

Computer Programs in Physics

A SPIRED code for the reconstruction of spin distribution ☆,☆☆

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ABSTRACT

In Nuclear Magnetic Resonance (NMR), it is of crucial importance to have an accurate knowledge of the spin probability distribution corresponding to inhomogeneities of the magnetic fields. An accurate identification of the sample distribution requires a set of experimental data that is sufficiently rich to extract all fundamental information. These data depend strongly on the control fields (and their number) used experimentally to perturb the spin system. In this work, we present and analyze a greedy reconstruction algorithm, and provide the corresponding SPIRED code, for the computation of a set of control functions allowing the generation of data that are appropriate for the accurate reconstruction of a sample probability distribution. In particular, the focus is on NMR and spin dynamics governed by the Bloch system with inhomogeneities in both the static and radio-frequency magnetic fields applied to the sample. We show numerically that the algorithm is able to reconstruct non trivial joint probability distributions of the two inhomogeneous Hamiltonian parameters. A rigorous convergence analysis of the algorithm is also provided.

Program summary

Program title: SPIRED

CPC Library link to program files: <https://doi.org/10.17632/6fsmzpg6srg.1>

Licensing provisions: CC by 4.0

Programming language: MATLAB

Nature of problem: Identify the sample probability distribution corresponding to inhomogeneities of the magnetic field in Nuclear Magnetic Resonance from experimental data. The data depends strongly on the control fields and their number, and needs to be sufficiently rich in order to extract all fundamental information.

Solution method: Use greedy reconstruction algorithms to compute a set of control functions that allows the generation of data that are appropriate for an accurate reconstruction of the sample distribution.

Additional comments including restrictions and unusual features: Some routines in the SPIRED code use MATLAB's *fmincon*-solver, which requires MATLAB's Optimization Toolbox to be installed.

1. Introduction

Quantum Control (QC) is nowadays a well-recognized area of research [1–5] with many applications ranging from magnetic resonance [6–8] and atomic and molecular physics [9–12] to quantum technologies [7,13,14]. In this context, the control of spin dynamics by means of magnetic fields is today a well-established tool in Nuclear

Magnetic Resonance (NMR) for the development of new techniques [7]. The goal of quantum control is generally to design external control fields to perform quantum operations on the studied system. A severe limitation of QC comes from measurement processes which are much more difficult to account for than their classical counterpart. This explains that a majority of QC protocols are performed in an open-loop framework without any feedback from the experiment when applying

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the control (see however [15] for a closed-loop example). A good agreement between theory and experiment is achieved if all the parameters of the model system are perfectly known within a given range of precision. In the context of NMR, the latter may be due to the spatial extension of the sample which means that in practice all its elements are not subjected to the same external electromagnetic field. The values of such parameters can be estimated experimentally but can also be actively found by using specifically adapted controls. To this aim, different approaches using quantum features have been developed recently with success [16–18]. Among others, we can mention inversion techniques [19], selective controls [20–23], the maximization of quantum Fisher information [24–28] and the fingerprinting approach [29,30]. Such methods allow one to estimate the value of the Hamiltonian parameter as well as its variation range. However, the latter is not the only interesting quantity and the probability distribution is also a key feature of the experimental sample. When controlling an ensemble of quantum systems, this distribution can be interpreted as the number of individual systems having a given value of the parameter. The distribution can have a simple form such as a Gaussian or a Lorentzian one. In this case, the identification is quite straightforward and can be done using standard techniques. However, the identification is much more difficult when the distribution has a complex structure with, e.g., several peaks.

In a previous work [31], we introduced a Greedy Reconstruction Algorithm (GRA) to identify in a systematic way the probability distribution of one given Hamiltonian parameter. This was based on the framework presented in [32,19]. In particular, we focused on an ensemble of spin 1/2 particles in liquid state NMR subjected to an inhomogeneous radio-frequency magnetic field [33,6,34,8,35,36], where the algorithm was successfully applied to identify the distribution of the scaling factor corresponding to the magnetic field inhomogeneity. More precisely, we considered an inhomogeneous ensemble of uncoupled spin systems in liquid state whose dynamics are governed by the Bloch equation, but subjected to a different radio-frequency magnetic field. Applying a given number of constant control fields to the sample during a fixed time, the algorithm was able to determine the probability distribution of the inhomogeneous parameter from the measure of the final ensemble magnetization. We showed numerically in [31] that GRA is able to identify with a good accuracy nontrivial probability distributions with several peaks or with a step variation. An optimized version of this algorithm was also proposed to further improve the identification process and to limit the number of required controlled processes. Notice that a convergence analysis of the algorithm was only briefly sketched in [31], without rigorous proof.

