chi/2$. For any $s \in (0, 1/2)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . Then

$$\liminf_{s \downarrow 0} \mathcal{C}_s \geq \frac{m-2}{m-1} \left(\frac{\chi - 2\beta}{2} \right)^{\frac{m-1}{m-2}}$$

where \mathcal{C}_s is defined in (2.15).

Proof. By (2.15) and (2.14) we obtain

$$\begin{aligned} \mathcal{C}_s &= \frac{1}{M(d-2s)} \left(\frac{dm-2d+2s}{m-1} \|\rho_s\|_m^m + 4\beta s \|\rho_s\|_2^2 + 2s \|\rho_s\|_m^m \right) \\ &> \frac{1}{M(d-2s)} \left(\frac{dm-2d+2s}{m-1} \|\rho_s\|_m^m + 2\beta s \|\rho_s\|_2^2 \right) = -\frac{1}{M} \mathcal{F}_s(\rho_s) \end{aligned}$$

and the result follows by the minimality of ρ_s and Lemma (3.2). \square

Lemma 3.6. *Let $0 \leq \beta < \chi/2$. For any $s \in (0, 1/2)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . Then there exists $R \in (0, +\infty)$ and $s_0 \in (0, 1/2)$ such that $\text{supp}(\rho_s) \subset B_R$ for any $s \in (0, s_0)$.*

Proof. We a slight abuse of notation we still denote by ρ_s the radial profile of ρ_s and we notice that since ρ_s is radially non-increasing there holds for any $R > 0$

$$M \geq \int_{B_R} \rho_s(x) dx \geq \int_{B_R} \rho_s(R) dx = \frac{1}{d} \rho_s(R) \sigma_d R^d,$$

which entails

$$\rho_s(x) \leq \frac{dM}{\sigma_d |x|^d} \quad \text{for any } x \in \mathbb{R}^d \setminus \{0\}. \quad (3.7)$$

Let $\text{supp}(\rho_s) =: B_{R_s}$. From (2.12) we deduce

$$\chi c_{d,s}(|\cdot|^{2s-d} * \rho_s)(R_s) = \mathcal{C}_s \quad \text{for any } s \in (0, s_0). \quad (3.8)$$

Notice that by (3.7) we have if $R_s > 1$

$$\begin{aligned} \chi c_{d,s}(|\cdot|^{2s-d} * \rho_s)(R_s) &\leq \chi c_{d,s} \int_{B_1} |y|^{2s-d} \rho_s(R_s - y) dy + \chi c_{d,s} \int_{B_1^c} |y|^{2s-d} \rho_s(R_s - y) dy \\ &\leq \chi c_{d,s} \frac{dM}{\sigma_d} \int_{B_1} \frac{|y|^{2s-d}}{|R_s - y|^d} dy + M \chi c_{d,s} \leq \chi c_{d,s} \frac{dM}{2s|R_s - 1|^d} + M \chi c_{d,s}. \end{aligned}$$

Assuming by contradiction that

$$\limsup_{s \downarrow 0} R_s = +\infty,$$

the above computation shows, recalling from (1.2) that $c_{d,s}/s$ is bounded on $(0, s_0)$, that

$$\liminf_{s \downarrow 0} \chi c_{d,s}(|\cdot|^{2s-d} * \rho_s)(R_s) = 0.$$

This contradicts (3.8), since $\liminf_{s \downarrow 0} \mathcal{C}_s > 0$ by Lemma 3.5. \square

Lemma 3.7. *Let $0 \leq \beta < \chi/2$. For any $s \in (0, 1/2)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . For any vanishing sequence $(s_n) \subset (0, 1/2)$, the sequence (ρ_{s_n}) admits limit points in the strong $L^p(\mathbb{R}^d)$ topology as $n \rightarrow +\infty$ for any $p \in [1, +\infty)$.*

Proof. We have $\rho_s \in W^{1,1}(\mathbb{R}^d)$, by reasoning as done in the proof of Proposition 2.10. We still denote by ρ_s the radial profile of ρ_s and we notice that $\nabla \rho_s(0) = 0$ and $\nabla \rho_s(x) = \rho'_s(|x|) \frac{x}{|x|}$ for $x \neq 0$. Let $R > 1$ and $s_0 \in (0, 1/2)$ such that $\text{supp}(\rho_s) \subset B_R$ for any $s \in (0, s_0)$. The existence of such R, s_0 is due to Lemma 3.6. We have

$$\int_{\mathbb{R}^d} |\nabla \rho_s(x)| dx = \int_{\mathbb{R}^d} |\rho'_s(|x|)| dx = -\sigma_d \int_0^R \rho'_s(r) r^{d-1} dr$$

If $d = 1$ we have $\sigma_d = 2$ and $\int_0^R \rho'(r) dr = -\rho_s(0)$. If $d \geq 2$ we have

$$\rho'_s(r) r^{d-1} = (\rho_s(r) r^{d-1})' - (d-1) \rho_s(r) r^{d-2},$$

thus

$$\begin{aligned} - \int_0^R \rho'_s(r) r^{d-1} dr &= - \int_0^R (\rho_s(r) r^{d-1})' dr + (d-1) \int_0^R \rho_s(r) r^{d-2} dr \\ &= (d-1) \int_0^R \rho_s(r) r^{d-2} dr \\ &\leq (d-1) \left(\int_0^1 \rho_s(r) dr + \int_1^R \rho_s(r) r^{d-1} dr \right) \leq (d-1)(\rho_s(0) + M). \end{aligned}$$

Therefore we always have

$$\sup_{s \in (0, s_0)} \int_{\mathbb{R}^d} |\nabla \rho_s| \leq d\sigma_d \sup_{s \in (0, s_0)} (\rho_s(0) + M) < +\infty,$$

where the finiteness is due to Lemma 3.4. The above uniform $BV(\mathbb{R}^d)$ estimate and the usual compact embedding $BV(\mathbb{R}^d)$ in $L^1_{loc}(\mathbb{R}^d)$, entails the strong sequential $L^1_{loc}(\mathbb{R}^d)$ compactness of the family (ρ_s) , which can be extended to the whole $L^1(\mathbb{R}^d)$ by the tightness due to Lemma 3.6. If ρ is a limit point along a vanishing sequence s_n , we also have $\rho \in L^\infty(\mathbb{R}^d)$ and $\rho_{s_n} \rightarrow \rho$ strongly in $L^p(\mathbb{R}^d)$ for any $p \in [1, +\infty)$, since the sequence (ρ_{s_n}) is also equi-bounded by Lemma 3.4. \square

Lemma 3.8. *Suppose that $\rho_s \in \mathcal{Y}_M$ for any $s > 0$ and that $\rho \in \mathcal{Y}_M$. If $\rho_s \rightarrow \rho$ strongly in $L^2(\mathbb{R}^d)$ as $s \downarrow 0$, then*

$$\lim_{s \downarrow 0} \int_{\mathbb{R}^{2d}} c_{d,s} |x - y|^{2s-d} \rho_s(x) \rho_s(y) dx dy = \int_{\mathbb{R}^d} \rho^2(x) dx.$$

Proof. By Plancherel theorem we have

$$\begin{aligned} & \left| \int_{\mathbb{R}^{2d}} c_{d,s} |x - y|^{2s-d} \rho_s(x) \rho_s(y) dx dy - \int_{\mathbb{R}^d} \rho^2(x) dx \right| \\ &= \frac{1}{(2\pi)^d} \left| \int_{\mathbb{R}^d} |\xi|^{-2s} |\hat{\rho}_s(\xi)|^2 d\xi - \int_{\mathbb{R}^d} |\hat{\rho}(\xi)|^2 d\xi \right| \\ &\leq \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} ||\xi|^{-2s} - 1| |\hat{\rho}_s(\xi)|^2 d\xi + \frac{1}{(2\pi)^d} \left| \int_{\mathbb{R}^d} |\hat{\rho}_s(\xi)|^2 d\xi - \int_{\mathbb{R}^d} |\hat{\rho}(\xi)|^2 d\xi \right|. \end{aligned} \quad (3.9)$$

About the first term in the right hand side, for any $R > 1$ we have, since $|\hat{\rho}_s(\xi)| \leq M$ for any $\xi \in \mathbb{R}^d$,

$$\begin{aligned} \int_{\mathbb{R}} ||\xi|^{-2s} - 1| |\hat{\rho}_s(\xi)|^2 d\xi &= \int_{B_R} ||\xi|^{-2s} - 1| |\hat{\rho}_s(\xi)|^2 d\xi + \int_{B_R^C} ||\xi|^{-2s} - 1| |\hat{\rho}_s(\xi)|^2 d\xi \\ &\leq M^2 \int_{B_R} ||\xi|^{-2s} - 1| d\xi + 2 \int_{B_R^C} |\hat{\rho}_s(\xi)|^2 d\xi \\ &\leq M^2 \int_{B_R} ||\xi|^{-2s} - 1| d\xi + 4 \int_{\mathbb{R}^d} |\hat{\rho}_s(\xi) - \hat{\rho}(\xi)|^2 d\xi \\ &\quad + 4 \int_{B_R^C} |\hat{\rho}(\xi)|^2 d\xi. \end{aligned}$$

As $s \downarrow 0$, the first term in the right hand side goes to zero by dominated convergence (as dominating function we take $|\xi|^{-2s_0} + 1$ for $|\xi| \leq 1$ and 2 for $1 < |\xi| \leq R$). Therefore, by the strong $L^2(\mathbb{R}^d)$ convergence of ρ_s to ρ , from (3.9) we get

$$\limsup_{s \downarrow 0} \left| \int_{\mathbb{R}^d} c_{d,s} |x - y|^{2s-d} \rho_s(x) \rho_s(y) dx dy - \int_{\mathbb{R}^d} \rho^2(x) dx \right| \leq 4 \int_{B_R^C} |\hat{\rho}(\xi)|^2 d\xi.$$

The result follows, since $\rho \in L^2(\mathbb{R}^d)$ and R is arbitrary. \square

Now we are in the position to state the main result about convergence of minimizers ρ_s towards ρ_0 , where ρ_0 is defined in (3.2).

Theorem 3.9. Assume $0 \leq \beta < \chi/2$. For any $s \in (0, 1/2)$, let $\rho_s \in \mathcal{Y}_M$ be the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M . Then, there exists $\rho \in \mathcal{Y}_M$ such that $\rho_s \rightarrow \rho$ strongly in $L^m(\mathbb{R}^d)$ as $s \downarrow 0$. Moreover, ρ is the unique radially decreasing minimizer of the functional (3.1) over \mathcal{Y}_M , given by (3.2).

Proof. Let $(s_n) \subset (0, 1/2)$ be a vanishing sequence. Let $\rho \in \mathcal{Y}_M$ be such that $\rho_{s_n} \rightarrow \rho$ strongly in $L^m(\mathbb{R}^d)$ as $n \rightarrow +\infty$. The existence of such a limit point follows from Lemma 3.7, and the convergence holds in $L^p(\mathbb{R}^d)$ for any $p \in [1, +\infty)$. Given $\tilde{\rho} \in \mathcal{Y}_M$, by the strong $L^m(\mathbb{R}^d)$ convergence and by Lemma 3.8 we have

$$\mathcal{F}_0[\rho] = \lim_{n \rightarrow +\infty} \mathcal{F}_{s_n}[\rho_{s_n}] \leq \lim_{n \rightarrow +\infty} \mathcal{F}_{s_n}[\tilde{\rho}] = \mathcal{F}_0[\tilde{\rho}].$$

By the arbitrariness of $\tilde{\rho}$, we conclude that ρ is a minimizer of \mathcal{F}_0 over \mathcal{Y}_M . Finally, the whole family (ρ_s) converges to ρ in $L^m(\mathbb{R}^d)$ as $s \downarrow 0$. \square

Remark 3.10. For $0 \leq \beta < \chi/2$, and given $M > 0$, we have in fact the Γ -convergence of functionals $\mathcal{F}_s : \mathcal{Y}_M \rightarrow \mathbb{R}$ to functional $\mathcal{F}_0 : \mathcal{Y}_M \rightarrow \mathbb{R}^d$ as $s \rightarrow 0$, with respect to the strong $L^2(\mathbb{R}^d)$ topology. Indeed, if $(\zeta_s)_{s \in (0, 1/2)} \subset \mathcal{Y}_M$, $\zeta \in \mathcal{Y}_M$ and $\zeta_s \rightarrow \zeta$ in $L^2(\mathbb{R}^d)$ as $s \rightarrow 0$, by Fatou's lemma and Lemma 3.8 we get $\mathcal{F}_0[\zeta] \leq \liminf_{s \rightarrow 0} \mathcal{F}_s[\zeta_s]$. On the other hand, still by Lemma 3.8, for every $\zeta \in \mathcal{Y}_M$ we have $\mathcal{F}_s[\zeta] \rightarrow \mathcal{F}_0[\zeta]$ as $s \rightarrow 0$.

Proof of Theorem 1.2. The case $0 \leq \beta < \chi/2$ has been treated in Theorem 3.9. Therefore, in order to conclude we consider the case $\beta \geq \chi/2$.

For every $\rho \in \mathcal{Y}_M$ and every $\lambda > 0$, let $\rho_\lambda(x) = \lambda^d \rho(\lambda x)$, $x \in \mathbb{R}^d$. We have

$$\mathcal{F}_s[\rho_\lambda] \geq \beta \lambda^d \int_{\mathbb{R}^d} \rho^2(x) dx + \lambda^{d-2s} \mathcal{W}_s[\rho]. \quad (3.10)$$

Similarly to the proof of Proposition 3.1, we minimize the right hand side with respect to λ and find a unique optimal value λ_* given by

$$\lambda_* = ((2s - d) \mathcal{W}_s[\rho])^{\frac{1}{2s}} \left(d \beta \int_{\mathbb{R}^d} \rho^2(x) dx \right)^{-\frac{1}{2s}},$$

and by inserting such a value in (3.10) we get

$$\begin{aligned} \mathcal{F}_s[\rho] &\geq -\frac{2s\beta}{d-2s} \left(\frac{\chi(d-2s)}{2d\beta} \right)^{\frac{d}{2s}} \left(\int_{\mathbb{R}^d} \rho^2(x) dx \right)^{1-\frac{d}{2s}} \\ &\quad \times \left(\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho(x) \rho(y) dx dy \right)^{\frac{d}{2s}}. \end{aligned}$$

But the Hardy-Littlewood-Sobolev inequality (2.6) along with interpolation of L^p norms, entails

$$\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho(x) \rho(y) dx dy \leq c_{d,s} H_{d,s} \|\rho\|_{\frac{2d}{d+2s}}^2 \leq c_{d,s} H_{d,s} M^{\frac{4s}{d}} \left(\int_{\mathbb{R}^d} \rho^2(x) dx \right)^{1-\frac{2s}{d}},$$

therefore we get the estimate

$$\mathcal{F}_s[\rho] \geq -\frac{2s\beta}{d-2s} \left(\frac{\chi(d-2s)}{2d\beta} \right)^{\frac{d}{2s}} M^2 (c_{d,s} H_{d,s})^{\frac{d}{2s}}$$

for every $\rho \in \mathcal{Y}_M$. It is not difficult to check that $(c_{d,s} H_{d,s})^{\frac{d}{2s}}$ converges to a finite limit as $s \downarrow 0$, therefore the above right hand side is negative and converges to 0 as $s \downarrow 0$ due to the condition $\beta \geq \chi/2$. If ρ_s denotes the unique minimizer of \mathcal{F}_s over \mathcal{Y}_M , we deduce from (2.14) that $\lim_{s \downarrow 0} \mathcal{F}_s[\rho_s] = 0$. This implies from (2.13) and (2.14) that $\rho_s \rightarrow 0$ in $L^m(\mathbb{R}^d)$ and $\mathcal{C}_s \rightarrow 0$ as $s \downarrow 0$.

Eventually, we prove that $\|\rho_s\|_\infty \rightarrow 0$ as $s \downarrow 0$. Since ρ_s is continuous and radially decreasing we have $\|\rho_s\|_\infty = \rho_s(0)$, and as seen in the proof of Lemma 3.4 we may take advantage of (2.12) and get

$$\frac{m}{m-1} \rho_s(0)^{m-1} + 2\beta \rho_s(0) = \chi c_{d,s} (|\cdot|^{2s-d} * \rho_s)(0) - \mathcal{C}_s \leq \frac{\chi c_{d,s} \sigma_d}{2s} \rho_s(0) + \chi c_{d,s} M. \quad (3.11)$$

Let $\bar{\rho} := \limsup_{s \downarrow 0} \rho_s(0)$. By taking (1.2) and $\sigma_d = 2\pi^{d/2}/\Gamma(d/2)$ into account, we have $\lim_{s \downarrow 0} (2s)^{-1} c_{d,s} \sigma_d = 1$, thus from (3.11) we deduce

$$\frac{m}{m-1} \bar{\rho}^{m-1} + 2\beta \bar{\rho} \leq \chi \bar{\rho}.$$

Since $m > 2$ and $\beta \geq \chi/2$, this forces $\bar{\rho} = 0$ which is the desired result. \square

4. Weak solutions for the aggregation-diffusion problem

The first objective of this section is the proof of the main existence result stated in Theorem (1.3). We mention that an alternative existence proof for $s > 1/2$ and $\beta = 0$ is found in [37].