The goal of the present paper is to extend the work [31] from three different points of view. Thus, the first novelty of this work is the extension of the GRA presented in [31], designed for the reconstruction a single parameter distribution, to the case of joint probability distributions of two distinct inhomogeneous Hamiltonian parameters. We consider numerically the case of both inhomogeneous static and radio-frequency magnetic fields applied to the sample. The second novelty is that we provide full MATLAB codes (SPIRED) implementing our GRA and its optimized version (denoted by OGRA) to find spin distributions. These codes can be directly used to solve the problems presented in [31] (in addition to the ones investigated in the present paper). The third novelty is that we prove theoretical convergence results for GRA. These cover also the ones only stated in [31]. As a result, this paper is a self-contained work that not only considers a more general problem than the one presented in [31], but also provides a full MATLAB code and a detailed and rigorous convergence analysis.

The paper is organized as follows. The identification problem of the spin distribution in NMR is presented in Sec. 2. The different variants of the greedy reconstruction algorithm are described in Sec. 3. Section 4 is dedicated to the description of the structure of the code SPIRED and its use. A convergence analysis of the algorithm is provided in Sec. 5. Numerical results are presented in Sec. 6. Conclusion

and prospective views are drawn in Sec. 7. Additional results are presented in Appendix A.

2. Identification of spin distribution

The framework of our SPIRED code is illustrated in a standard control problem in liquid-state NMR, i.e., a spin ensemble subjected to both inhomogeneous static and radio-frequency magnetic fields [33,37,8]. In a given rotating frame, each isochromat is characterized by a Bloch vector $\mathbf{M} = [M_x, M_y, M_z]^T$, evolving in time according to the equations

$$\begin{cases} \dot{M}_x = -\omega M_y + (1 + \alpha)\omega_y M_z, \\ \dot{M}_y = \omega M_x - (1 + \alpha)\omega_x M_z, \\ \dot{M}_z = (1 + \alpha)\omega_x M_y - (1 + \alpha)\omega_y M_x. \end{cases}$$

Notice that the components of $\mathbf{M}$ satisfy $M_x^2 + M_y^2 + M_z^2 = M_0^2$, with M_0 the equilibrium magnetization. Here, ω_x and ω_y are time-dependent controls corresponding to the components of the radio-frequency magnetic field along the x - and y - directions. The parameters ω and α correspond to offset, i.e., to the inhomogeneity of the static magnetic field applied along the z - direction, and radio-frequency magnetic field inhomogeneities, respectively [6]. In standard experiments, we have $\frac{\omega}{2\pi} \in [-20, 20]$ Hz and $\alpha \in [-0.2, 0.2]$. For the purpose of this paper, we assume that the probability densities of ω and α are unknown. The controls $\frac{\omega_x}{2\pi}$ and $\frac{\omega_y}{2\pi}$ are expressed in Hz. We consider a typical field amplitude ω_0 that can be fixed, for instance, to $\omega_0 = 2\pi \times 100$ Hz. We introduce normalized coordinates as

$$u_x = 2\pi \frac{\omega_x}{\omega_0}; \quad u_y = 2\pi \frac{\omega_y}{\omega_0}; \quad t' = \frac{\omega_0}{2\pi} t; \quad \Delta = 2\pi \frac{\omega}{\omega_0}; \quad \mathbf{X} = \frac{\mathbf{M}}{M_0}.$$

In what follows, we omit the prime to simplify the notations. We deduce that the differential system can be expressed in normalized units as

$$\begin{cases} \dot{x} = -\Delta y + (1 + \alpha)u_y z, \\ \dot{y} = \Delta x - (1 + \alpha)u_x z, \\ \dot{z} = (1 + \alpha)u_x y - (1 + \alpha)u_y x \end{cases} \quad (1)$$

with $x^2 + y^2 + z^2 = 1$. The initial state of the dynamics for each spin is the thermal equilibrium point, i.e., $\mathbf{X}_0 = [0, 0, 1]^T$. We consider a control time of the order of 100 ms, that corresponds to a normalized time t'_f of the order of 10. The range of variation of the parameter Δ is $\Delta_0 + 2\pi[-0.2, 0.2]$, where Δ_0 is a frequency value that can be used to shift arbitrarily the interval. For the purpose of this paper, we assume that $\Delta_0 \geq 0.4\pi$, meaning that $\Delta \geq 0$.