We fix $M > 0$. For $\rho^0 \in \mathcal{Y}_{M,2}$, and we consider the Cauchy problem (1.1) We shall construct a weak solution to problem (1.1) by an application of the JKO scheme to the functional \mathcal{F}_s defined by (1.3). Therefore, for a discrete time step $\tau > 0$, we consider the minimization problem

$$\min_{\rho \in \mathcal{Y}_{M,2}} \left(\mathcal{F}_s[\rho] + \frac{1}{2\tau} W_2^2(\rho, \rho^0) \right). \quad (4.1)$$

Proposition 4.1 (Existence of discrete minimizers). *Let $\tau > 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. The minimization problem (4.1) admits solutions.*

Proof. It is clear from Lemma 2.5 that the functional to be minimized over $\mathcal{Y}_{M,2}$ is bounded from below. Let $(\rho_n) \subset \mathcal{Y}_{M,2}$ be one of its minimizing sequences. The sequence (ρ_n) has uniformly bounded $L^m(\mathbb{R}^d)$ norm by inequality (2.7). It also has uniformly bounded second moment, thanks to the uniform bound for $W_2(\rho_n, \rho^0)$, which follows from the fact that ρ_n is a minimizing sequence and again from (2.7) and (2.10). Hence, up to subsequences, it converges to ρ_* weakly in $L^p(\mathbb{R}^d)$ for every $p \in [1, m]$ and narrowly by Prokhorov's theorem (see e.g. [1, Theorem 5.1.3]). This implies that ρ_* has mass M and by [1, Lemma 5.17] that ρ_* has 0 center of mass, therefore $\rho_* \in \mathcal{Y}_M$. Furthermore, by the lower semicontinuity of the second moments with respect to the narrow topology we have $\rho_* \in \mathcal{Y}_{M,2}$. We notice that also the sequence of product measures $\rho_n(x) dx \rho_n(y) dy$ is narrowly converging to $\rho_*(x) dx \rho_*(y) dy$, see for instance [2, Theorem 2.8].

Let $\varepsilon > 0$ and let $\eta_\varepsilon : \mathbb{R}^d \rightarrow [0, 1]$ be a smooth cutoff function such that $\eta_\varepsilon(x) = 1$ if $|x| < \varepsilon$ and $\eta_\varepsilon(x) = 0$ if $|x| > 2\varepsilon$. As $K_s(1 - \eta_\varepsilon)$ is a bounded continuous function over \mathbb{R}^d we obtain

$$\begin{aligned} \lim_{n \rightarrow +\infty} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) (1 - \eta_\varepsilon(x-y)) \rho_n(x) \rho_n(y) dx dy \\ = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) (1 - \eta_\varepsilon(x-y)) \rho_*(x) \rho_*(y) dx dy \end{aligned}$$

for every $\varepsilon > 0$. On the other hand, by Cauchy-Schwarz and Young inequality we have

$$\begin{aligned} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \eta_\varepsilon(x-y) \rho_n(x) \rho_n(y) dx dy &\leq \|\rho_n\|_{L^2(\mathbb{R}^d)} \|(\eta_\varepsilon K_s) * \rho_n\|_{L^2(\mathbb{R}^d)} \\ &\leq \|\rho_n\|_{L^2(\mathbb{R}^d)}^2 \|\eta_\varepsilon K_s\|_{L^1(\mathbb{R}^d)} \\ &\leq C \|\eta_\varepsilon K_s\|_{L^1(\mathbb{R}^d)} \end{aligned}$$

for every $n \in \mathbb{N}$ and every $\varepsilon > 0$, where C is a constant that does not depend on n , since (ρ_n) is a bounded sequence in $L^2(\mathbb{R}^d)$. A combination of the two above relations yields

$$\begin{aligned} \limsup_{n \rightarrow +\infty} \left| \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho_n(x) \rho_n(y) dx dy - \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho_*(x) \rho_*(y) dx dy \right| \\ \leq \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \eta_\varepsilon(x-y) \rho_*(x) \rho_*(y) dx dy + C \|\eta_\varepsilon K_s\|_{L^1(\mathbb{R}^d)}, \end{aligned}$$

for every $\varepsilon > 0$. By dominated convergence, the two terms in the right hand side vanish as $\varepsilon \rightarrow 0$, so that the arbitrariness of ε entails

$$\lim_{n \rightarrow +\infty} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho_n(x) \rho_n(y) dx dy = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) \rho_*(x) \rho_*(y) dx dy. \quad (4.2)$$

By the weak lower semicontinuity of the $L^m(\mathbb{R}^d)$ and of the $L^2(\mathbb{R}^d)$ norms, by the narrow lower semicontinuity of $W_2(\cdot, \rho^0)$ (see [1, Proposition 7.1.3]) and thanks to (4.2) we conclude that

$$\mathcal{F}_s[\rho_*] + \frac{1}{2\tau} W_2^2(\rho_*, \rho^0) \leq \liminf_{n \rightarrow +\infty} \left(\mathcal{F}_s[\rho_n] + \frac{1}{2\tau} W_2^2(\rho_n, \rho^0) \right).$$

Since (ρ_n) is a minimizing sequence, we conclude that ρ_* is a solution to problem (4.1). \square

Once existence of a discrete solution is established, we perform a recursive minimization and apply standard arguments from the theory of minimizing movements to obtain convergence of the scheme and existence of a limit curve, as summarized in the next two statements.

Proposition 4.2 (*Basic estimate of minimizing movements*). *Let $\tau > 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. We let $\rho_\tau^0 := \rho^0$ and for every $k \in \mathbb{N}$, we take recursively*

$$\rho_\tau^k \in \operatorname{argmin}_{\mathcal{Y}_{M,2}} \left(\mathcal{F}_s[\cdot] + \frac{1}{2\tau} W_2^2(\cdot, \rho_\tau^{k-1}) \right), \quad (4.3)$$

thus defining a sequence $(\rho_\tau^k)_{k \geq 1}$ of discrete minimizers, whose existence is ensured by Proposition 4.1. For every $k \in \mathbb{N}$ there hold

$$\mathcal{F}_s[\rho_\tau^k] + \frac{1}{2\tau} \sum_{h=1}^k W_2^2(\rho_\tau^h, \rho_\tau^{h-1}) \leq \mathcal{F}_s[\rho^0] \leq \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0(x))^m dx + \beta \int_{\mathbb{R}^d} (\rho^0(x))^2 dx, \quad (4.4)$$

$$\frac{1}{2(m-1)} \int_{\mathbb{R}^d} (\rho_\tau^k(x))^m dx \leq \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0(x))^m dx + \beta \int_{\mathbb{R}^d} (\rho^0(x))^2 dx + \bar{C}, \quad (4.5)$$

and

$$\int_{\mathbb{R}^d} |x|^2 (\rho_\tau^k(x)) dx \leq 4k\tau \mathcal{F}_s[\rho^0] + 2 \int_{\mathbb{R}^d} |x|^2 \rho^0(x) dx, \quad (4.6)$$

where $\bar{C} = \bar{C}(\chi, m, s, d, M)$ is the constant given by (2.11), only depending on χ, m, s, d, M .

Proof. Estimate (4.4) directly follows from the minimality property of ρ_τ^k , which is defined by (4.3). Estimate (4.5) follows from (2.7) and (4.4). Moreover, we have by triangle inequality and by Cauchy Schwarz inequality

$$\begin{aligned} \int_{\mathbb{R}^d} |x|^2 \rho_\tau^k(x) dx &= W_2^2(\rho_\tau^k, M\delta_0) \leq \left(\sum_{h=1}^k W_2(\rho_\tau^h, \rho_\tau^{h-1}) + W_2(\rho^0, M\delta_0) \right)^2 \\ &\leq 2k\tau \sum_{h=1}^k \frac{1}{\tau} W_2^2(\rho_\tau^h, \rho_\tau^{h-1}) + 2W_2(\rho^0, M\delta_0), \end{aligned}$$

which entails (4.6) by means of (4.4). \square

Proposition 4.3 (Convergence of the scheme). *Let $\rho^0 \in \mathcal{Y}_{M,2}$. For every $\tau > 0$, let us consider a sequence $(\rho_\tau^k)_{k \geq 0}$ of discrete minimizers defined by (4.3) and define the piecewise constant interpolation*

$$\rho_\tau(t, \cdot) := \rho_\tau^{\lceil t/\tau \rceil}(\cdot), \quad t \geq 0, \quad (4.7)$$

where $\lceil x \rceil := \min\{h \in \mathbb{N} : h \geq x\}$. Then there exist a vanishing sequence $(\tau_n)_{n \in \mathbb{N}} \subset (0, 1)$ and a limit function $\rho \in L^\infty((0, +\infty); L^m(\mathbb{R}^d))$ such that $[0, +\infty) \ni t \mapsto \rho(t, \cdot) \in \mathcal{Y}_{M,2}$ is narrowly continuous and such that $\rho_{\tau_n}(t, \cdot)$ narrowly converge to $\rho(t, \cdot)$ for every $t \geq 0$. Furthermore, $t \mapsto \rho(t, \cdot)$ is a $AC^2([0, +\infty))$ curve with respect to the Wasserstein distance, i.e., there exists $g \in L^2(0, +\infty)$ such that $W_2(\rho(t_1, \cdot), \rho(t_2, \cdot)) \leq \int_{t_1}^{t_2} g(r) dr$ for every $0 \leq t_1 < t_2 < +\infty$.

Proof. The convergence to a narrowly continuous curve along a vanishing sequence (τ_n) follows from the standard convergence arguments for minimizing movements from [1]. In order to obtain it, we let $g_\tau : [0, +\infty) \rightarrow [0, +\infty)$ be defined as $g_\tau(t) = \tau^{-1} W_2(\rho_\tau(t, \cdot), \rho_\tau(t - \tau, \cdot))$, with the convention $\rho_\tau(r, \cdot) = \rho^0(\cdot)$ if $r < 0$, and we have from (4.7)

$$\frac{1}{2} \int_0^{+\infty} g_\tau^2(t) dt = \frac{1}{2\tau^2} \sum_{h=1}^{+\infty} \int_{(h-1)\tau}^{h\tau} W_2^2(\rho_\tau(t, \cdot), \rho_\tau(t - \tau, \cdot)) dt = \frac{1}{2\tau} \sum_{h=1}^{+\infty} W_2^2(\rho_\tau^h, \rho_\tau^{h-1}).$$

On the other hand, (4.4), (4.5) and (2.10) entail for every integer $k \geq 1$

$$\frac{1}{2\tau} \sum_{h=1}^k W_2^2(\rho_\tau^h, \rho_\tau^{h-1}) \leq \mathcal{F}_s[\rho^0] - \mathcal{F}_s[\rho_\tau^k] \leq 2 \left(\bar{C} + \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0)^m + \beta \int_{\mathbb{R}^d} (\rho^0)^2 \right),$$

where \bar{C} is defined by (2.11) and (2.8). We deduce that

$$\frac{1}{2} \int_0^{+\infty} g_\tau^2(t) dt \leq 2 \left(\bar{C} + \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0)^m + \beta \int_{\mathbb{R}^d} (\rho^0)^2 \right)$$

so that $g_\tau \in L^2(0, +\infty)$ there exists a vanishing sequence $(\tau_n)_{n \in \mathbb{N}} \subset (0, 1)$ such that $g_{\tau_n} \rightharpoonup g$ weakly in $L^2((0, +\infty))$ for some $g \in L^2((0, +\infty))$. For arbitrary $T > 0$, the family of functions $\{\rho_{\tau_n}(t, \cdot) : n \in \mathbb{N}, t \in [0, T]\} \subset \mathcal{Y}_{M,2}$ has uniformly bounded second moments thanks to (4.6), hence it is narrowly relatively compact, and moreover if $0 \leq t_1 < t_2 \leq T$ we may apply the triangle inequality to find the Wasserstein equi-continuity estimate

$$\begin{aligned} \limsup_{n \rightarrow +\infty} W_2(\rho_{\tau_n}(t_2, \cdot), \rho_{\tau_n}(t_1, \cdot)) &\leq \limsup_{n \rightarrow +\infty} \sum_{h=\lceil t_1/\tau_n \rceil + 1}^{\lceil t_2/\tau_n \rceil} W_2(\rho_{\tau_n}^h, \rho_{\tau_n}^{h-1}) \\ &= \limsup_{n \rightarrow +\infty} \int_{(\lceil t_1/\tau_n \rceil + 1)\tau_n}^{\lceil t_2/\tau_n \rceil \tau_n} g_{\tau_n}(t) dt = \int_{t_1}^{t_2} g(t) dt, \end{aligned}$$

so that we can apply the abstract Ascoli-Arzelà theorem from [1, Proposition 3.3.1] and deduce that there exists a narrowly continuous curve $[0, T] \ni t \mapsto \rho(t, \cdot)$ such that, up to extraction of a not relabeled subsequence, $\rho_{\tau_n}(t, \cdot) \rightarrow \rho(t, \cdot)$ narrowly as $n \rightarrow +\infty$ for every $t \in [0, T]$. Eventually, by the above estimate and the narrow lower semicontinuity of the Wasserstein distance (see [1, Proposition 7.1.3]), we deduce $W_2(\rho(t_2, \cdot), \rho(t_1, \cdot)) \leq \int_{t_1}^{t_2} g(t) dt$, which is the desired AC^2 property. By the uniform estimate (4.5) and the narrow lower semicontinuity of the L^m norm (see for instance [30, Proposition 7.7] we also deduce that $\|\rho(t, \cdot)\|_{L^m(\mathbb{R}^d)}^m \leq 2C(m-1)$ for every $t \geq 0$, where C is the right hand side of (4.5). \square

The curve $t \mapsto \rho(t, \cdot)$ that was obtained in Proposition 4.3 will be shown to be a weak solution to (1.1). The first step towards this goal is to obtain a first order optimality condition for discrete minimizers of the JKO scheme.

Proposition 4.4 (Euler-Lagrange equation for discrete minimizers). *Let $\tau > 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. If ρ_* is a solution to problem (4.1), then there holds*

$$\begin{aligned} \frac{1}{\tau} \int_{\mathbb{R}^d} (T_*(x) - x) \cdot \nabla \zeta(x) \rho_*(x) dx &= - \int_{\mathbb{R}^d} \Delta \zeta(x) (\rho_*(x)^m + \beta \rho_*(x)^2) dx \\ &+ \frac{(d-2s)\chi_{c_{d,s}}}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \zeta(x) - \nabla \zeta(y)) \cdot (x-y) |x-y|^{2s-d-2} \rho_*(x) \rho_*(y) dx dy \end{aligned} \quad (4.8)$$

for every $\zeta \in C^\infty(\mathbb{R}^d)$ such that $\nabla \zeta \in W^{1,\infty}(\mathbb{R}^d; \mathbb{R}^d)$. Here, T_* is the unique optimal transport map (for the quadratic cost) from $\rho_*(x) dx$ to $\rho^0(x) dx$.

Proof. Let $\zeta \in C^\infty(\mathbb{R}^d)$ be such that $\nabla\zeta \in W^{1,\infty}(\mathbb{R}^d; \mathbb{R}^d)$ and such that

$$\int_{\mathbb{R}^d} \rho_* \nabla \zeta \, dx = 0. \quad (4.9)$$

Let $\rho_\varepsilon(x) \, dx$ be the push-forward measure of $\rho_*(x) \, dx$ through the map $x \mapsto x + \varepsilon \nabla \zeta$, defined for any $\varepsilon \in \mathbb{R}$ such that $|\varepsilon| < 1/L_\zeta$, where L_ζ the Lipschitz constant of $\nabla \zeta$. It is clear that $\rho_\varepsilon \in \mathcal{Y}_M$. By the change of variables formula and by the Taylor expansion

$$\det(I + \varepsilon \nabla^2 \zeta) = 1 + \varepsilon \Delta \zeta + o(\varepsilon)$$

we get

$$\begin{aligned} \int_{\mathbb{R}^d} \rho_\varepsilon(x)^m \, dx &= \int_{\mathbb{R}^d} \frac{\rho_*(x)^m}{(\det(I + \varepsilon \nabla^2 \zeta(x)))^{m-1}} \, dx \\ &= \int_{\mathbb{R}^d} \rho_*(x)^m \, dx - (m-1)\varepsilon \int_{\mathbb{R}^d} \Delta \zeta(x) \rho_*(x)^m \, dx + o(\varepsilon) \end{aligned}$$

where I is the identity matrix and ∇^2 is the Hessian operator. Therefore we have

$$\frac{d}{d\varepsilon} \left(\frac{1}{m-1} \int_{\mathbb{R}^d} \rho_\varepsilon(x)^m \, dx \right) \Big|_{\varepsilon=0} = - \int_{\mathbb{R}^d} \Delta \zeta(x) \rho_*(x)^m \, dx,$$

which is of course still true if m is replaced by 2. On the other hand, the definition of push-forward entails

$$\begin{aligned} &\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho_\varepsilon(x) \rho_\varepsilon(y) \, dx \, dy \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y + \varepsilon(\nabla \zeta(x) - \nabla \zeta(y))|^{2s-d} \rho_*(x) \rho_*(y) \, dx \, dy \end{aligned}$$

so that

$$\begin{aligned} &\frac{1}{\varepsilon} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho_\varepsilon(x) \rho_\varepsilon(y) \, dx \, dy - \frac{1}{\varepsilon} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x-y|^{2s-d} \rho_*(x) \rho_*(y) \, dx \, dy \\ &= \frac{1}{\varepsilon} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (|x-y + \varepsilon(\nabla \zeta(x) - \nabla \zeta(y))|^{2s-d} - |x-y|^{2s-d}) \rho_*(x) \rho_*(y) \, dx \, dy. \end{aligned}$$

It is clear that

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (|x - y + \varepsilon(\nabla\zeta(x) - \nabla\zeta(y))|^{2s-d} - |x - y|^{2s-d}) \\ = (2s-d)|x - y|^{2s-d-2}(x - y) \cdot (\nabla\zeta(x) - \nabla\zeta(y)) \end{aligned}$$

for every $x, y \in \mathbb{R}^d$, $x \neq y$. On the other hand, since $|\varepsilon| < 1/L_\zeta$ we have $|\varepsilon||\nabla\zeta(x) - \nabla\zeta(y)| < |x - y|$ and then we can obtain the estimate

$$|x - y + \varepsilon(\nabla\zeta(x) - \nabla\zeta(y))|^{2s-d} - |x - y|^{2s-d} \leq ((1 - |\varepsilon|L_\zeta)^{2s-d} - 1) |x - y|^{2s-d}$$

for every $x, y \in \mathbb{R}^d$, $x \neq y$. Thus, Bernoulli inequality entails, for every $x, y \in \mathbb{R}^d$, $x \neq y$ and every $\varepsilon \in \mathbb{R}$ such that $2|\varepsilon| < 1/L_\zeta$,

$$|x - y + \varepsilon(\nabla\zeta(x) - \nabla\zeta(y))|^{2s-d} - |x - y|^{2s-d} < |\varepsilon| d L_\zeta 2^{d-2s} |x - y|^{2s-d},$$

therefore by dominated convergence we get

$$\begin{aligned} \frac{d}{d\varepsilon} \left(\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x - y|^{2s-d} \rho_\varepsilon(x) \rho_\varepsilon(y) dx dy \right) \Big|_{\varepsilon=0} \\ = (2s-d) \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |x - y|^{2s-d-2} (x - y) \cdot (\nabla\zeta(x) - \nabla\zeta(y)) \rho_*(x) \rho_*(y) dx dy. \end{aligned}$$

For the derivative of the Wasserstein distance, by a standard result (see [36, Theorem 8.13]) we get

$$\frac{d}{d\varepsilon} \left(\frac{1}{2\tau} W_2^2(\rho_\varepsilon, \rho^0) \right) \Big|_{\varepsilon=0} = \frac{1}{\tau} \int_{\mathbb{R}^d} (x - T_*(x)) \cdot \nabla\zeta(x) \rho_*(x) dx$$

Since ρ_* is a minimizer, the derivative with respect to ε of $\mathcal{F}_s[\rho_\varepsilon] + \frac{1}{2\tau} W_2^2(\rho_\varepsilon, \rho^0)$ needs to vanish at $\varepsilon = 0$. We obtain the result.

If we wish to remove the compatibility condition (4.9), we just replace ζ with

$$\tilde{\zeta}(x) = \zeta(x) - \frac{1}{M} \left(\int_{\mathbb{R}^d} \rho_* \nabla\zeta dx \right) \cdot x,$$

in order to have

$$\nabla\tilde{\zeta}(x) = \nabla\zeta(x) - \frac{1}{M} \int_{\mathbb{R}^d} \rho_* \nabla\zeta dx,$$

hence $\tilde{\zeta}$ satisfies (4.9). Inserting $\tilde{\zeta}$ in (4.8) and taking into account that

$$\int_{\mathbb{R}^d} T_*(x) \cdot \left(\int_{\mathbb{R}^d} \rho_* \nabla \zeta \, dx \right) \rho_*(x) dx = \left(\int_{\mathbb{R}^d} \rho_* \nabla \zeta \, dx \right) \cdot \int_{\mathbb{R}^d} x \rho_0(x) dx = 0,$$

we have that (4.8) holds for any test function $\zeta \in C^\infty(\mathbb{R}^d)$ such that $\nabla \zeta \in W^{1,\infty}(\mathbb{R}^d; \mathbb{R}^d)$. \square

The next result is based on a different perturbation of ρ_* , which gets perturbed along the solution of the heat equation originating from it. For nonnegative $L^1(\mathbb{R}^d)$ functions u with finite second moment on \mathbb{R}^d we introduce the entropy functional

$$\mathcal{G}[u] := \int_{\mathbb{R}^d} u(x) \log u(x) \, dx,$$

which is a displacement convex functional in the sense of McCann [24]. We recall that the solution u of the heat equation $\partial_t u = \Delta u$ with initial datum $\rho^0 \in \mathcal{Y}_{M,2}$ is the Wasserstein gradient flow of \mathcal{G} and it satisfies the evolution variational inequalities

$$\frac{1}{2} W_2^2(u(t, \cdot), w) - \frac{1}{2} W_2^2(\rho^0, w) \leq t (\mathcal{G}(w) - \mathcal{G}(u(t, \cdot)))$$

for every $w \in \mathcal{Y}_{M,2}$ and every $t > 0$,

see [1, Chapter 11]. This allows to take advantage of the flow interchange lemma introduced in [23], as we do in the next proof.

Proposition 4.5 (improved regularity of discrete minimizers). *Let $\tau > 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. Suppose that $\beta > 0$. If ρ_* is a solution to problem (4.1), then $\rho_*^{m/2} \in H^1(\mathbb{R}^d)$ and*

$$\frac{4}{m} \int_{\mathbb{R}^d} |\nabla \rho_*^{m/2}|^2 \, dx \leq \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|\rho_*\|_{L^2(\mathbb{R}^d)}^2 + \frac{\mathcal{G}[\rho^0] - \mathcal{G}[\rho_*]}{\tau}.$$

Proof. Let us introduce the Cauchy problem

$$\begin{cases} \partial_t u = \Delta u \\ u(0) = \rho_*. \end{cases} \quad (4.10)$$

The unique solution to the heat equation with initial datum in $\rho_* \in \mathcal{Y}_{M,2}$ is given by $u(t, \cdot) = \Gamma_t * \rho_*$, where $\Gamma_t(x) := (4\pi t)^{-d/2} \exp\{-|x|^2/(4t)\}$ is the Gaussian kernel. $u(t, \cdot)$ is smooth, positive for every $t > 0$, and moreover a direct computation by means of integration by parts shows that for any $t > 0$ there holds

$$\frac{1}{m-1} \frac{d}{dt} \int_{\mathbb{R}^d} u(t, x)^m \, dx = -\frac{4}{m} \int_{\mathbb{R}^d} \left| \nabla u(t, x)^{\frac{m}{2}} \right|^2 \, dx \quad (4.11)$$

and in particular

$$\beta \frac{d}{dt} \int_{\mathbb{R}^d} u(t, x)^2 dx = -2\beta \int_{\mathbb{R}^d} |\nabla u(t, x)|^2 dx. \quad (4.12)$$

Similarly, thanks to [22, Lemma 4.5] we have for any $t > 0$

$$\frac{\chi}{2} \frac{d}{dt} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) u(t, x) u(t, y) dx dy = -\chi \|u(t, \cdot)\|_{\dot{H}^{1-s}(\mathbb{R}^d)}. \quad (4.13)$$

Here, $\dot{H}^r(\mathbb{R}^d)$ denotes the homogeneous Sobolev space of order $r \in \mathbb{R}$, i.e., the completion of $C_c^\infty(\mathbb{R}^d)$ with respect to the norm $\|w\|_{\dot{H}^r(\mathbb{R}^d)}^2 := (2\pi)^{-d} \int_{\mathbb{R}^d} |\xi|^{2r} |\hat{w}(\xi)|^2 d\xi$. By the interpolation inequality $\|w\|_{\dot{H}^{1-s}(\mathbb{R}^d)} \leq \|w\|_{L^2(\mathbb{R}^d)}^{1-s} \|w\|_{\dot{H}^1(\mathbb{R}^d)}^s$ and by Young inequality we deduce, for given $\alpha > 0$,

$$\begin{aligned} \chi \|u(t, \cdot)\|_{\dot{H}^{1-s}(\mathbb{R}^d)}^2 &\leq \chi \|u(t, \cdot)\|_{L^2(\mathbb{R}^d)}^{2s} \|u(t, \cdot)\|_{\dot{H}^1(\mathbb{R}^d)}^{2(1-s)} \\ &\leq \chi s \alpha^{-1/s} \|u(t, \cdot)\|_{L^2(\mathbb{R}^d)}^2 + \chi(1-s) \alpha^{1/(1-s)} \|u(t, \cdot)\|_{\dot{H}^1(\mathbb{R}^d)}^2. \end{aligned} \quad (4.14)$$

The choice $\alpha = (\chi(1-s)/(2\beta))^{s-1}$ in (4.14) entails, together with (4.11) and (4.12),

$$\frac{d}{dt} \mathcal{F}_s(u(t, \cdot)) \leq -\frac{4}{m} \int_{\mathbb{R}^d} \left| \nabla u(t, x)^{\frac{m}{2}} \right|^2 dx + \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|u(t, \cdot)\|_{L^2(\mathbb{R}^d)}^2 \quad (4.15)$$

for every $t > 0$. Moreover, the maps

$$\begin{aligned} [0, +\infty) \ni t &\mapsto \frac{1}{m-1} \int_{\mathbb{R}^d} u(t, x)^m dx, & [0, +\infty) \ni t &\mapsto \int_{\mathbb{R}^d} u(t, x)^2 dx, \\ [0, +\infty) \ni t &\mapsto \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) u(t, x) u(t, y) dx dy \end{aligned}$$

are continuous up to $t = 0$. They are also differentiable at any $t > 0$ with derivatives given by (4.11), (4.12) and (4.13), therefore by Lagrange mean value theorem, for every $t > 0$ there exists $\theta(t) \in (0, t)$ such that

$$\frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} = -\frac{d}{dt} \mathcal{F}_s[u(t, \cdot)] \Big|_{t=\theta(t)},$$

so that by applying (4.15) we obtain

$$\frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} \geq \frac{4}{m} \int_{\mathbb{R}^d} \left| \nabla u(\theta(t), x)^{\frac{m}{2}} \right|^2 dx - \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|u(\theta(t), \cdot)\|_{L^2(\mathbb{R}^d)}^2$$

Since the $L^2(\mathbb{R}^d)$ norm decreases along the solution to the heat equation (4.10), we deduce

$$\begin{aligned} \limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} &\geq \limsup_{t \downarrow 0} \frac{4}{m} \int_{\mathbb{R}^d} |\nabla u(\theta(t), x)|^{\frac{m}{2}} dx \\ &\quad - \chi^s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|\rho_*\|_{L^2(\mathbb{R}^d)}^2, \end{aligned}$$

showing that

$$\limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} > -\infty.$$

Therefore, we can apply the flow interchange lemma from [23], in its version from [22, Proposition 4.3] and deduce

$$\begin{aligned} \frac{\mathcal{G}[\rho^0] - \mathcal{G}[\rho_*]}{\tau} &\geq \limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} \\ &\geq \limsup_{t \downarrow 0} \frac{4}{m} \int_{\mathbb{R}^d} |\nabla u(\theta(t), x)|^{\frac{m}{2}} dx - \chi^s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|\rho_*\|_{L^2(\mathbb{R}^d)}^2, \end{aligned}$$

thus showing that the spatial gradient of $u(\theta(t), \cdot)^{m/2}$ stays bounded in $L^2(\mathbb{R}^d)$ as $t \downarrow 0$. But $u(\theta(t), \cdot)^{m/2}$ is also bounded in $L^1(\mathbb{R}^d)$ as $t \downarrow 0$, since $\|u(\theta(t), \cdot)\|_{L^{m/2}(\mathbb{R}^d)} \leq \|\rho_*\|_{L^{m/2}(\mathbb{R}^d)}$. Thus Sobolev embedding shows that $u(\theta(t), \cdot)^{m/2}$ is in fact bounded in $H^1(\mathbb{R}^d)$ as $t \downarrow 0$. Since $\theta(t) \rightarrow 0$ as $t \downarrow 0$, and since $u(\theta(t), \cdot) \rightarrow \rho_*$ pointwise a.e. as $t \downarrow 0$, by the weak lower semicontinuity of the $H^1(\mathbb{R}^d)$ norm we finally deduce

$$\chi^s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \|\rho_*\|_{L^2(\mathbb{R}^d)}^2 + \frac{\mathcal{G}[\rho^0] - \mathcal{G}[\rho_*]}{\tau} \geq \frac{4}{m} \int_{\mathbb{R}^d} |\nabla \rho_*^{\frac{m}{2}}|^2 dx,$$

which is the desired result. \square

Remark 4.6. The constant $\chi^s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}}$ appearing in the above result is bounded as $s \rightarrow 0$ if and only if $\beta \geq \chi/2$.

Corollary 4.7. Let $\rho^0 \in \mathcal{Y}_{M,2}$ and $\tau > 0$. Let $\beta > 0$. Let us consider the sequence $(\rho_\tau^k)_{k \geq 0}$ of discrete minimizers defined by (4.3) and the piecewise constant interpolation ρ_τ defined by (4.7). There holds for any $k \in \mathbb{N}$

$$\frac{4}{m} \int_{\mathbb{R}^d} |\nabla (\rho_\tau^k(x)^{m/2})|^2 dx \leq \frac{\mathcal{G}[\rho_\tau^{k-1}] - \mathcal{G}[\rho_\tau^k]}{\tau} + \chi^s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \int_{\mathbb{R}^d} (\rho_\tau^k(x))^2 dx.$$

Moreover, for every $T > 0$ there holds the time integrated estimate

$$\frac{4}{m} \int_0^T \int_{\mathbb{R}^d} |\nabla(\rho_\tau(t, x))^{m/2}|^2 dx dt \leq C_1^* + C_2^*(T + \tau) + C_3^*(T + \tau) \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}}, \quad (4.16)$$

where C_i^* , $i = 1, 2, 3$, are suitable explicit constants, only depending on χ, M, m, s, d, β , and on ρ^0 .

Proof. The first estimate in the statement is a direct consequence of Proposition 4.5, and it implies that for every $T > 0$ we have

$$\begin{aligned} \frac{4}{m} \int_0^T \int_{\mathbb{R}^d} |\nabla(\rho_\tau(t, x)^{m/2})|^2 dx dt &\leq \frac{4}{m} \int_0^{\lceil T/\tau \rceil \tau} \int_{\mathbb{R}^d} |\nabla(\rho_\tau(t, x)^{m/2})|^2 dx dt \\ &= \frac{4\tau}{m} \sum_{k=1}^{\lceil T/\tau \rceil} \int_{\mathbb{R}^d} |\nabla(\rho_\tau^k(x)^{m/2})|^2 dx \\ &\leq \mathcal{G}[\rho^0] - \mathcal{G}[\rho_\tau^{\lceil T/\tau \rceil}] + \tau \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}} \sum_{k=1}^{\lceil T/\tau \rceil} \int_{\mathbb{R}^d} (\rho_\tau^k(x))^2 dx. \end{aligned} \quad (4.17)$$

By (4.6) and by the standard estimate from [4, Lemma 2.2], and since $\mathcal{G}[\rho_\tau^k] \leq \int_{\mathbb{R}^d} (\rho_\tau^k)^m dx$ (recall that $m > 2$), we have for every $k \in \mathbb{N}$

$$\begin{aligned} |\mathcal{G}[\rho_\tau^k]| &\leq \mathcal{G}[\rho_\tau^k] + \frac{1}{2} \int_{\mathbb{R}^d} |x|^2 \rho_\tau^k(x) dx + d \log(4\pi) + 2e^{-1} \\ &\leq \int_{\mathbb{R}^d} \rho_\tau^k(x)^m dx + 2\tau k \mathcal{F}_s[\rho^0] + \int_{\mathbb{R}^d} |x|^2 \rho^0(x) dx + d \log(4\pi) + 2e^{-1}. \end{aligned}$$

From (4.5) we also have for every $k \in \mathbb{N}$

$$\begin{aligned} \int_{\mathbb{R}^d} (\rho_\tau^k(x))^2 dx &\leq M + \int_{\mathbb{R}^d} (\rho_\tau^k(x))^m dx \\ &\leq M + 2 \int_{\mathbb{R}^d} (\rho^0)^m dx + 2(m-1) \left(\bar{C} + \beta \int_{\mathbb{R}^d} (\rho^0)^2 dx \right) =: C_3^*, \end{aligned} \quad (4.18)$$

We insert the latter two estimates, combined with (4.4)-(4.5), into (4.17): since $\lceil T/\tau \rceil \tau \leq T + \tau$ we deduce

$$\frac{4}{m} \int_0^T \int_{\mathbb{R}^d} |\nabla(\rho_\tau(t, x)^{m/2})|^2 dx dt \leq C_1^* + C_2^*(T + \tau) + C_3^*(T + \tau) \chi s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}},$$

where

$$C_1^* := 3 \int_{\mathbb{R}^d} (\rho^0)^m dx + 2(m-1) \left(\bar{C} + \beta \int_{\mathbb{R}^d} (\rho^0)^2 dx \right) + d \log(4\pi) + 2e^{-1} + \int_{\mathbb{R}^d} |x|^2 \rho^0 dx,$$

where C_2^* is twice the right hand side of (4.4) and where C_3^* is defined in (4.18). The proof is concluded. \square

In the case $\beta = 0$ we provide an alternative estimate under the restriction $s \in [1/2, 1)$.

Proposition 4.8. *Let $d \geq 2$, $s \in [1/2, 1)$ and $\beta = 0$. Let $\tau > 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. If ρ_* is a solution to problem (4.1), then $\rho_*^{m-1} \in H^1(\mathbb{R}^d)$ and*

$$\begin{aligned} \frac{m}{2(m-1)} \int_{\mathbb{R}^d} |\nabla \rho_*^{m-1}|^2 dx &\leq \frac{\chi^2 S_{d,2s-1}^2 (m-1)}{2m} \|\rho_*\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2 \\ &\quad + \frac{1}{\tau(m-2)} \left(\int_{\mathbb{R}^d} (\rho^0(x))^{m-1} dx - \int_{\mathbb{R}^d} (\rho_*(x))^{m-1} dx \right), \end{aligned}$$

where $S_{d,2s-1}$ is defined by (2.8) if $1/2 < s < 1$ and $S_{d,0} := 1$.