The goal of our SPIRED code is to estimate simultaneously the distributions for the parameters α and Δ by designing specific controls (u_x, u_y) . We consider an ensemble of N spins whose dynamics are governed by Eq. (1). We assume that the control amplitudes (u_x, u_y) belong to the admissible set $\mathcal{U} = \{(u_x, u_y) \in \mathbb{R}^2 \mid |u_x| \leq u_m, |u_y| \leq u_m\}$, where u_m is the maximum amplitude of each component. A simple way to proceed can be described as follows. We consider that the system of N spins is divided into K_Δ groups, and we associate with the ℓ -th subgroup a certain value Δ_ℓ and the corresponding probability $P_\star^\Delta(\ell) = \frac{N_{\Delta,\ell}}{N}$, $\ell = 1, \dots, K_\Delta$, with $\sum_{\ell=1}^{K_\Delta} P_\star^\Delta(\ell) = 1$. The probability $P_\star^\Delta(\ell)$ is unknown, which means that the number of elements $N_{\Delta,\ell}$ of each group is to be found. Similarly, for the parameter α , we have K_α groups with the probabilities $P_\star^\alpha(\ell) = \frac{N_{\alpha,\ell}}{N}$, $\ell = 1, \dots, K_\alpha$, with $\sum_{\ell=1}^{K_\alpha} P_\star^\alpha(\ell) = 1$ to estimate.

This problem can be viewed as a natural extension of the work [31] and leads to the identification of two independent discrete distributions. However, this approach has two main drawbacks. First, the two random variables Δ and α are assumed to be independent. This is a limitation when trying to reconstruct the two unknown distributions, since any possible correlation is a priori neglected. Second, the final identification problem is nonlinear, since the product of the two distributions would

appear. This is in contrast with the case of the reconstruction of one single distribution, where the identification problem is quadratic [31]. For these reasons, rather than considering two independent distributions, we work directly with the joint distribution, i.e., the system of N spins is divided into K groups and we associate to each subgroup a pair $(\alpha, \Delta)_\ell$ and the corresponding joint probability $P_\star(\ell) = \frac{N_\ell}{N}$, $\ell = 1, \dots, K$, with $\sum_{\ell=1}^K P_\star(\ell) = 1$. Now, the joint probability $P_\star(\ell)$ is unknown, namely the number of elements N_ℓ affected by the pairs $(\alpha, \Delta)_\ell$. This approach has the advantage of taking into account correlation effects and the final identification problem remains quadratic. It should be noted that these are acquired at the cost of an increase in the dimension of the unknown object(s), i.e., from two one-dimensional functions to a two-dimensional function. Finally, we point out that two independent distributions can also be treated as a specific case of joint distributions.

Since we are dealing with an inverse problem, we need to define what quantities can be observed in an experimental setting. In NMR, only the first two coordinates of the magnetization vector can be directly measured. We do not have accessed directly to the z component due to the strong constant magnetic field applied along this direction [6]. We denote by $\mathbf{Y}_{\mathbf{u},(\Delta,\alpha)}(t) = [x(t), y(t)]^\top$ the projection of the Bloch vector onto the first two coordinates. Here, the dependence on $\mathbf{u}$ and (Δ, α) has been explicitly mentioned. The corresponding experimental realization of this controlled dynamic is obtained at $t = t_f$ and leads to $\mathbf{Y}_{\mathbf{u}}^{\text{exp}}(t_f) = [x_{\mathbf{u}}^{\text{exp}}(t_f), y_{\mathbf{u}}^{\text{exp}}(t_f)]^\top$, where $\mathbf{Y}_{\mathbf{u}}^{\text{exp}}(t_f)$ can be viewed as the average at time t_f of the experimental measures of all the spins of the set subjected to the control $\mathbf{u}$. The coordinates $x_{\mathbf{u}}^{\text{exp}}$ and $y_{\mathbf{u}}^{\text{exp}}$ are those of this measured magnetization vector.

The relation between the theoretical description of the dynamical system to the experimental outcome can be expressed as

$$\mathbf{Y}_{\mathbf{u}}^{\text{exp}}(t_f) = \sum_{\ell=1}^K P_\star(\ell) \mathbf{Y}_{\mathbf{u},(\Delta,\alpha)_\ell}(t_f), \quad (2)$$

in which the two sides of the equation crucially depend on the control $\mathbf{u}$.

In general, one control protocol is not sufficient to obtain an appropriate identification of the unknown $P_\star$, but a set of $\tilde{K}$ control processes with $\tilde{K}$ different control functions denoted $\mathbf{u}_k$, $k = 1, \dots, \tilde{K}$, needs to be used. On the basis of the experimental outputs, a straightforward way to determine $P_\star$ is to solve the minimization problem

$$\min_{P \in \mathbb{P}} \sum_{k=1}^K \|\mathbf{Y}_{\mathbf{u}_k}^{\text{exp}}(t_f) - \sum_{\ell=1}^K P(\ell) \mathbf{Y}_{\mathbf{u}_k,(\Delta,\alpha)_\ell}(t_f)\|^2, \quad (3)$$

where $\|\cdot\|$ denotes the standard Euclidean vector norm, and $\mathbb{P}$ is the convex and closed set of all probability distributions P that satisfy $P(\ell) \geq 0$ for $1 \leq \ell \leq K$ and $\sum_{\ell=1}^K P(\ell) = 1$. At this point, it is clear that a key ingredient of the accuracy of the identification process rests on the choice of a set of $\tilde{K}$ controls $\mathbf{u}_k$. The identification of the number $\tilde{K}$ of control functions is a difficult task. The theoretical analysis presented in Sec. 5 shows that the choice $\tilde{K} = K$ is sufficient. The GRA algorithm computes exactly $\tilde{K} = K$ control fields. However, we will show that OGRA is capable of reducing (halving) the number $\tilde{K}$ of control fields while guaranteeing an accurate identification.

Let us now rewrite (3) in a form that we consider in our SPIRED implementation. We introduce a set $\Phi := \{\phi_j\}_{j=1}^K$ of linearly independent functions $\phi_j : \{1, \dots, K\} \rightarrow \mathbb{R}$ such that $\mathbb{P} \subset \text{span}(\Phi)$, where span denotes the vector space generated by the functions. Expressing P as $P(\ell) = \sum_{j=1}^K \beta_j \phi_j(\ell)$, the minimization problem (3) becomes

$$\min_{\beta \in \mathbb{R}^K} \sum_{k=1}^K \|\mathbf{Y}_{\mathbf{u}_k}^{\text{exp}}(t_f) - \sum_{\ell=1}^K \beta_j \phi_j(\ell) \mathbf{Y}_{\mathbf{u}_k,(\Delta,\alpha)_\ell}(t_f)\|^2, \quad (4)$$

where the vector $\beta = (\beta_j)_{j=1}^K$ is taken in $\hat{\mathbb{R}}^K$, a subset of $\mathbb{R}^K$, so that $P = \sum_j \beta_j \phi_j$ is a probability distribution.

We show in this study that GRA allows us to design a set of controls $\mathbf{u}_k$ that makes (4) solvable and well conditioned. The algorithm is composed of two steps, namely an offline and an online steps. In the offline step, GRA computes the controls $\mathbf{u}_k$. In this step, only the theoretical model is needed without any experimental input. The derived controls are used in the online step in which the different magnetization vectors are measured and the minimization problem (3) is solved. Note that the controls are the same for any probability distribution to identify and only depend on the model system under study. Finally, we point out that in our algorithms the duration of each control pulse is considered as a variable to be optimized together with its amplitude. In particular, we assume that the controls are constant in time, i.e., $\mathbf{u}(t) \equiv \mathbf{u} \in \mathbb{R}^2$, and that we can freely choose the control time up to a fixed maximum value t_f . Since the initial state is an equilibrium point, this is equivalent to turning on the control at a time $t \geq 0$. We show in Sec. 5 that these hypotheses are sufficient for the different examples to identify the probability distributions. The generality of GRA allows one to tackle this situation in a straightforward manner.