Proof. Let us introduce the auxiliary functional

$$\mathcal{G}_m[u] := \frac{1}{m-2} \int_{\mathbb{R}^d} u(x)^{m-1} dx$$

and the Cauchy problem

$$\begin{cases} \partial_t u = \Delta u^{m-1} \\ u(0) = \rho_*. \end{cases} \quad (4.19)$$

This is a standard porous media equation with initial datum in $\mathcal{Y}_{M,2}$ and it enjoys the following properties, for which we refer to [35, Theorem 9.12, Proposition 9.13]: there exists a unique strong solution u (meaning that the equation $\partial_t u = \Delta u^{m-1}$ is satisfied pointwise a.e. in space-time) such that $u \in C^0([0, +\infty); L^1(\mathbb{R}^d))$ and $\nabla u^{m-1} \in L^2((0, T) \times (\mathbb{R}^d))$ for every $T > 0$. Moreover, the map $t \mapsto \int_{\mathbb{R}^d} u(t, x)^m dx$ is a nonincreasing absolutely continuous map on $[0, T]$ and

$$\frac{1}{m-1} \frac{d}{dt} \int_{\mathbb{R}^d} u(t, x)^m dx = -\frac{m}{m-1} \int_{\mathbb{R}^d} |\nabla u(t, x)^{m-1}|^2 dx \quad \text{for a.e. } t > 0. \quad (4.20)$$

Moreover, the solution is the Wasserstein gradient flow of the displacement convex functional \mathcal{G}_m , see [1, Theorem 11.2.5].

If we multiply (4.19) by $K_s * u$, an integration by parts argument shows that for every $\eta \in C^\infty((0, +\infty))$

$$\begin{aligned} \frac{1}{2} \int_0^{+\infty} \eta'(t) \int_{\mathbb{R}^d} (K_s * u)(t, x) u(t, x) dx dt &= - \int_0^{+\infty} \int_{\mathbb{R}^d} \eta(t) (K_s * u)(t, x) \partial_t u(t, x) dx dt \\ &= - \int_0^{+\infty} \int_{\mathbb{R}^d} \eta(t) \Delta u(t, x)^{m-1} (K_s * u)(t, x) dx dt \\ &= \int_0^{+\infty} \eta'(t) \int_{\mathbb{R}^d} \nabla u(t, x)^{m-1} \cdot \nabla (K_s * u)(t, x) dx dt \end{aligned} \quad (4.21)$$

We notice that for any $s \in (1/2, 1)$ and any $v \in L_+^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)$, by Plancherel theorem and the Hardy-Littlewood-Sobolev inequality (2.6), with the notation (2.8), there holds

$$\int_{\mathbb{R}^d} |\nabla K_s * v|^2 = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} |\xi|^{1-2s} \hat{v}(\xi)^2 d\xi = \int_{\mathbb{R}^d} v K_{2s-1} * v \leq S_{d,2s-1} \|v\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2,$$

thus $\nabla K_s * v \in (L^2(\mathbb{R}^d))^d$. If $s = 1/2$, the above formula directly shows that $\|K_s * v\|_{L^2(\mathbb{R}^d)} = \|v\|_{L^2(\mathbb{R}^d)}$. Since $1 < \frac{2d}{d+4s-2} \leq 2 < m$ and $u(t, \cdot) \in L^1 \cap L^m(\mathbb{R}^d)$ for every $t \geq 0$ with $\|u(t, \cdot)\|_{L^m(\mathbb{R}^d)} \leq \|\rho_*\|_{L^m(\mathbb{R}^d)}$, we deduce that $\nabla K_s * u \in (L^2((0, T) \times \mathbb{R}^d))^d$ and then we get $\nabla u^{m-1} \cdot \nabla K_s * u \in L^1((0, T) \times \mathbb{R}^d)$. Therefore from (4.21) we see that the map

$$t \mapsto \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) u(t, x) u(t, y) dx dy$$

is in $AC([0, T])$ with a.e. derivative given by

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(x-y) u(t, x) u(t, y) dx dy &= - \int_{\mathbb{R}^d} \nabla u(t, x)^{m-1} \cdot \nabla (K_s * u)(t, x) dx \\ &= - \langle u(t, \cdot), u(t, \cdot)^{m-1} \rangle_{\dot{H}^{1-s}(\mathbb{R}^d)}, \end{aligned} \quad (4.22)$$

where the last equality is due to Plancherel theorem, having introduced the following scalar product on $\dot{H}^r(\mathbb{R}^d)$, $r \in (0, 1)$,

$$\langle v, w \rangle_{\dot{H}^r} := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} |\xi|^{2r} \hat{v}(\xi) \overline{\hat{w}(\xi)} d\xi,$$

so that by Cauchy-Schwarz inequality there holds

$$\langle v, w \rangle_{\dot{H}^r} \leq \|v\|_{\dot{H}^{2r-1}(\mathbb{R}^d)} \|w\|_{\dot{H}^1(\mathbb{R}^d)}.$$

By taking advantage of the latter inequality, we deduce the following estimate for the scalar product $\langle u(t, \cdot), u(t, \cdot)^{m-1} \rangle_{\dot{H}^{1-s}(\mathbb{R}^d)}$. Indeed, if $1/2 < s < 1$, by (2.6), by (2.8) and by Young inequality we have

$$\begin{aligned} \chi \langle u(t, \cdot), u(t, \cdot)^{m-1} \rangle_{\dot{H}^{1-s}(\mathbb{R}^d)} &\leq \chi \|u(t, \cdot)\|_{\dot{H}^{1-2s}(\mathbb{R}^d)} \|u(t, \cdot)^{m-1}\|_{\dot{H}^1(\mathbb{R}^d)} \\ &\leq \chi S_{d,2s-1} \|u(t, \cdot)\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)} \|u(t, \cdot)^{m-1}\|_{\dot{H}^1(\mathbb{R}^d)} \\ &\leq Q_{\chi,m,s,d} \|u(t, \cdot)\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2 + \frac{m}{2(m-1)} \int_{\mathbb{R}^d} |\nabla u(t, \cdot)^{m-1}|^2 dx, \end{aligned} \quad (4.23)$$

where $Q_{\chi,m,s,d} := \chi^2 S_{d,2s-1}^2 (m-1)/(2m)$, which is readily seen to hold also for $s = 1/2$, with the convention $S_{d,0} := 1$.

We have shown the absolute continuity of the map $t \mapsto \mathcal{F}_s[u(t, \cdot)]$, which together with (4.20) and (4.22) entails for every $t > 0$

$$\begin{aligned} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} &= -\frac{1}{t} \int_0^t \mathcal{F}_s[u(r, \cdot)] dr \\ &= \frac{m}{m-1} \frac{1}{t} \int_0^t \int_{\mathbb{R}^d} |\nabla u(r, x)^{m-1}|^2 dx dr - \frac{\chi}{t} \int_0^t \langle u(r, \cdot), u(r, \cdot)^{m-1} \rangle_{\dot{H}^{1-s}(\mathbb{R}^d)} dr. \end{aligned}$$

By applying (4.23), since the $L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)$ norm decreases along the solution to the porous media equation (4.19), we deduce

$$\begin{aligned} \limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} \\ \geq \limsup_{t \rightarrow 0} \frac{m}{2(m-1)} \frac{1}{t} \int_0^t \int_{\mathbb{R}^d} |\nabla u(r, x)^{m-1}|^2 dx dr - Q_{\chi,m,s,d} \|\rho_*\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2, \end{aligned}$$

showing that

$$\limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} > -\infty.$$

Therefore, we can apply the flow interchange lemma, in its version from [22, Proposition 4.3] and deduce

$$\begin{aligned} \frac{\mathcal{G}_m[\rho^0] - \mathcal{G}_m[\rho_*]}{\tau} &\geq \limsup_{t \rightarrow 0} \frac{\mathcal{F}_s[\rho_*] - \mathcal{F}_s[u(t, \cdot)]}{t} \\ &\geq \limsup_{t \rightarrow 0} \frac{m}{2(m-1)} \frac{1}{t} \int_0^t \int_{\mathbb{R}^d} |\nabla u(r, x)^{m-1}|^2 dx dr \\ &\quad - Q_{\chi, m, s, d} \|\rho_*\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2. \end{aligned}$$

By the absolute continuity of the map $t \mapsto \int_0^t \int_{\mathbb{R}^d} |\nabla u(r, x)^{m-1}|^2 dx dr$, we may apply l'Hospital rule and get

$$Q_{\chi, m, s, d} \|\rho_*\|_{L^{\frac{2d}{d+4s-2}}(\mathbb{R}^d)}^2 + \frac{\mathcal{G}_m[\rho^0] - \mathcal{G}_m[\rho_*]}{\tau} \geq \liminf_{t \rightarrow 0} \frac{m}{2(m-1)} \int_{\mathbb{R}^d} |\nabla u(t, x)^{m-1}|^2 dx.$$

By taking a suitable vanishing sequence $(t_n)_{n \in \mathbb{N}}$ of positive numbers, the above bound shows that $\nabla u(t_n, \cdot)^{m-1} \rightarrow u_*$ weakly in $L^2(\mathbb{R}^d)$ as $n \rightarrow +\infty$. But $u(t_n, \cdot)$ strongly converge to ρ_* as $n \rightarrow +\infty$, hence up to subsequences we also have that $u(t_n, \cdot)^{m-1} \rightarrow \rho_*^{m-1}$ pointwise a.e. and weakly in $L^{\frac{m}{m-1}}(\mathbb{R}^d)$ (since $u(t_n, \cdot)$ is bounded in $L^m(\mathbb{R}^d)$). This allows to conclude that $u_* = \nabla \rho_*^{m-1}$, and the weak lower semicontinuity of the $L^2(\mathbb{R}^d)$ norm yields the desired estimate. Since $\nabla \rho_*^{m-1} \in L^2(\mathbb{R}^d)$ and since $\rho_*^{m-1} \in L^{\frac{m}{m-1}}(\mathbb{R}^d)$, $m > 2$, by the Gagliardo-Nirenberg inequality we have $\rho_*^{m-1} \in L^2(\mathbb{R}^d)$ with the inequality

$$\|\rho_*^{m-1}\|_{L^2(\mathbb{R}^d)} \leq C \|\nabla \rho_*^{m-1}\|_{L^2(\mathbb{R}^d)}^\alpha \|\rho_*^{m-1}\|_{L^{\frac{m}{m-1}}(\mathbb{R}^d)}^{1-\alpha},$$

where $\alpha = \frac{d(m-2)}{2m+d(m-2)} \in (0, 1)$. \square

By arguing as done for proving Corollary (4.7), from (4.8) we immediately deduce the following

Corollary 4.9. *Let $d \geq 2$, $s \in [1/2, 1)$ and $\beta = 0$. Let $T > 0$. Let $\rho^0 \in \mathcal{Y}_{M,2}$ and $\tau > 0$. Let us consider the sequence $(\rho_\tau^k)_{k \geq 0}$ of discrete minimizers defined by (4.3) and the piecewise constant interpolation ρ_τ defined by (4.7). Then*

$$\int_0^T \int_{\mathbb{R}^d} |\nabla(\rho_\tau(t, x))^{m-1}|^2 dx dt \leq C_1^{**} + (T + \tau) C_2^{**} \quad (4.24)$$

where C_1^{**} , C_2^{**} are a suitable explicit constants, only depending on χ, M, m, s, d and the initial datum ρ^0 .

Proposition 4.10 (Improved space-time compactness). *Let $\beta \geq 0$ and $\rho^0 \in \mathcal{Y}_{M,2}$. If $\beta = 0$, assume in addition that $d \geq 2$ and $1/2 \leq s < 1$. Let us consider a sequence $(\rho_\tau^k)_{k \geq 0}$ of discrete minimizers defined by (4.3), the piecewise constant interpolation ρ_τ defined by (4.7), and let ρ be a limit function obtained from Proposition 4.3 along a vanishing sequence $(\tau_n)_{n \in \mathbb{N}} \subset (0, 1)$. Then $\rho_{\tau_n} \rightarrow \rho$ strongly in $L^2((0, T) \times \mathbb{R}^d)$ for any $T > 0$.*

Proof. Let T_τ^{k+1} be the unique optimal transport map (for the quadratic cost) from $\rho_\tau^{k+1}(x) dx$ to $\rho_\tau^k(x) dx$. Let $\varphi \in W^{2,\infty}(\mathbb{R}^d)$ be smooth, let $h > 0$, and notice that

$$\begin{aligned} & \left| \int_{\mathbb{R}^d} \varphi(x) (\rho_\tau(t+h, x) - \rho_\tau(t, x)) dx \right| = \left| \sum_{k=\lceil t/\tau \rceil}^{\lceil (t+h)/\tau \rceil - 1} \int_{\mathbb{R}^d} \varphi(x) (\rho_\tau^{k+1}(x) - \rho_\tau^k(x)) dx \right| \\ &= \left| \sum_{k=\lceil t/\tau \rceil}^{\lceil (t+h)/\tau \rceil - 1} \int_{\mathbb{R}^d} (\varphi(\mathsf{T}_\tau^{k+1}(x)) - \varphi(x)) \rho_\tau^{k+1}(x) dx \right| \\ &\leq \sum_{k=\lceil t/\tau \rceil}^{\lceil (t+h)/\tau \rceil - 1} \left| \int_{\mathbb{R}^d} \nabla \varphi(x) \cdot (\mathsf{T}_\tau^{k+1}(x) - x) \rho_\tau^{k+1}(x) dx \right| \\ &+ \left| \sum_{k=\lceil t/\tau \rceil}^{\lceil (t+h)/\tau \rceil - 1} \frac{1}{2} \int_0^1 (1-\xi)^2 d\xi \int_{\mathbb{R}^d} (\mathsf{T}_\tau^{k+1}(x) - x) \nabla^2 \varphi(\vartheta_\xi) \cdot (\mathsf{T}_\tau^{k+1}(x) - x) \rho_\tau^{k+1}(x) dx \right| \end{aligned} \quad (4.25)$$

where we used the Taylor expansion formula

$$\begin{aligned} \varphi(\mathsf{T}_\tau^{k+1}(x)) &= \varphi(x) + \nabla \varphi(x) \cdot (\mathsf{T}_\tau^{k+1}(x) - x) \\ &+ \frac{1}{2} \int_0^1 [(\mathsf{T}_\tau^{k+1}(x) - x) \nabla^2 \varphi(\vartheta_\xi) \cdot (\mathsf{T}_\tau^{k+1}(x) - x)] (1-\xi)^2 d\xi, \end{aligned}$$

where $\vartheta_\xi = \xi \mathsf{T}_\tau^{k+1}(x) + (1-\xi)x$. Let us separately treat the two terms in the right hand side. For the first, by (4.8) we have

$$\begin{aligned}
& \left| \int_{\mathbb{R}^d} \nabla \varphi(x) \cdot (\mathbf{T}_\tau^{k+1}(x) - x) \rho_\tau^{k+1}(x) dx \right| \leq \left| \tau \int_{\mathbb{R}^d} \Delta \varphi(x) (\rho_\tau^{k+1}(x)^m + \beta \rho_\tau^{k+1}(x)^2) dx \right| \\
& + \left| \tau c_{d,s} \chi \frac{d-2s}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x-y) |x-y|^{2s-d-2} \rho_\tau^{k+1}(x) \rho_\tau^{k+1}(y) dx dy \right| \\
& \leq \tau \|\varphi\|_{W^{2,\infty}(\mathbb{R}^d)} \left(\int_{\mathbb{R}^d} (\rho_\tau^{k+1}(x)^m + \beta \rho_\tau^{k+1}(x)^2) dx \right. \\
& \quad \left. + \frac{\chi}{2} (d-2s) \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(|x-y|) \rho_\tau^{k+1}(x) \rho_\tau^{k+1}(y) dx dy \right) \leq K \tau \|\varphi\|_{W^{2,\infty}(\mathbb{R}^d)}
\end{aligned}$$

where K is an explicit constant, only depending on M, m, s, d, β, χ and ρ^0 , which can be obtained by applying Proposition 4.2, and in particular by combining (2.10) and (4.5). Concerning the second, we have

$$\begin{aligned}
& \left| \frac{1}{2} \int_0^1 (1-\xi)^2 d\xi \int_{\mathbb{R}^d} \nabla^2 \varphi(\vartheta_\xi) (\mathbf{T}_\tau^{k+1}(x) - x)^T \cdot (\mathbf{T}_\tau^{k+1}(x) - x) \rho_\tau^{k+1}(x) dx \right| \\
& \leq \frac{1}{2} \|\varphi\|_{W^{2,\infty}(\mathbb{R}^d)} W_2^2(\rho_\tau^{k+1}, \rho_\tau^k).
\end{aligned}$$

By inserting the latter two estimates in (4.25) we deduce that for every smooth function $\varphi \in W^{2,\infty}(\mathbb{R}^d)$

$$\begin{aligned}
\int_{\mathbb{R}^d} \varphi(x) (\rho_\tau(t+h, x) - \rho_\tau(t, x)) dx & \leq \|\varphi\|_{W^{2,\infty}(\mathbb{R}^d)} \sum_{k=\lceil t/\tau \rceil}^{\lceil (t+h)/\tau \rceil - 1} \left(K \tau + \frac{1}{2} W_2^2(\rho_\tau^{k+1}, \rho_\tau^k) \right) \\
& \leq (K(\tau+h) + 2\tau(\bar{C} + \mathcal{H}_m[\rho^0])) \|\varphi\|_{W^{2,\infty}(\mathbb{R}^d)},
\end{aligned}$$

where \mathcal{H}_m is defined by (2.5) and \bar{C} is defined by (2.11), and where the latter inequality follows again from Proposition 4.2 together with (2.10), and from the basic inequalities $\lceil \frac{t+h}{\tau} \rceil \leq \frac{t+h}{\tau} + 1$, $\lceil \frac{t}{\tau} \rceil \geq \frac{t}{\tau}$. Therefore, fixing $q \in \mathbb{N}$ with $q > 2 + d/2$, so that the continuous embedding given by Morrey's theorem $H^q(\mathbb{R}^d) \subset W^{2,\infty}(\mathbb{R}^d)$ holds with constant Q , we deduce

$$\begin{aligned}
\|\rho_\tau(t+h, \cdot) - \rho_\tau(t, \cdot)\|_{H^{-q}(\mathbb{R}^d)} & = \sup_{\|\varphi\|_{H^q(\mathbb{R}^d)} \leq 1} \int_{\mathbb{R}^d} \varphi(x) (\rho_\tau(t+h, x) - \rho_\tau(t, x)) dx \\
& \leq Q (K(\tau+h) + 2\tau(\bar{C} + \mathcal{H}_m[\rho^0])).
\end{aligned}$$

In particular

$$\limsup_{\tau \rightarrow 0} \|\rho_\tau(t+h, \cdot) - \rho_\tau(t, \cdot)\|_{H^{-q}(\mathbb{R}^d)} \leq QKh. \quad (4.26)$$

The conclusion is similar to the one in [6, Proposition 14]. Let (τ_n) and ρ be the vanishing sequence and the limit function in the statement. Let us consider the set of functions $\{\rho_{\tau_n}(t, \cdot) : t \in [0, T], n \in \mathbb{N}\}$. This is a set of functions having uniformly bounded second moments and uniformly bounded $L^1 \cap L^2(\mathbb{R}^d)$ norm, thanks to the estimates (4.5) and (4.6). Hence, it is relatively compact in $H^{-q}(\mathbb{R}^d)$ by Lemma 4.11 below. Thanks to this fact and to the equicontinuity estimate (4.26), we may apply [1, Proposition 3.3.1] and deduce that there exists a vanishing subsequence (τ_{n_k}) and $\bar{\rho} \in C([0, T]; H^{-q}(\mathbb{R}^d))$ such that

$$\rho_{\tau_{n_k}}(t, \cdot) \rightarrow \bar{\rho}(t, \cdot) \quad \text{in } H^{-q}(\mathbb{R}^d) \quad \text{for every } t \in [0, T]. \quad (4.27)$$

By uniqueness of the limit we have $\bar{\rho} \equiv \rho$ and the above convergence holds along the original sequence (τ_n) .

The conclusion of the proof is split in two cases. We first consider the case $\beta > 0$, and we will take advantage of Corollary (4.7). We observe that for every $f \in H^1(\mathbb{R}^d)$ there holds $|f|^{\frac{2}{m}} \in X_m := W^{\frac{2}{m}, m}(\mathbb{R}^d)$ since $m > 2$: this is shown in [25] along with the estimate $\| |f|^{\frac{2}{m}} \|_{X_m}^m \leq c \|\nabla f\|_{L^2(\mathbb{R}^d)}^2$ for a suitable constant c . Therefore, from (4.16) and Jensen inequality we deduce that

$$\left(\frac{1}{T} \int_0^T \|\rho_{\tau_n}(t, \cdot)\|_{X_m}^2 dt \right)^{m/2} \leq \frac{1}{T} \int_0^T \|\rho_{\tau_n}(t, \cdot)\|_{X_m}^m dt \leq C_T \quad (4.28)$$

where C_T is a constant depending only on $M, m, s, d, \beta, \chi, T$ and the initial datum ρ^0 . This shows that the sequence ρ_{τ_n} is bounded in $L^2((0, T); X_m)$. Since $L^2(\mathbb{R}^d)$ continuously embeds in $H^{-q}(\mathbb{R}^d)$, by the uniform $L^2(\mathbb{R}^d)$ bound deduced from (4.5) and by (4.27) we may apply the dominated convergence theorem and obtain the convergence of ρ_{τ_n} to ρ in $L^2((0, T); H^{-q}(\mathbb{R}^d))$. The latter convergence, together with (4.28) allows for an application of the compactness result in the space $L^2((0, T); L^2(\mathbb{R}^d))$ from [32, Lemma 9] so that we conclude that $\rho_{\tau_n} \rightarrow \rho$ strongly in $L^2((0, T) \times \mathbb{R}^d)$. Here, [32, Lemma 9] is applied by using the Banach triple $Y \subset L^2(\mathbb{R}^d) \subset H^{-q}(\mathbb{R}^d)$, where $Y = X_m \cap L^1(\mathbb{R}^d, (1+|x|^2)dx)$, and where the first embedding is compact by Lemma 4.11. The proof for the case $\beta > 0$ is concluded.

Eventually, let us consider the case $\beta = 0$, $d \geq 2$, $1/2 \leq s < 1$. In this case we change the definition of the Sobolev space X_m and we let $X_m = W^{\frac{1}{m-1}, 2m-2}(\mathbb{R}^d)$, and by invoking Corollary (4.9) instead of Corollary (4.7), we conclude by repeating the same argument that we have used for $\beta > 0$. \square

Lemma 4.11. *Let $m > 2$. The spaces $Y_1 := W^{\frac{2}{m}, m}(\mathbb{R}^d) \cap L^1(\mathbb{R}^d, (1+|x|^2)dx)$ and $Y_2 := W^{\frac{1}{m-1}, 2m-2}(\mathbb{R}^d) \cap L^1(\mathbb{R}^d, (1+|x|^2)dx)$ are compactly embedded into $L^2(\mathbb{R}^d)$. The space $L^2(\mathbb{R}^d) \cap L^1(\mathbb{R}^d, (1+|x|^2)dx)$ is compactly embedded into $H^{-q}(\mathbb{R}^d)$ if $q > d/2$.*

Proof. Let us give the proof for Y_1 (the argument for Y_2 is analogous). Let us consider a sequence $(u_n)_{n \in \mathbb{N}}$ which is bounded in Y_1 , thus in particular it is bounded in $W^{\frac{2}{m}, m}(\mathbb{R}^d)$ and in $L^m \cap L^1(\mathbb{R}^d)$. By fractional Sobolev embedding, since $m > 2$ we have that $W^{\frac{2}{m}, m}(\mathbb{R}^d)$ is embedded compactly in $L^2(B)$ for every ball in \mathbb{R}^d , so that there is $u \in L^2(\mathbb{R}^d)$ and a not relabeled subsequence (u_n) such that $u_n \rightarrow u$ strongly in $L^2(B)$ for every ball B and weakly in $L^m(\mathbb{R}^d)$. Let $\varepsilon > 0$ and choose $B = B_\varepsilon$ to be a large enough ball, such that $\int_{\mathbb{R}^d \setminus B} |u| + \int_{\mathbb{R}^d \setminus B} |u_n| < \varepsilon$ for every $n \in \mathbb{N}$: this is possible thanks to the tightness of the sequence (u_n) , which has uniformly bounded second moments by assumption. We have

$$\|u_n - u\|_{L^2(\mathbb{R}^d)} \leq \|u_n - u\|_{L^1(\mathbb{R}^d)}^{1-\theta} \|u_n - u\|_{L^m(\mathbb{R}^d)}^\theta, \quad \theta := \frac{m}{2m-2}. \quad (4.29)$$

We also have

$$\int_{\mathbb{R}^d} |u_n - u| dx \leq \int_B |u_n - u| dx + \int_{\mathbb{R}^d \setminus B} |u_n - u| dx \leq \int_B |u_n - u| dx + \varepsilon$$

and since $u_n \rightarrow u$ strongly in $L^1(B)$, the arbitrariness of ε shows that $u_n \rightarrow u$ strongly in $L^1(\mathbb{R}^d)$ as well. Therefore, the boundedness of (u_n) in $L^m(\mathbb{R}^d)$ and (4.29) show that $u_n \rightarrow u$ strongly in $L^2(\mathbb{R}^d)$.

Similarly, let $(u_n)_{n \in \mathbb{N}}$ be a bounded sequence in $L^2(\mathbb{R}^d) \cap L^1(\mathbb{R}^d, (1 + |x|^2) dx)$, and given $\varepsilon > 0$ let as above $B = B_\varepsilon$ such that $\int_{\mathbb{R}^d \setminus B} |u| + \int_{\mathbb{R}^d \setminus B} |u_n| < \varepsilon$ for every $n \in \mathbb{N}$. By Sobolev embedding, $H^q(B)$ compactly embeds into $L^2(B)$ and then (by Schauder's theorem) $L^2(B)$ compactly embeds in the dual space $H^q(B)^*$, therefore up to subsequences we have $u_n \rightarrow u$ weakly in $L^2(\mathbb{R}^d)$ and strongly in $H^q(B)^*$. We have

$$\begin{aligned} \|u_n - u\|_{H^{-q}(\mathbb{R}^d)} &\leq \sup_{\|\varphi\|_{H^q(\mathbb{R}^d)} \leq 1} \left| \int_B \varphi(u_n - u) \right| + \sup_{\|\varphi\|_{H^q(\mathbb{R}^d)} \leq 1} \left| \int_{\mathbb{R}^d \setminus B} \varphi(u_n - u) \right| \\ &\leq \sup_{\|\varphi\|_{H^q(\mathbb{R}^d)} \leq 1} \|\varphi\|_{H^q(B)} \|u_n - u\|_{H^q(B)^*} + C\varepsilon \leq \|u_n - u\|_{H^q(B)^*} + C\varepsilon, \end{aligned}$$

where we have also used the continuous embedding $\|\cdot\|_{L^\infty(\mathbb{R}^d)} \leq C\|\cdot\|_{H^q(\mathbb{R}^d)}$, since $q > d/2$. Taking the limit as $n \rightarrow +\infty$, since $u_n \rightarrow u$ strongly in $H^q(B)^*$ and since ε is arbitrary, we deduce that $\|u_n - u\|_{H^{-q}(\mathbb{R}^d)} \rightarrow 0$. \square

We are ready to prove Theorem (1.3), recalling that by a gradient flow solution $\rho = \rho(t, x)$ to (1.1) we mean a weak solution according to (1.6) which is a limit of the JKO scheme, i.e., ρ is a limit function (obtained from Proposition 4.3 along a vanishing sequence $(\tau_n)_{n \in \mathbb{N}} \subset (0, 1)$) of the piecewise constant interpolations ρ_τ (defined by (4.7)) constructed from a sequence $(\rho_\tau^k)_{k \geq 0}$ of discrete minimizers from (4.3).

Proof of Theorem 1.3. We apply Proposition 4.10: we have $\rho_{\tau_n} \rightarrow \rho$ strongly in $L^2((0, T) \times \mathbb{R}^d)$. In particular, up to extracting a further subsequence, we have the pointwise a.e. space-time convergence of ρ_{τ_n} to ρ and thus of $\rho_{\tau_n}^m$ to ρ^m . Thanks to the Sobolev embedding $H^1(\mathbb{R}^d) \subset L^{\frac{2d}{d-2}}(\mathbb{R}^d)$ if $d \geq 3$ (resp. $H^1(\mathbb{R}^d) \subset L^4(\mathbb{R}^d)$ if $d = 1, 2$) we deduce from (4.16) if $\beta > 0$ and from (4.24) if $\beta = 0$ that the sequence $(\rho_{\tau_n}^m)_{n \in \mathbb{N}}$ is also bounded in $L^{\frac{d}{d-2}}((0, T) \times \mathbb{R}^d)$ if $d \geq 3$ (resp. in $L^2((0, T) \times \mathbb{R}^d)$ if $d = 1, 2$), hence by interpolation we also find that $(\rho_{\tau_n}^2)_{n \in \mathbb{N}}$ is bounded in the same space. Thus up the extraction of one more subsequence we get

$$\begin{aligned} & \lim_{n \rightarrow +\infty} \int_0^T \int_{\mathbb{R}^d} \eta(t) \Delta \varphi(x) (\rho_{\tau_n}(t, x)^m + \beta \rho_{\tau_n}(t, x)^2) dx dt \\ &= \int_0^T \int_{\mathbb{R}^d} \eta(t) \Delta \varphi(x) (\rho(t, x)^m + \beta \rho(t, x)^2) dx dt, \end{aligned}$$

for every $\eta \in C_c^\infty(0, T)$ and every $\varphi \in C_c^\infty(\mathbb{R}^d)$. The weak $L^2(\mathbb{R}^d)$ convergence of $\rho_{\tau_n}(t, \cdot)$ to $\rho(t, \cdot)$ is sufficient for obtaining

$$\begin{aligned} & \lim_{n \rightarrow +\infty} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho_{\tau_n}(t, x) \rho_{\tau_n}(t, y) dx dy \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(t, x) \rho(t, y) dx dy \end{aligned}$$

for every t , by making use of the same argument of the proof of Proposition 4.1. By dominated convergence, the associated time integrals on $(0, T)$ also converge. Finally, we have for every $\eta \in C_c^\infty(0, T)$, by the convergence properties of minimizing movements, see for instance [1, Theorem 11.1.6],

$$\lim_{n \rightarrow +\infty} \frac{1}{\tau_n} \int_0^T \eta(t) \int_{\mathbb{R}^d} (\mathsf{T}_{\tau_n}(t, x) - x) \cdot \nabla \varphi(x) \rho_{\tau_n}(t, x) dx dt = \int_0^T \partial_t \eta(t) \int_{\mathbb{R}^d} \varphi(x) \rho(t, x) dx dt,$$

where $\mathsf{T}_{\tau_n}(t, \cdot)$ is the unique optimal transport map from $\rho_{\tau_n}(t, x) dx$ to $\rho_{\tau_n}^{\lceil t/\tau_n \rceil - 1}(x) dx$. Recalling the definition of ρ_τ as the piecewise constant interpolation of discrete minimizers, by writing (4.8) for ρ_τ we have

$$\begin{aligned} & \frac{1}{\tau_n} \int_{\mathbb{R}^d} (\mathsf{T}_{\tau_n}(t, x) - x) \cdot \nabla \varphi(x) \rho_{\tau_n}(t, x) dx = - \int_{\mathbb{R}^d} \Delta \varphi(x) (\rho_{\tau_n}(t, x)^m + \beta \rho_{\tau_n}(t, x)^2) dx \\ &+ \frac{(d-2s)\chi c_{d,s}}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \zeta(x) - \nabla \zeta(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho_{\tau_n}(t, x) \rho_{\tau_n}(t, y) dx dy \end{aligned}$$

for every $t > 0$. By multiplying the latter by $\eta \in C_c^\infty(0, T)$ and by integrating on $(0, T)$, we may therefore pass to the limit along the sequence $(\tau_n)_n$ and conclude that ρ is a weak solution to problem (1.1). \square

Let us collect some properties of the constructed solution.

Proposition 4.12. *Let $\beta \geq 0$. If $\beta = 0$, assume in addition that $d \geq 2$, $1/2 \leq s < 1$. Let $\rho^0 \in \mathcal{Y}_{M,2}$. Let $(\rho_\tau^k)_{k \geq 0}$ be a sequence of discrete minimizers defined by (4.3), let ρ_τ be the piecewise constant interpolation defined by (4.7), and let ρ be a limit function obtained from Proposition 4.3 along a vanishing sequence $(\tau_n)_{n \in \mathbb{N}} \subset (0, 1)$. Then the following properties hold.*

- (i) *The function $[0, +\infty) \ni t \mapsto \rho(t, \cdot) \in \mathcal{Y}_{M,2}$ is absolutely continuous with respect to the Wasserstein distance W_2 .*
- (ii) *The following $L^m(\mathbb{R}^d)$ estimate holds for every $t > 0$*

$$\frac{1}{2(m-1)} \int_{\mathbb{R}^d} (\rho(t, x))^m dx \leq \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0(x))^m dx + \beta \int_{\mathbb{R}^d} (\rho^0(x))^2 dx + \bar{C},$$

where \bar{C} is defined by (2.11) and (2.8).

- (iii) *if $\beta > 0$, then $\rho^{m/2} \in L^2((0, T); H^1(\mathbb{R}^d))$ for every $T > 0$ along with the estimate*

$$\frac{4}{m} \int_0^T \int_{\mathbb{R}^d} |\nabla(\rho(t, x))^{m/2}|^2 dx dt \leq C_1^* + TC_2^* + TC_3^* \chi_s \left(\frac{\chi(1-s)}{2\beta} \right)^{\frac{1-s}{s}},$$

where C_i^* are the explicit constants defined in the proof of Proposition 4.7.

- (iv) *for every $t > 0$ there holds*

$$\begin{aligned} \int_{\mathbb{R}^d} |x|^2 \rho(t, x) dx &\leq 4t \mathcal{F}_s[\rho^0] + 2 \int_{\mathbb{R}^d} |x|^2 \rho^0(x) dx \\ &\leq 4t \left(\frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0(x))^m dx + \beta \int_{\mathbb{R}^d} (\rho^0(x))^2 dx \right) + 2 \int_{\mathbb{R}^d} |x|^2 \rho^0(x) dx. \end{aligned}$$

- (v) *if $\beta = 0$, then $\rho^{m-1} \in L^2((0, T); H^1(\mathbb{R}^d))$ for every $T > 0$ along with the estimate*

$$\int_0^T \int_{\mathbb{R}^d} |\nabla(\rho(t, x))^{m-1}|^2 dx dt \leq C_1^{**} + TC_2^{**},$$

where C_i^{**} are the explicit constants appearing in Corollary 4.9.