3. Greedy reconstruction algorithms

We present in this section the GRA in its classical form and in an optimized extension called optimized GRA (OGRA).

GRA computes the controls $\mathbf{u}_k$ and the corresponding control times t_k by solving a sequence of fitting-step and discriminatory-step problems. The goal of the fitting step is to identify a defect of the system, namely a nontrivial kernel of a certain matrix W introduced below, while the discriminatory step designs a new control which is aimed to correct this discrepancy and to eliminate the identified nontrivial kernel. The explicit formulation of GRA is presented in Algorithm 1 and is given in terms of the function $\mathbf{h}^{(k)}$ defined as

$$\mathbf{h}^{(k)}(\beta, \mathbf{u}, t) := \sum_{\ell=1}^K \sum_{j=1}^k \beta_j \phi_j(\ell) \mathbf{Y}_{\mathbf{u},(\Delta,\alpha)_\ell}(t), \quad (5)$$

for any β in $\mathbb{R}^k$. Notice that the fitting step minimizes over the full space $\mathbb{R}^k$, meaning that $\sum_j \beta_j \phi_j$ does not have to be a probability dis-

Greedy Reconstruction Algorithm (GRA)

Require: A set of K linearly independent functions $\Phi = \{\phi_1, \dots, \phi_K\}$.

1: Compute the control $\mathbf{u}_1$ and the control time t_1 by solving

$$\max_{\substack{\mathbf{u} \in \mathcal{U} \\ t \in [0, t_f]}} \|\mathbf{h}^{(1)}(1, \mathbf{u}, t)\|^2. \quad (6)$$

2: **for** $k = 1, \dots, K-1$ **do**

3: Fitting step: Find $\beta^k = (\beta_j^k)_{j=1, \dots, k}$ that solves

$$\min_{\beta \in \mathbb{R}^k} \sum_{m=1}^k \|\mathbf{h}^{(K)}(\mathbf{e}_{k+1}, \mathbf{u}_m, t_m) - \mathbf{h}^{(k)}(\beta, \mathbf{u}_m, t_m)\|^2, \quad (7)$$

where $\mathbf{e}_{k+1}$ is the $(k+1)$ -th canonical vector in $\mathbb{R}^K$.

4: Discriminatory step: Find $\mathbf{u}_{k+1}$ and t_{k+1} that solves

$$\max_{\substack{\mathbf{u} \in \mathcal{U} \\ t \in [0, t_f]}} \|\mathbf{h}^{(K)}(\mathbf{e}_{k+1}, \mathbf{u}, t) - \mathbf{h}^{(k)}(\beta^k, \mathbf{u}, t)\|^2. \quad (8)$$

5: **end for**

Algorithm 1: GRA consists of three steps. First, the initialization problem (6) computes the control $\mathbf{u}_1$ and control time t_1 . Afterwards, the algorithm iterates between the fitting and discriminatory step. The fitting step attempts to find two different probability distributions such that their final ensemble magnetization can not be distinguished by any previously computed control. Then, the discriminatory step finds a new control $\mathbf{u}_{k+1}$ and control time t_{k+1} that maximize the difference between these final ensemble magnetizations. A graphical representation of the workflow of GRA can be found in Fig. 1.