Proof. Point (i) was shown in Proposition 4.3. Point (ii) follows from the uniform $L^m(\mathbb{R}^d)$ bound from (4.5) along with the narrow lower semicontinuity of the L^m norm. Points (iii) and (iv) respectively follow from the uniform bounds (4.16) and (4.6), again by lower semicontinuity properties. Similarly, point (v) follows from (4.24). \square

We conclude this section by proving Theorem 1.4. Before giving the proof, we include a couple of technical lemmas, whose proofs are postponed to the Appendix.

Lemma 4.13. *Let $\rho \in L^1 \cap L^2(\mathbb{R}^d)$. Then for every $\varphi \in C_c^\infty(\mathbb{R}^d)$ there holds*

$$\begin{aligned} \lim_{s \downarrow 0} (d - 2s) c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(x) \rho(y) dx dy \\ = \int_{\mathbb{R}^d} \rho^2(x) \Delta \varphi(x) dx. \end{aligned}$$

A generalization of the previous lemma is the following

Lemma 4.14. *Let $\rho \in L^1 \cap L^2(\mathbb{R}^d)$ and let $(\rho_s)_{s \in (0, 1/2)} \subset L^1 \cap L^2(\mathbb{R}^d)$ be a family of functions such that $\rho_s \rightarrow \rho$ in $L^2(\mathbb{R}^d)$ as $s \downarrow 0$ and $\|\rho_s\|_{L^1(\mathbb{R}^d)} = \|\rho\|_{L^1(\mathbb{R}^d)}$ for every $s \in (0, 1/2)$. Then for every $\varphi \in C_c^\infty(\mathbb{R}^d)$ there holds*

$$\begin{aligned} \lim_{s \downarrow 0} (d - 2s) c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho_s(x) \rho_s(y) dx dy \\ = \int_{\mathbb{R}^d} \rho^2(x) \Delta \varphi(x) dx. \end{aligned}$$

We are ready for the proof of our last result.

Proof of Theorem 1.4. Let $s = s_n$. Let $\varphi \in C_c^\infty(\mathbb{R}^d)$. Since $[0, +\infty) \ni t \mapsto \rho_n(t, \cdot)$ is absolutely continuous with respect to W_2 as recalled in point (i) of Proposition 4.12, the map $t \mapsto \int_{\mathbb{R}^d} \varphi(x) \rho_n(t, x) dx$ is absolutely continuous on $[0, +\infty)$ and we may write the time integrated version of (1.6), i.e.,

$$\begin{aligned} \int_{\mathbb{R}^d} (\rho_n(t_2, x) - \rho_n(t_1, x)) \varphi(x) dx &= \int_{t_1}^{t_2} \int_{\mathbb{R}^d} \Delta \varphi(x) (\rho_n^m(t, x) + \beta \rho_n^2(t, x)) dx dt \\ &\quad - \frac{\chi}{2} (d - 2s_n) c_{d,s_n} \int_{t_1}^{t_2} dt \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) \\ &\quad \times |x - y|^{2s_n-d-2} \rho_n(t, x) \rho_n(t, y) dx dy, \end{aligned}$$

for every $0 \leq t_1 < t_2 < +\infty$. We estimate the last term as done in (2.9)-(2.10), obtaining

$$\begin{aligned} & \left| \frac{\chi}{2} (d - 2s_n) c_{d,s_n} \int_{t_1}^{t_2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) \right. \\ & \quad \left. \times |x - y|^{2s_n - d - 2} \rho_n(t, x) \rho_n(t, y) dx dy dt \right| \\ & \leq \frac{\chi}{2} (d - 2s_n) \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \int_{\mathbb{R}^d} (K_{s_n} * \rho_n)(t, x) \rho_n(t, x) dx dt \\ & \leq d \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \left(\bar{C}(\chi, m, s_n, d, M) + \frac{1}{2(m-1)} \int_{\mathbb{R}^d} \rho_n^m(t, x) dx \right) dt, \end{aligned}$$

where \bar{C} is defined by (2.11) and (2.8), therefore we have by $L^1 - L^2 - L^m$ interpolation

$$\begin{aligned} & \left| \int_{\mathbb{R}^d} (\rho_n(t_2, x) - \rho_n(t_1, x)) \varphi(x) dx \right| \leq \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \int_{\mathbb{R}^d} (\rho_n^m(t, x) + \beta \rho_n^2(t, x)) dx dt \\ & \quad + \frac{\chi}{2} (d - 2s_n) \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \int_{\mathbb{R}^d} (K_{s_n} * \rho_n)(t, x) \rho_n(t, x) dx, \\ & \leq \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \int_{\mathbb{R}^d} \rho_n^m(t, x) dx dt \\ & \quad + \beta \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \left(\frac{m-2}{m-1} M + \frac{1}{m-1} \int_{\mathbb{R}^d} \rho_n^m(t, x) dx \right) dt \\ & \quad + d \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \int_{t_1}^{t_2} \left(\bar{C}(\chi, m, s_n, d, M) + \frac{1}{2(m-1)} \int_{\mathbb{R}^d} \rho_n^m(t, x) dx \right) dt. \end{aligned}$$

It is immediate to check that $\sup_{s \in (0, 1/2)} \bar{C}(\chi, m, s, d, M) < +\infty$, therefore by including the estimate in point (ii) of Proposition 4.12 we deduce that

$$\left| \int_{\mathbb{R}^d} (\rho_n(t_2, x) - \rho_n(t_1, x)) \varphi(x) dx \right| \leq \tilde{C}(t_2 - t_1) \|\varphi\|_{W^{2, \infty}(\mathbb{R}^d)}$$

where \tilde{C} is a suitable constant, depending on m, d, χ, β, M and ρ^0 , but not on n . We take $q \in \mathbb{N}$, $q > 2 + d/2$, so that we have the continuous embedding $H^q(\mathbb{R}^d) \subset W^{2,\infty}(\mathbb{R}^d)$ with embedding constant Q , and we deduce the time equi-Lipschitz estimate

$$\|\rho_n(t_2, \cdot) - \rho_n(t_1, \cdot)\|_{H^{-q}(\mathbb{R}^d)} \leq Q\tilde{C}(t_2 - t_1) \quad \text{for every } n \in \mathbb{N}.$$

Letting $T > 0$, we repeat the same arguments of the proof of Proposition 4.10: indeed, the family of functions $\{\rho_n(t, \cdot) : t \in [0, T], n \in \mathbb{N}\}$ has uniformly bounded second moments and $L^1 \cap L^2(\mathbb{R}^d)$ norms, which is seen by applying to ρ_n the estimates of points (ii) and (iv) of Proposition 4.12 (as already noticed, the right hand side of (ii) can be estimated uniformly with respect to $n \in \mathbb{N}$). Therefore by Lemma 4.11 such a family of functions is relatively compact in $H^{-q}(\mathbb{R}^d)$, so that in view of the above equi-Lipschitz estimate we may apply [1, Lemma 3.3.1] to find $\rho \in C^0([0, T]; H^{-q}(\mathbb{R}^d))$ such that for every $t \in [0, T]$ there holds $\rho_n(t, \cdot) \rightarrow \rho(t, \cdot)$ in $H^{-q}(\mathbb{R}^d)$ as $n \rightarrow +\infty$. By the dominated convergence theorem, we also get $\rho_n \rightarrow \rho \in L^2((0, T); H^{-q}(\mathbb{R}^d))$ as $n \rightarrow +\infty$: here, the dominating function for showing that $\int_0^T \|\rho_n(t, \cdot) - \rho(t, \cdot)\|_{H^{-q}(\mathbb{R}^d)}^2 dt \rightarrow 0$ as $n \rightarrow +\infty$ is obtained by using the continuous embedding of $L^2(\mathbb{R}^d)$ into $H^{-q}(\mathbb{R}^d)$ and the estimate in point (ii) of Proposition 4.12, where again the right hand side is uniformly bounded with respect to n . The same estimate and the estimate in point (iv) of the same Proposition imply that $\rho_n(t, \cdot)$ converges weakly to $\rho(t, \cdot)$ in $L^1 \cap L^m(\mathbb{R}^d)$, therefore the weak lower semicontinuity of the L^m norm also shows that $\rho \in L^\infty((0, +\infty); L^m(\mathbb{R}^d))$.

The estimate in point (iii) of Proposition 4.12, applied to ρ_n , implies

$$\sup_{n \in \mathbb{N}} \int_0^T \int_{\mathbb{R}^d} |\nabla(\rho_n(t, x))^{m/2}|^2 dx dt \leq \frac{m}{4} \sup_{n \in \mathbb{N}} \left(C_1^* + TC_2^* + TC_3^* \chi s_n \left(\frac{\chi(1-s_n)}{2\beta} \right)^{\frac{1-s_n}{s_n}} \right),$$

and the supremum in the right hand side is finite, thanks to the crucial assumption $\beta \geq \chi/2$ (as observed in Remark 4.6). Therefore, we may reason as done in the proof of Proposition 4.10 to get that the sequence (ρ_n) enjoys a uniform $L^2((0, T); W^{2/m, m}(\mathbb{R}^d))$ bound. Similarly, in view of point (iv) of Proposition 4.12, the sequence (ρ_n) is also uniformly bounded in $L^2((0, T); L^1(\mathbb{R}^d, (1+|x|^2) dx))$. By the same argument at the end of the proof of Proposition 4.10, we have $\rho_n \rightarrow \rho$ strongly in $L^2((0, T) \times \mathbb{R}^d)$.

Let us conclude by passing to the limit in the equation. Let $\varphi \in C_c^\infty(\mathbb{R}^d)$ and $\eta \in C_c^\infty((0, T))$. By definition of weak solution, for each $n \in \mathbb{N}$ we have that ρ_n satisfies

$$\begin{aligned} - \int_0^T \int_{\mathbb{R}^d} \rho_n(t, x) \partial_t \eta(t) \varphi(x) dx dt &= \int_0^T \int_{\mathbb{R}^d} \eta(t) \Delta \varphi(x) (\rho_n(t, x)^m + \beta \rho_n(t, x)^2) dx dt \\ &\quad - \frac{(d-2s_n)c_{d,s_n}\chi}{2} \int_0^T \eta(t) \int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{(\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x-y)}{|x-y|^{d+2-2s_n}} \rho_n(t, x) \rho_n(t, y) dx dy dt. \end{aligned} \quad (4.30)$$

Since $\rho_n \rightarrow \rho$ strongly in $L^2((0, T) \times \mathbb{R}^d)$, up to taking another subsequence we have $\rho_n(t, \cdot) \rightarrow \rho(t, \cdot)$ in $L^2(\mathbb{R}^d)$ for a.e. $t \in (0, T)$. An application of Lemma 4.14 entails therefore

$$\begin{aligned} & \lim_{n \rightarrow +\infty} \frac{\chi}{2} (d - 2s_n) c_{d, s_n} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) \\ & \quad \times |x - y|^{2s_n - d - 2} \rho_n(t, x) \rho_n(t, y) dx dy \\ & = \frac{\chi}{2} \int_{\mathbb{R}^d} \rho^2(t, x) \Delta \varphi(x) dx \quad \text{for a.e. } t \in (0, T). \end{aligned}$$

After multiplying by η and integrating on $(0, T)$, the time integral passes to the limit by dominated convergence: a dominating function is obtained by the usual estimates of the form (2.9)-(2.10), yielding for a.e. $t \in (0, T)$

$$\begin{aligned} & \left| \frac{\chi}{2} (d - 2s_n) c_{d, s_n} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s_n - d - 2} \rho_n(t, x) \rho_n(t, y) dx dy \right| \\ & \leq d \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \|\eta\|_{L^\infty(\mathbb{R})} \left(\bar{C}(\chi, m, s_n, d, M) + \frac{1}{2(m-1)} \int_{\mathbb{R}^d} \rho_n^m(t, x) dx \right) \\ & \leq d \|\nabla^2 \varphi\|_{L^\infty(\mathbb{R}^d)} \|\eta\|_{L^\infty(\mathbb{R})} \left(2\bar{C} + \frac{1}{m-1} \int_{\mathbb{R}^d} (\rho^0(x))^m dx + \beta \int_{\mathbb{R}^d} (\rho^0(x))^2 dx \right), \end{aligned}$$

where we have also used point (ii) of Proposition 4.12, and $\bar{C} = \bar{C}(\chi, m, s_n, d, M)$ stays bounded as $n \rightarrow +\infty$.

Eventually, we take advantage of the previously obtained uniform $L^2((0, T); W^{2/m, m}(\mathbb{R}^d))$ estimate: as in the proof of Theorem 1.3, by Sobolev embedding it implies that $(\rho_n^m)_{n \in \mathbb{N}}$ is also uniformly bounded in $L^{\frac{d}{d-2}}((0, T) \times \mathbb{R}^d)$ if $d \geq 3$ (and in $L^2((0, T) \times \mathbb{R}^d)$ if $d = 1, 2$). Therefore, up to subsequences, we have $\rho_n^m \rightarrow \rho^m$ weakly in $L^{\frac{d}{d-2}}((0, T) \times \mathbb{R}^d)$ if $d \geq 3$ (weakly in $L^2((0, T) \times \mathbb{R}^d)$ if $d = 1, 2$) which allow to pass to the limit in the other two terms of (4.30). \square

5. Some qualitative properties of solutions

In this section we present some numerical simulations about the evolution problem (1.1), using the scheme developed in [13]. The time evolution is shown in Fig. 2 and 3 for different initial data in one dimension, providing a numerical illustration of the expected asymptotic behavior, i.e., solutions approaching the unique stationary states. In general, if s is not too close to zero, the stationary states are reached quickly, otherwise the

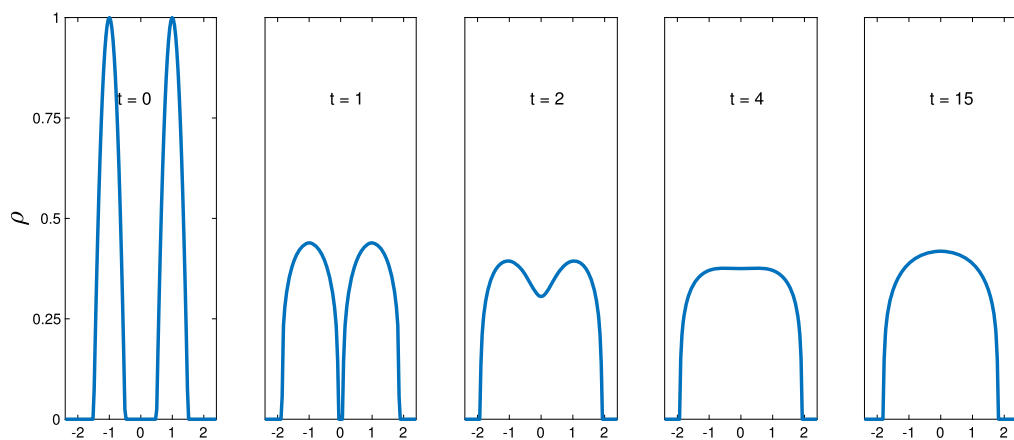


Fig. 2. The evolution of the solution starting with two bumps with parameters $m = 3, \chi = 1, s = 0.1$ and $\beta = 0.2$, reaching the stationary state reasonably fast.

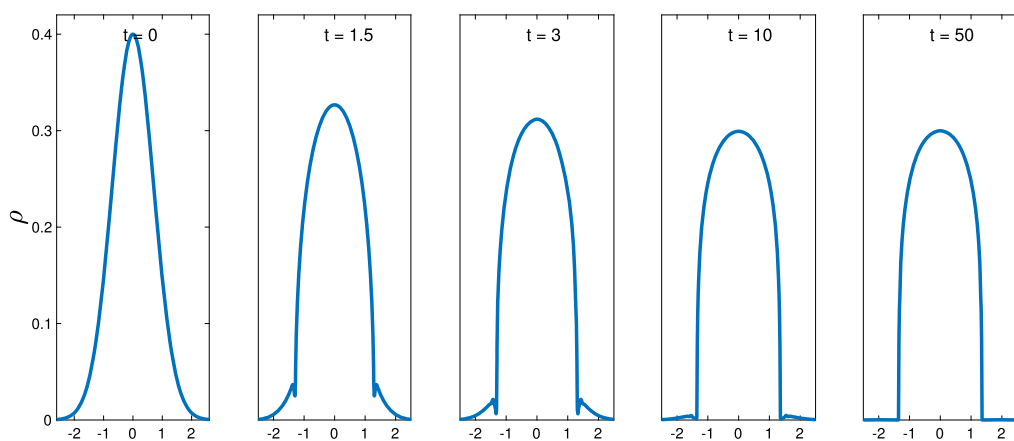


Fig. 3. The evolution starting with a rescaled Gaussian, with $m = 3, s = 0.08, \chi = 1$ and $\beta = 0.4$, where the solution does converge to the expected stationary state.

convergence may take longer with the appearance of “disturbances” near the boundary of the support as in Fig. 3.

A big open problem concerning the Cauchy problem (1.1) is the *uniqueness of the solution*. Assuming that this property holds true, by the rotationally invariant property of the main equation (1.1) it follows that for a given radial initial datum $\rho_0 = \rho_0(|x|)$ we have that the density solution ρ is radial w.r.t. x , *i.e.* $\rho = \rho(|x|, t)$. But the property of being radially decreasing may not be preserved during the evolution, even with initial data (and limiting steady states) sharing this property. The following counterexample is an adaptation of the one contained in [19, Proposition 4.3]. Set

$$\rho_{0,\varepsilon}(x) = \varphi_\varepsilon * (\delta_0 + \varepsilon^\alpha \mathbb{1}_{B(0,1)}) = \varphi_\varepsilon(x) + \varepsilon^\alpha \varphi_\varepsilon * \mathbb{1}_{B(0,1)}, \quad (5.1)$$

being $\varepsilon > 0$, φ a mollifier with mass 1 supported in the ball $B(0, 1)$, being $\varphi_\varepsilon(x) = \frac{1}{\varepsilon^d} \varphi\left(\frac{x}{\varepsilon}\right)$ and $\alpha > d + 2$. Assume that there exists a radial solution $\rho(x, t)$ to (1.1) with datum $\rho_{0,\varepsilon}$ and suppose we know that the solution $\rho(x, t)$ is smooth enough up to $t = 0$. Then it is possible to show that ρ is not radially decreasing. Indeed, it is immediate to see that for ε small and $\varepsilon < |x| < 1 - \varepsilon$ we have

$$\rho_{0,\varepsilon}(x) = \varepsilon^\alpha,$$

while $\rho_{0,\varepsilon}$ is supported in the ball $B(0, 1 + \varepsilon)$. Taking two points x_1, x_2 such that $\varepsilon < |x_1| < |x_2| < 1 - \varepsilon$ and taking into account that $\rho_{0,\varepsilon}$ is constant in the interval $(\varepsilon, 1 - \varepsilon)$, we have for $|x_1| \leq |x| \leq |x_2|$

$$\partial_t \rho(x, 0) = -\rho_{0,\varepsilon}(x) \Delta(K_s * \rho_{0,\varepsilon}). \quad (5.2)$$

Now, observe that since $\varphi_\varepsilon \rightarrow \delta_0$ in $\mathcal{D}'(\mathbb{R}^d)$ for $\varepsilon \rightarrow 0$, we have

$$K_s * \rho_{0,\varepsilon} \rightarrow K_s \quad \text{in } \mathcal{D}'(\mathbb{R}^d).$$

as $\varepsilon \rightarrow 0$. But

$$\Delta(K_s * \rho_{0,\varepsilon}) = K_s * \Delta \rho_{0,\varepsilon} = K_s * \Delta \varphi_\varepsilon + K_s * \varepsilon^\alpha (\Delta \varphi_\varepsilon * \mathbb{1}_{B(0,1)}). \quad (5.3)$$

Since

$$\Delta \varphi_\varepsilon \rightarrow \Delta \delta_0 \quad \text{in } \mathcal{D}'(\mathbb{R}^d)$$

and $\Delta \delta_0$ is supported at 0, taking a cutoff function η_ε such that $\eta_\varepsilon = 1$ in a $B(0, \varepsilon)$ we find

$$(K_s * \Delta \rho_{0,\varepsilon})(x) = \int_{\mathbb{R}^d} (\Delta \varphi_\varepsilon(y)) \eta_\varepsilon(y) K_s(x-y) dy \rightarrow \langle \Delta \delta_0, \eta_\varepsilon K_s(x-\cdot) \rangle = \Delta(\eta_\varepsilon K_s(x-\cdot))(0)$$

as $\varepsilon \rightarrow 0$ and an easy computation shows that $\Delta(\eta_\varepsilon K_s(x-\cdot))(0) = \Delta K_s(x)$. Now, since

$$|\Delta \varphi_\varepsilon(x)| = \varepsilon^{-d-2} |\Delta \varphi(x/\varepsilon)| \leq C \varepsilon^{-d-2},$$

we have

$$\|\Delta \varphi_\varepsilon * \mathbb{1}_{B(0,1)}\|_{L^\infty} \leq \|\Delta \varphi_\varepsilon\|_{L^\infty} \|\mathbb{1}_{B(0,1)}\|_{L^1} \leq C \varepsilon^{-d-2}$$

and by Young inequality

$$\|\Delta \varphi_\varepsilon * \mathbb{1}_{B(0,1)}\|_{L^1} \leq \|\Delta \varphi_\varepsilon\|_{L^1} \|\mathbb{1}_{B(0,1)}\|_{L^1} = C \varepsilon^{-2} \|\Delta \varphi\|_{L^1}.$$

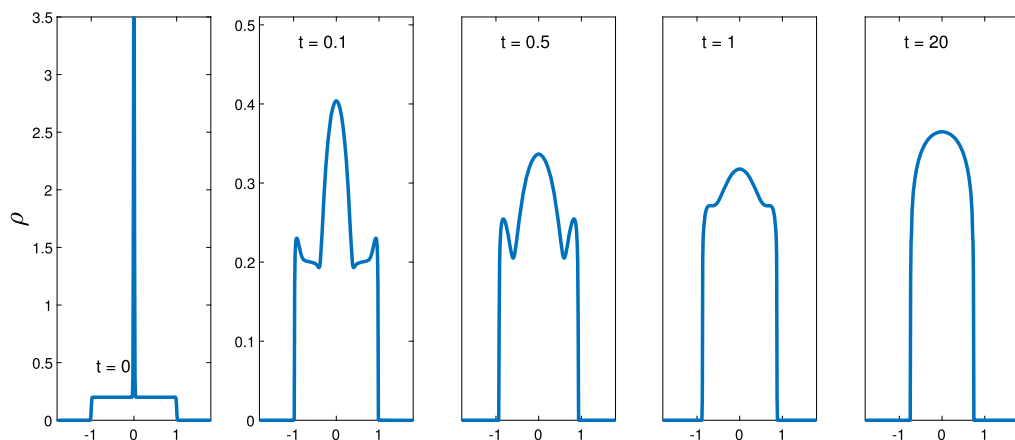


Fig. 4. Numerical demonstration of the fact that radially decreasing initial data does not necessarily remain radially decreasing. The initial condition is the one in (5.1), with the parameters $m = 3$, $\chi = 1$, $s = 0.1$ and $\beta = 0$.

Therefore,

$$\begin{aligned} |K_s * (\Delta\varphi_\varepsilon * \mathbb{1}_{B(0,1)})| &\leq C(\|\Delta\varphi_\varepsilon * \mathbb{1}_{B(0,1)}\|_{L^\infty} + \|\Delta\varphi_\varepsilon * \mathbb{1}_{B(0,1)}\|_{L^1}) \\ &\leq C\varepsilon^{-d-2}(1 + \varepsilon^N). \end{aligned}$$

Thus choosing $\alpha > d + 2$, from (5.3) we have

$$\Delta(K_s * \rho_{0,\varepsilon})(x) \rightarrow \Delta K_s(x) = c(d, s)|x|^{2s-d-2},$$

as $\varepsilon \rightarrow 0$. This implies that for ε small,

$$\Delta(K_s * \rho_{0,\varepsilon})(x_1) > \Delta(K_s * \rho_{0,\varepsilon})(x_2),$$

hence (5.2) gives (recalling that $\rho_{0,\varepsilon}$ is radially decreasing)

$$\partial_t \rho(x_1, 0) < \partial_t \rho(x_2, 0),$$

meaning that the radially decreasing monotonicity is *not* preserved for small times. The non-monotonicity of the solution is shown in the simulation from Fig. 4.

Fig. 5 shows a further simulation, which takes into account a characteristic function of a symmetric interval as initial datum: also in this case the radial monotonicity is not preserved.

6. Comments, extensions and open problems

- As mentioned in the previous section, an open problem concerns the *uniqueness* of solutions, which would give radially of solutions with radial initial data as a direct

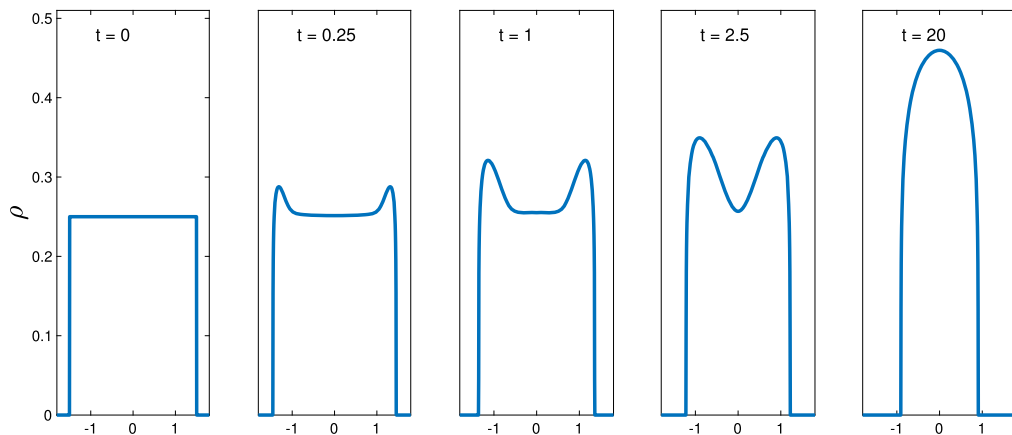


Fig. 5. Another example: again the radially decreasing initial datum does not remain radially decreasing. The initial condition is $\rho_0(X) = \frac{1}{4} \mathbb{1}_{|x| < 3/2}$, with the parameters $m = 3, \chi = 1, s = 0.1$ and $\beta = 0$.

consequence.

- A second open problem is to rigorously prove that every solution to the evolution problem (1.1) does converge to the unique stationary state provided by Theorem 1.1. We mention that a similar result is available in the two dimensional setting, in the case of aggregation with the Newtonian potential instead of the Riesz potential, with $\beta = 0$ and $m > 1$ (i.e., diffusion-dominated regime), see [11].

- Concerning Theorem 1.4, uniqueness of the distributional solution for the Cauchy problem for (1.7) with $\beta \geq \chi/2$ is known under additional conditions. For instance, according to the classical result by [26], uniqueness holds among distributional solutions that are essentially bounded on any strip $\mathbb{R}^d \times (\tau, T)$, for all $T > 0$ and $\tau \in (0, T)$. Therefore, in order to obtain a unique limit as $s \rightarrow 0$ for families of gradient flow solutions ρ_s to (1.1), further a-priori L^∞ bounds (uniformly in s) should be established for ρ_s .

- Another interesting open problem is to show that the family of solutions ρ_s to problem (1.1) converges as $s \rightarrow 0$ to a solution (in an appropriate sense) to the equation (1.7) even in the case $\beta < \chi/2$. We notice that such equation has the form $\partial_t \rho = \Delta \varphi(\rho)$, and if $\beta < \chi/2$ the nonlinearity φ is nonmonotone and equation (1.7) is of *forward-backward* type, with the *unstable* phase given by the interval $[0, (\frac{\chi-2\beta}{m})^{1/(m-2)}]$ and the *stable* phase by $[(\frac{\chi-2\beta}{m})^{1/(m-2)}, +\infty)$. The nontrivial zero of φ is

$$\rho = \left(\frac{\chi - 2\beta}{2} \right)^{1/(m-2)}$$

which coincides exactly with the height of the minimizer of the free energy limit functional \mathcal{F}_0 given in (3.2). If $\beta < \chi/2$, we would like to consider equation (1.7) as a *singular*

limit as $s \rightarrow 0$ of the main equation in (1.1). An existence theory for equation (1.7) supplemented with an initial condition $\rho(x, 0) = \rho_0(x)$ could be given in the setting of Young measure solutions, see for instance [27], where such notion of solution is recovered for cubic-like nonlinearities φ as vanishing limit as $\varepsilon \rightarrow 0$ of a third order pseudo-parabolic regularization $\partial_t \rho = \Delta \varphi(\rho) + \varepsilon \Delta \partial_t \rho$, see also [28], [33] and the references therein. It would be interesting to show that even a weak limit ρ as $s \rightarrow 0$ of a family of densities ρ_s solving the equation (1.1) in the sense of Theorem 1.3 fits the above mentioned existence theory.

Appendix

We provide the proof of Lemma 2.1, Lemma 4.13 and Lemma 4.14.

Proof of Lemma 2.1. First of all, $K_s * \rho \in L^\infty(\mathbb{R}^d)$ by [12, Lemma 1], and let us observe that, for each $i = 1, \dots, d$, $K_s * \rho_{x_i} \in L^1_{loc}(\mathbb{R}^d)$. Indeed, K_s can be written as the sum of two functions, $K_s(x) = \mathbb{1}_{B_1}(x)K_s(x) + (1 - \mathbb{1}_{B_1}(x))K_s(x)$, supported on the unit ball $B_1(x)$ around x and its complement. By Young convolution inequality we have $(\mathbb{1}_{B_1}K_s) * \rho_{x_i} \in L^1(\mathbb{R}^d)$ and $((1 - \mathbb{1}_{B_1})K_s) * \rho_{x_i} \in L^\infty(\mathbb{R}^d)$. Moreover, for any smooth compactly supported test function φ we have $(K_s * \rho_{x_i})\varphi \in L^1(\mathbb{R}^d)$ and we have, by the symmetry of K_s ,

$$\begin{aligned} \int_{\mathbb{R}^d} (K_s * \rho_{x_i})\varphi \, dx &= c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \frac{\rho_{x_i}(y)\varphi(x)}{|x-y|^{d-2s}} \, dx \, dy \\ &= c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \frac{\rho_{x_i}(x)\varphi(y)}{|x-y|^{d-2s}} \, dx \, dy = \int_{\mathbb{R}^d} \rho_{x_i}(K_s * \varphi) \, dx \end{aligned} \quad (6.1)$$

and similarly

$$\int_{\mathbb{R}^d} (K_s * \varphi_{x_i})\rho \, dx = \int_{\mathbb{R}^d} (K_s * \rho)\varphi_{x_i} \, dx. \quad (6.2)$$

If we take a large ball B_R centered at the origin, an integration by parts leads to

$$\int_{B_R} (K_s * \varphi)\rho_{x_i} \, dx = - \int_{B_R} \rho(x)(K_s * \varphi_{x_i}) \, dx + \int_{\partial B_R} \rho(K_s * \varphi)\nu_i \, d\sigma \quad (6.3)$$

and since $|(K_s * \varphi)(x)| \leq C/|x|^{d-2s}$ for large x (as φ is compactly supported), thanks to the continuity and boundedness of ρ we have if $s < 1/2$

$$\int_{\partial B_R} \rho(K_s * \varphi)\nu_i \, d\sigma \leq C\|\rho\|_{L^\infty(\mathbb{R}^d)}R^{2s-1} \rightarrow 0 \quad \text{as } R \rightarrow +\infty,$$

along with

$$\int_{\tilde{B}_R} (K_s * \varphi) \rho_{x_i} dx \rightarrow \int_{\mathbb{R}^d} (K_s * \varphi) \rho_{x_i} dx \quad \text{and} \quad \int_{\tilde{B}_R} \rho(K_s * \varphi_{x_i}) dx \rightarrow \int_{\mathbb{R}^d} \rho(K_s * \varphi_{x_i}) dx$$

as $R \rightarrow +\infty$, which hold by dominated convergence due to the fact that $\rho(K_s * \varphi_{x_i}) \in L^1(\mathbb{R}^d)$ and $(K_s * \varphi) \rho_{x_i} \in L^1(\mathbb{R}^d)$. Hence, we may pass to the limit in (6.3) and get

$$\int_{\mathbb{R}^d} (K_s * \varphi) \rho_{x_i} dx = - \int_{\mathbb{R}^d} \rho(x) (K_s * \varphi_{x_i}) dx,$$

which can be combined with (6.1) and (6.2) to imply

$$\int_{\mathbb{R}^d} (K_s * \rho_{x_i}) \varphi dx = - \int_{\mathbb{R}^d} (K_s * \rho) \varphi_{x_i} dx.$$

Therefore, we have that $K_s * \rho \in W_{loc}^{1,1}(\mathbb{R}^d)$ and $\nabla(K_s * \rho) = K_s * \nabla \rho$ if $s < 1/2$. On the other hand, if $d \geq 2$ and $s > 1/2$ we even obtain $K_s * \rho \in W^{1,\infty}(\mathbb{R}^d)$, see [12, Lemma 1]. Else if $d \geq 2$ and $s = 1/2$ we obtain $K_s * \rho \in W^{1,p}$ for every $p \in (\frac{d}{d-1}, +\infty)$: indeed, since $\rho \in L^1(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$, by the Hardy-Littlewood-Sobolev inequality we get $K_s * \rho \in L^p(\mathbb{R}^d)$ for every $p \in (\frac{d}{d-1}, +\infty)$, thus $K_s * \rho$ belongs to the Bessel potential space $\mathcal{L}^{1,p}$ defined as $\mathcal{L}^{1,p} := L^p(\mathbb{R}^d) \cap \{K_{1/2} * g : g \in L^p(\mathbb{R}^d)\}$, which coincides with $W^{1,p}(\mathbb{R}^d)$, see [34, Theorem 3, pp 135].

The fact that $\rho^{m-1} \in W^{1,1}(\mathbb{R}^d)$ follows from the chain rule in Sobolev spaces, since $m > 2$ and $\rho \in L^\infty(\mathbb{R}^d) \cap W^{1,1}(\mathbb{R}^d)$. Therefore $\psi \in W_{loc}^{1,1}(\mathbb{R}^d)$.