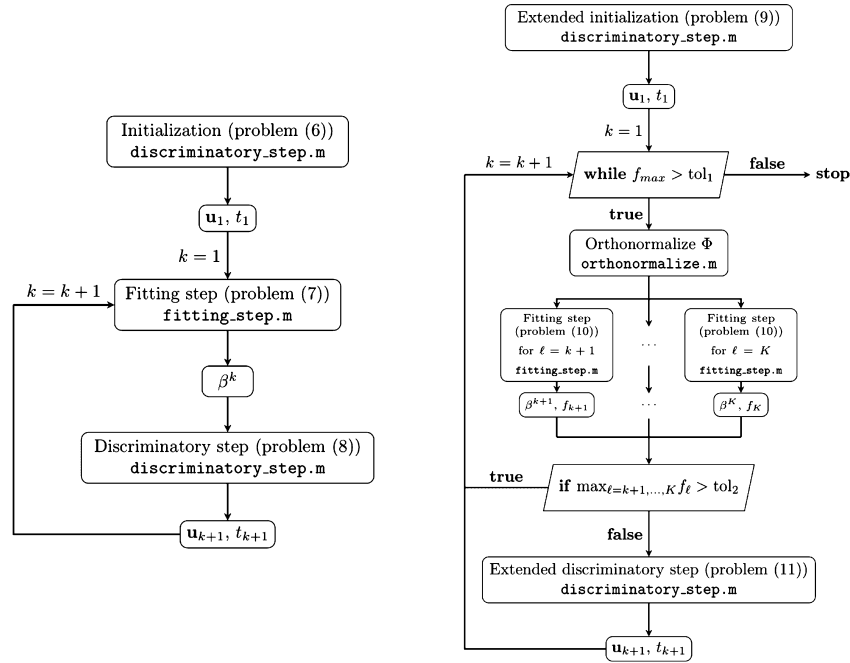


Fig. 1. Graphical representation of the workflow of GRA (left) and OGRA (right). The main difference between the two strategies are the **while**-condition that allows the algorithm to stop when it can not separate any more states by a new control, and the **if**-condition that allows it to skip the discriminatory step if one of the previously computed controls already separates the two states computed by one of the fitting step problems.

tribution. However, this is a restrictive condition. On the contrary, it allows the algorithm to find and correct more nontrivial kernels than might be necessary.

One main characteristic of GRA is that the set Φ and its order have to be fixed a-priori. However, the choice and order of Φ can have a crucial impact on the outcome of the algorithm as shown in [32, Sec. 5.3]. Hence, the idea of OGRA, which is stated in Algorithm 2, is to make the algorithm independent of the choice and order of the set Φ . Additionally, it aims at avoiding the computation of unnecessary control functions. This is achieved by two adaptations in GRA. The first one is that in each step one does not only consider the next element (the next canonical vector $\mathbf{e}_{k+1}$) in the set, but all remaining elements (the canonical vectors $\mathbf{e}_{k+\ell}$ for all $1 \leq \ell \leq K - k$) in parallel. Hence, it is also possible to enlarge the set Φ (and thus enlarge K) by additional functions ϕ_k which do not have to be linearly independent. In order to progressively remove linearly dependent functions in the set and to avoid scaling issues, all remaining basis elements are orthonormalized against the already selected ones after each iteration. The second adaptation is the introduction of two tolerances $\text{tol}_1, \text{tol}_2 > 0$. The first tolerance tol_1 is used as a stopping criterion. The algorithm terminates if the function value in the initialization or any of the discriminatory steps (denoted by f_ℓ in Algorithm 2) is too small, thus not adding new information. The second tolerance tol_2 is used to skip the computation of a new control field in the discriminatory step, if the minimum cost function value computed by the fitting step is not small enough. If this function value is large, then there already exists a control function in the current set of controls that discriminates between the two distributions $\phi_{k+\ell}$ and $\sum_{j=1}^{\ell} \beta_j^\ell \phi_j$. Notice that setting tol_2 to a very small value is reasonable if the final identification problem is quadratic. In this case, one can prove that a nonzero cost function value in the fitting step implies that one does not need to compute a new control for the corresponding set element (compare with [32]). However, if the final identification problem is not quadratic, then it can make sense to set tol_2 to a larger value.

Optimized Greedy Reconstruction Algorithm (OGRA)

Require: A set of K functions $\Phi = \{\phi_1, \dots, \phi_K\}$ and two tolerances $\text{tol}_1, \text{tol}_2 > 0$.

- 1: Compute $\mathbf{u}_1, t_1$ and the index ℓ_1 by solving the initialization problem

$$\max_{\ell \in \{1, \dots, K\}} \max_{\mathbf{u} \in U} \min_{t \in [0, t_f]} \|\mathbf{h}^{(1)}(\mathbf{e}_\ell, \mathbf{u}, t)\|^2, \quad (9)$$

- 2: Swap ϕ_1 and ϕ_{ℓ_1} in Φ , and set $k = 1, \tilde{k} = 1$, and $f_{max} = \|\mathbf{h}^{(1)}(\mathbf{e}_{\ell_1}, \mathbf{u}_1, t_1)\|^2$.
- 3: **while** $k \leq K - 1$ and $f_{max} > \text{tol}_1$ **do**
- 4: Orthonormalize all elements $(\phi_{k+1}, \dots, \phi_K)$ with respect to $(\phi_1, \dots, \phi_k)$, remove any that are linearly dependent and update K accordingly.
- 5: **for** $\ell = k + 1, \dots, K$ **do**
- 6: Fitting step: Find $(\beta_j^\ell)_{j=1, \dots, k}$ that solve the problem

$$\min_{\beta \in \mathbb{R}^k} \sum_{m=1}^{\tilde{k}} \|\mathbf{h}^{(K)}(\mathbf{e}_{k+\ell}, \mathbf{u}_m, t_m) - \mathbf{h}^{(k)}(\beta, \mathbf{u}_m, t_m)\|^2, \quad (10)$$

$$\text{and set } f_\ell = \sum_{m=1}^{\tilde{k}} \|\mathbf{h}^{(K)}(\mathbf{e}_{k+\ell}, \mathbf{u}_m, t_m) - \mathbf{h}^{(k)}(\beta^\ell, \mathbf{u}_m, t_m)\|^2.$$

- 7: **end for**
- 8: **if** $\max_{\ell=k+1, \dots, K} f_\ell > \text{tol}_2$ **then**
- 9: Set $\ell_{k+1} = \arg \max_{\ell=k+1, \dots, K} f_\ell$.
- 10: **else**
- 11: Extended discriminatory step: Find $\mathbf{u}_{\tilde{k}+1}, t_{\tilde{k}+1}$ and ℓ_{k+1} that solve

$$\max_{\ell \in \{k+1, \dots, K\}} \max_{\mathbf{u} \in U} \min_{t \in [0, t_f]} \|\mathbf{h}^{(K)}(\mathbf{e}_{k+\ell}, \mathbf{u}, t) - \mathbf{h}^{(k)}(\beta^\ell, \mathbf{u}, t)\|^2. \quad (11)$$
- 12: Set $\tilde{k} = \tilde{k} + 1$.
- 13: **end if**
- 14: Swap ϕ_{k+1} and $\phi_{\ell_{k+1}}$ in Φ .
- 15: Set $f_{max} = \|\mathbf{h}^{(K)}(\mathbf{e}_{k+\ell_{k+1}}, \mathbf{u}_{\tilde{k}+1}, t_{\tilde{k}+1}) - \mathbf{h}^{(k)}(\beta^{\ell_{k+1}}, \mathbf{u}_{\tilde{k}+1}, t_{\tilde{k}+1})\|^2$.
- 16: Set $k = k + 1$.
- 17: **end while**

Algorithm 2: OGRA consists of the same three main substeps as GRA, with the difference that OGRA considers all remaining basis elements simultaneously in each substep. Additionally, OGRA is able to skip discriminatory steps if no new control is required (line 8-9), or stop if no control that adds new information is found at all (line 3). The orthonormalization of the set Φ in line 4 removes redundant elements and prevents numerical instabilities. A graphical representation of the workflow of OGRA can be found in Fig. 1.

Table 1

Routines related to the greedy reconstruction algorithm.

GRA routines	Description
main	Main function used to run the code.
GRA	Greedy reconstruction algorithm.
OGRA	Optimized greedy reconstruction algorithm.
discriminatory_step	Routine that solves the initialization and discriminatory step problem using MATLAB's <i>fmincon</i> -solver.
fitting_step	Routine that solves the fitting step problem.
orthonormalize	Routine that orthonormalizes all remaining basis elements after each iteration of OGRA.
SVD_solver	Routine that solves the fitting step problem using the singular value decomposition (SVD).

Table 2

Routines related to the reconstruction of the probability distribution.

Reconstruction routines	Description
generate_data	Routine that generates the experimental realizations for all computed controls.
reconstr	Routine that either solves the final identification problem (4) using a second order interior-point algorithm, or the compact form (12) in Section 5) using a solver based on the SVD.

In conclusion, the main adaptations of OGRA in lines 3 and 8-9 allow the algorithm to reduce the number of computed controls $\tilde{K}$, meaning that $\tilde{K} < K$. On the other hand, as we mentioned before, GRA is designed to always compute exactly $\tilde{K} = K$ controls.