In order to conclude, we write for every $\varphi \in C_c^\infty(\mathbb{R}^d)$

$$\begin{aligned} & \int_{\mathbb{R}^d} dx \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(x) \rho(y) dy \\ &= \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} dx \int_{|x-y|>\varepsilon} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(x) \rho(y) dy, \end{aligned}$$

then using the antisymmetry of the gradient of K_s we have

$$\begin{aligned} c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(x) \rho(y) dx dy \\ = \frac{2}{2s-d} \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_{s,\varepsilon} * \rho) dx \end{aligned}$$

where $K_{s,\varepsilon}$ is the following truncation of K_s :

$$K_{s,\varepsilon}(x) := \begin{cases} K_s(x), & \text{if } |x| > \varepsilon, \\ c_{d,s}\varepsilon^{2s-d}, & \text{if } |x| \leq \varepsilon. \end{cases}$$

At this point, we observe that

$$\begin{aligned} & \left| \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_{s,\varepsilon} * \rho) dx - \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_s * \rho) dx \right| \\ & \leq C \int_{\mathbb{R}^d} |\nabla \varphi(x) \cdot (K_{s,\varepsilon} - K_s) * \nabla \rho| dx \leq C \int_{\mathbb{R}^d} |(K_{s,\varepsilon} - K_s) * \nabla \rho| dx, \end{aligned}$$

thus by Young convolution inequality

$$\begin{aligned} & \left| \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_{s,\varepsilon} * \rho) dx - \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_s * \rho) dx \right| \\ & \leq C \|\nabla \rho\|_{L^1} \int_{\mathbb{R}^d} |K_{s,\varepsilon} - K_s| dx = C \|\nabla \rho\|_{L^1} \int_{|x| \leq \varepsilon} (|x|^{2s-d} - \varepsilon^{2s-d}) \rightarrow 0 \end{aligned}$$

as $\varepsilon \rightarrow 0$. Therefore we can write

$$\begin{aligned} & c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} \rho(x) \rho(y) dx dy \\ & = \frac{2}{2s-d} \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_s * \rho) dx. \end{aligned} \tag{6.4}$$

Since (2.1) implies

$$\chi \int_{\mathbb{R}^d} \rho(x) \nabla \varphi \cdot \nabla (K_s * \rho) dx = \int_{\mathbb{R}^d} \rho \nabla \left(\frac{m}{m-1} \rho^{m-1} + 2\beta \rho \right) \cdot \nabla \varphi dx = \int_{\mathbb{R}^d} \nabla (\rho^m + \beta \rho^2) \cdot \nabla \varphi dx,$$

by (6.4) the identity (2.2) follows. Vice versa, if ρ verifies (2.2), the same computation gives that ψ solves (2.1). \square

Proof of Lemma 4.13. Through the proof, for every $f \in L^2(\mathbb{R}^d)$ and every $\varphi \in C_c^\infty(\mathbb{R}^d)$ we shall use the notation

$$\mathcal{I}_s(f; \varphi) := (d - 2s) c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (\nabla \varphi(x) - \nabla \varphi(y)) \cdot (x - y) |x - y|^{2s-d-2} f(x) f(y) dx dy$$

The result is true if $\rho \in C_c^\infty(\mathbb{R}^d)$, since in this case we may apply (6.4), and we may integrate by parts and take advantage of the fact that $K_s \rightarrow \delta_0$ in the sense of distributions to get

$$\begin{aligned} \lim_{s \downarrow 0} \mathcal{I}_s(\rho; \varphi) &= -2 \lim_{s \downarrow 0} \int_{\mathbb{R}^d} \rho(x) \nabla \varphi(x) \cdot \nabla (K_s * \rho)(x) dx = -2 \int_{\mathbb{R}^d} \rho(x) \nabla \varphi(x) \cdot \nabla \rho(x) dx \\ &= - \int_{\mathbb{R}^d} \nabla \varphi(x) \cdot \nabla (\rho^2(x)) dx = \int_{\mathbb{R}^d} \Delta \varphi(x) \rho^2(x) dx. \end{aligned}$$

In order to obtain the result for $\rho \in L^1 \cap L^2(\mathbb{R}^d)$, let $(\rho_n)_{n \in \mathbb{N}} \subset C_c^\infty(\mathbb{R}^d)$ be a sequence that converges to ρ in $L^2(\mathbb{R}^d)$ and in $L^1(\mathbb{R}^d)$ as $n \rightarrow +\infty$. Thanks to (2.6) and by interpolation of L^p norms, we have

$$\begin{aligned} &|\mathcal{I}_s(\rho; \varphi) - \mathcal{I}_s(\rho_n; \varphi)| \\ &\leq (d-2s) c_{d,s} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |\nabla \varphi(x) - \nabla \varphi(y)| |x-y|^{2s-d-1} |\rho(x)\rho(y) - \rho_n(x)\rho_n(y)| dx dy \\ &\leq (d-2s) \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} |K_s(|x-y|) (\rho(x)\rho(y) - \rho_n(x)\rho_n(y))| dx dy \\ &\leq d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(|x-y|) |\rho(x)| |\rho(y) - \rho_n(y)| dx dy \\ &\quad + d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} K_s(|x-y|) |\rho(x) - \rho_n(x)| |\rho_n(y)| dx dy \\ &\leq 2d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| S_{d,s} \|\rho\|_{L^{\frac{2d}{d+2s}}(\mathbb{R}^d)} \|\rho - \rho_n\|_{L^{\frac{2d}{d+2s}}(\mathbb{R}^d)} \\ &\leq 2d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| S_{d,s} \|\rho\|_{L^1(\mathbb{R}^d)}^{\frac{4s}{d}} \|\rho - \rho_n\|_{L^1(\mathbb{R}^d)}^{\frac{4s}{d}} \|\rho\|_{L^2(\mathbb{R}^d)}^{\frac{d-2s}{d}} \|\rho - \rho_n\|_{L^2(\mathbb{R}^d)}^{\frac{d-2s}{d}}, \end{aligned} \tag{6.5}$$

so that for every $n \in \mathbb{N}$ and every $\varphi \in C_c^\infty(\mathbb{R}^d)$

$$\limsup_{s \downarrow 0} |\mathcal{I}_s(\rho; \varphi) - \mathcal{I}_s(\rho_n; \varphi)| \leq 2d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| \|\rho\|_{L^2(\mathbb{R}^d)} \|\rho - \rho_n\|_{L^2(\mathbb{R}^d)}. \tag{6.6}$$

Moreover, for every $n \in \mathbb{N}$ and every $\varphi \in C_c^\infty(\mathbb{R}^d)$, since $\rho_n \in C_c^\infty(\mathbb{R}^d)$, we have

$$\lim_{s \downarrow 0} \mathcal{I}_s(\rho_n; \varphi) = \int_{\mathbb{R}^d} \Delta \varphi(x) \rho_n^2(x) dx,$$

which together with (6.6) entails

$$\begin{aligned}
& \limsup_{s \downarrow 0} \left| \mathcal{I}_s(\rho; \varphi) - \int_{\mathbb{R}^d} \rho^2 \Delta \varphi \, dx \right| \\
& \leq \limsup_{s \downarrow 0} \left(|\mathcal{I}_s(\rho; \varphi) - \mathcal{I}_s(\rho_n; \varphi)| + \left| \mathcal{I}_s(\rho_n; \varphi) - \int_{\mathbb{R}^d} \rho_n^2 \Delta \varphi \right| + \int_{\mathbb{R}^d} |\rho^2 - \rho_n^2| \Delta \varphi \right) \\
& \leq 2d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| (\|\rho\|_{L^2(\mathbb{R}^d)} + \|\rho + \rho_n\|_{L^2(\mathbb{R}^d)}) \|\rho - \rho_n\|_{L^2(\mathbb{R}^d)}.
\end{aligned}$$

Since $\rho_n \rightarrow \rho$ in $L^2(\mathbb{R}^d)$, by taking the limit as $n \rightarrow +\infty$ we finally obtain $\mathcal{I}_s(\rho; \varphi) \rightarrow \int_{\mathbb{R}^d} \rho^2 \Delta \varphi$ as $s \downarrow 0$, for every $\varphi \in C_c^\infty(\mathbb{R}^d)$. \square

Proof of Lemma 4.14. With the same notation of the previous proof, we have for every $\varphi \in C_c^\infty(\mathbb{R}^d)$

$$\begin{aligned}
& \limsup_{s \downarrow 0} \left| \mathcal{I}_s(\rho_s; \varphi) - \int_{\mathbb{R}^d} \rho^2 \Delta \varphi \right| \\
& \leq \limsup_{s \downarrow 0} |\mathcal{I}_s(\rho_s; \varphi) - \mathcal{I}_s(\rho; \varphi)| + \limsup_{s \downarrow 0} \left| \mathcal{I}_s(\rho; \varphi) - \int_{\mathbb{R}^d} \rho^2 \Delta \varphi \right|,
\end{aligned}$$

therefore in view of Lemma 4.13, it will be enough to prove that $|\mathcal{I}_s(\rho_s; \varphi) - \mathcal{I}_s(\rho; \varphi)| \rightarrow 0$ as $s \downarrow 0$. But the very same estimates of (6.5) allow to obtain

$$\begin{aligned}
|\mathcal{I}_s(\rho_s; \varphi) - \mathcal{I}_s(\rho; \varphi)| & \leq 2d \sup_{x \in \mathbb{R}^d} |\nabla^2 \varphi(x)| S_{d,s} \|\rho\|_{L^1(\mathbb{R}^d)}^{\frac{4s}{d}} \|\rho - \rho_s\|_{L^1(\mathbb{R}^d)}^{\frac{4s}{d}} \|\rho\|_{L^2(\mathbb{R}^d)}^{\frac{d-2s}{d}} \\
& \quad \times \|\rho - \rho_s\|_{L^2(\mathbb{R}^d)}^{\frac{d-2s}{d}}
\end{aligned}$$

where the right hand side vanishes as $s \downarrow 0$ thanks to the assumptions on the family (ρ_s) . \square

Data availability

No data was used for the research described in the article.

Acknowledgments

The authors wish to thank Flavia Smarrazzo, Giuseppe Savaré and Yao Yao for fruitful discussions and suggestions.

E.M. acknowledges support from the MIUR-PRIN project No 2017TEXA3H and the MIUR-PRIN project No 202244A7YL. E.M. and B.V. are members of the GNAMPA group of the Istituto Nazionale di Alta Matematica (INdAM). The work of J. L. Vázquez

was funded by grant PGC2018-098440-B-I00 from the Spanish Government. He is an Honorary Professor at Univ. Complutense de Madrid.

References

- [1] L. Ambrosio, N. Gigli, G. Savaré, *Gradient Flows in Metric Spaces and in the Space of Probability Measures*, Lectures in Mathematics, Birkhäuser Verlag, Basel, 2008.
- [2] P. Billingsley, *Convergence of Probability Measures*, 2nd ed., Wiley & Sons, New York, 1999.
- [3] A. Blanchet, A gradient flow approach to the Keller-Segel systems, in: *RIMS Kokyuroku's Lecture Note*, vol. 1837, 2013, pp. 52–73.
- [4] A. Blanchet, V. Calvez, J.A. Carrillo, Convergence of the mass-transport steepest descent scheme for the sub-critical Patlak-Keller-Segel model, *SIAM J. Numer. Anal.* 46 (2) (2008) 691–721.
- [5] A. Blanchet, J.A. Carrillo, D. Kinderlehrer, M. Kowalczyk, P. Laurençot, S. Lisini, A hybrid variational principle for the Keller-Segel system in \mathbb{R}^2 , *ESAIM: Math. Model. Numer. Anal.* 49 (6) (2015) 1553–1576.
- [6] A. Blanchet, P. Laurençot, The parabolic-parabolic Keller-Segel system with critical diffusion as a gradient flow in R^d , $d \geq 3$, *Commun. Partial Differ. Equ.* 38 (2013) 658–686.
- [7] L. Caffarelli, J.L. Vázquez, Nonlinear porous medium flow with fractional potential pressure, *Arch. Ration. Mech. Anal.* 202 (2011) 537–565.
- [8] V. Calvez, J.A. Carrillo, F. Hoffmann, Equilibria of homogeneous functionals in the fair-competition regime, *Nonlinear Anal.* 159 (2017) 85–128.
- [9] V. Calvez, J.A. Carrillo, F. Hoffmann, The geometry of diffusing and self-attracting particles in a one-dimensional fair-competition regime, in: *Nonlocal and Nonlinear Diffusions and Interactions: New Methods and Directions*, in: *Lecture Notes in Math.*, vol. 2186, Springer, Cham, 2017, pp. 1–71.
- [10] V. Calvez, J.A. Carrillo, F. Hoffmann, Uniqueness of stationary states for singular Keller-Segel type models, *Nonlinear Anal.* 205 (2021) 112222.
- [11] J.A. Carrillo, S. Hittmeir, B. Volzone, Y. Yao, Nonlinear aggregation-diffusion equations: radial symmetry and long time asymptotics, *Invent. Math.* 218 (3) (2019) 889–977.
- [12] J.A. Carrillo, F. Hoffmann, E. Mainini, B. Volzone, Ground states in the diffusion-dominated regime, *Calc. Var. Partial Differ. Equ.* 57 (5) (2018) 127.
- [13] J.A. Carrillo, A. Chertock, Y. Huang, A finite-volume method for nonlinear nonlocal equations with a gradient flow structure, *Commun. Comput. Phys.* 17 (1) (2015) 233–258.
- [14] J.A. Carrillo, D. Castorina, B. Volzone, Ground states for diffusion dominated free energies with logarithmic interaction, *SIAM J. Math. Anal.* 47 (1) (2015) 1–25.
- [15] H. Chan, M.d.M. González, Y. Huang, E. Mainini, B. Volzone, Uniqueness of entire ground states for the fractional plasma problem, *Calc. Var. Partial Differ. Equ.* 59 (2020) 195.
- [16] G. De Figueiredo, E.M. Dos Santos, O.H. Miyagaki, Sobolev spaces of symmetric functions and applications, *J. Funct. Anal.* 261 (12) (2011) 3735–3770.
- [17] M.G. Delgadino, X. Yan, Y. Yao, Uniqueness and non-uniqueness of steady states of aggregation-diffusion equations, *Commun. Pure Appl. Math.* 75 (1) (2022) 3–59.
- [18] R. Jordan, D. Kinderlehrer, F. Otto, The variational formulation of the Fokker-Planck equation, *SIAM J. Math. Anal.* 29 (1998) 1–17.
- [19] I. Kim, Y. Yao, The Patlak-Keller-Segel model and its variations: properties of solutions via maximum principle, *SIAM J. Math. Anal.* 44 (2) (2012) 568–602.
- [20] E.H. Lieb, M. Loss, *Analysis*, second edition, Graduate Studies in Mathematics, vol. 14, American Mathematical Society, Providence, RI, 2001.
- [21] P.L. Lions, The concentration-compactness principle in the calculus of variations. The locally compact case, part 1, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* 1 (2) (1984) 109–145.
- [22] S. Lisini, E. Mainini, A. Segatti, A gradient flow approach to the porous medium equation with fractional pressure, *Arch. Ration. Mech. Anal.* 227 (2) (2018) 567–606.
- [23] D. Matthes, R.J. McCann, G. Savaré, A family of nonlinear fourth order equations of gradient flow type, *Commun. Partial Differ. Equ.* 34 (2009) 1352–1397.
- [24] R.J. McCann, A convexity principle for interacting gases, *Adv. Math.* 128 (1997) 153–179.
- [25] P. Mironescu, Superposition with subunitary powers in Sobolev spaces, *C. R. Math. Acad. Sci. Paris* 353 (6) (2015) 483–487.
- [26] M. Pierre, Uniqueness of the solutions of $u_t - \Delta\phi(u) = 0$ with initial datum a measure, *Nonlinear Anal., Theory Methods Appl.* 6 (1982) 175–187.

- [27] P.I. Plotnikov, Passing to the limit with respect to viscosity in an equation with variable parabolicity direction, *Differ. Equ.* 30 (4) (1994) 614–622.
- [28] P.I. Plotnikov, Forward-backward parabolic equations and hysteresis, *J. Math. Sci.* 93 (1999) 747–766.
- [29] X. Ros-Oton, J. Serra, Regularity theory for general stable operators, *J. Differ. Equ.* 260 (12) (2016) 8675–8715.
- [30] F. Santambrogio, Optimal transport for applied mathematicians. Calculus of variations, PDEs, and modeling, in: *Prog. Nonlinear Differ. Equ. Appl.*, Birkhäuser/Springer, 2015, 87.
- [31] T. Senba, T. Suzuki, Weak solutions to a parabolic-elliptic system of Chemotaxis, *J. Funct. Anal.* 191 (2002) 17–51.
- [32] J. Simon, Compact sets in the space $L^p(0, T; B)$, *Ann. Mat. Pura Appl.* (4) 146 (1987) 65–96.
- [33] F. Smarrazzo, A. Terracina, Sobolev approximation for two-phase solutions of forward-backward parabolic problems, *Discrete Contin. Dyn. Syst.* 33 (4) (2013) 1657–1697.
- [34] E.M. Stein, *Singular Integrals and Differentiability Properties of Functions*, Princeton Mathematical Series, vol. 30, Princeton University Press, Princeton, N.J., 1970.
- [35] J.L. Vázquez, *The Porous Medium Equation. Mathematical Theory*, Oxford University Press, Oxford, 2007.
- [36] C. Villani, *Topics in Optimal Transportation*, Graduate Studies in Mathematics, vol. 58, American Mathematical Society, Providence, RI, 2003.
- [37] Y.P. Zhang, On a class of diffusion-aggregation equations, *Discrete Contin. Dyn. Syst.* 40 (2) (2020) 907–932.