Flowcharts describing GRA and OGRA can be found in Fig. 1. The numerical implementation of GRA and OGRA is presented and discussed in the following sections.

4. The SPIRED code

4.1. Structure of the code

In this section, we provide a full list of all MATLAB functions contained in the SPIRED code. Inside the SPIRED folder the user can find the main routine used to run the SPIRED code, as well as the routines that run GRA and OGRA, and that solve their sub-problems (see Table 1). Additionally, the SPIRED folder contains the routines that generate the synthetic experimental data for the true parameter probability distribution, and that solve the final identification problem (3) (see Table 2). Notice that the both the *discriminatory_step* and the *reconstr* routine use MATLAB's *fmincon*-solver, which requires MATLAB's Optimization Toolbox to be installed.

There are also three subfolders labeled “Test1”, “Test2” and “Test3”. These contain three test problems the user can choose from. “Test1” corresponds to the problem discussed in this paper. “Test2” is the same as “Test1”, but only considers a control in the x direction (in other words $u_y = 0$ for all control fields). Finally, “Test3” corresponds to the problem investigated in [31], where the resonance offset Δ is fixed and one attempts to reconstruct only the control inhomogeneity parameter α . Each of these “Test” folders contains routines to set the input variables, describe the cost function and its gradient for the discriminatory-step problem, and solve the corresponding dynamical system (see Table 3). They also each contain two routines used to plot the results and condition number of the reconstruction process (see Table 4).

4.2. Usage of the code

Here we illustrate the working procedure of the SPIRED code with an example. The user needs to define the test problem in the function *main*, which is used to initialize the procedure.

```
function [ controls, bases, model, results ] = main
```

Table 3

Routines related to the test problems.

Test routines	Description
starting	Input function.
fun_discriminatory	Function computing the cost functional and gradient for the discriminatory step problem.
NMR_solver	Routine that solves the (normalized) dynamical system via direct calculations of the exponential matrix (compare the proof of Theorem 2).

Table 4

Routines plotting the results for the test problem.

Plotting routines	Description
plot_reconstr	Routine that plots the true and reconstructed probability distributions for the two control sets.
plot_condition	Routine that produces a table containing the condition numbers corresponding to the reconstruction process.

```
% STEP 1: Choice of the Problem;
addpath('Test1');
% STEP 2: Assemble problem variables;
[ model, bases, options ] = starting ( );
% STEP 3: Run!
[ controls.GRA, results.GRA ] = GRA( bases.GRA, model,
    options );
[ controls.OGRA, bases.OGRA, results.OGRA ] = OGRA( bases.
    OGRA, model, options );
% STEP 4: Compute (synthetic) experimental data;
Y_exp.GRA = generate_data( controls.GRA, model );
Y_exp.OGRA = generate_data( controls.OGRA, model );
% STEP 5: Solve the final identification problem;
...
```

In particular, at the “STEP 1” the user needs to define the path of the folder containing the test routines. Then, the user can define the input variables in the function *starting*, which is listed exemplary for the first test problem in the following.

```
function [ model, bases, options ] = starting ( )
% STEP 1: Input variables;
% control bounds and maximum control time;
um = 10;
tf = 16;
% variables for the unknown parameters
Delta0 = 4*pi;
Delta1 = 0.2;
Delta_interval = Delta0 + 2*pi.*[-Delta1, Delta1];
alpha_interval = [-0.2, 0.2];
% number of grid points for the unknown parameters
nr_alphas = 10;
nr_Deltas = 10;
% open the file to get the input probability distribution
load('Test1/Distributions/Gaussian.mat', 'P_star')
% number of spins;
nr_spins = 100000;
% options for GRA and OGRA
iterations = nr_alphas*nr_Deltas;
Display_GRA = 'off';
flag_orth = 1;
% numerical parameters for OGRA;
tol_OGRA_fit = 1e-4;
tol_OGRA_discr = 1e-14;
% tolerance for the SVD solver in the fitting step (and
    optionally for the reconstruction solver)
tol_svd = 1e-10;
% optimization method for the final identification problem
solver = 'fmincon';
```

At “STEP 1” in this function, the user can define the input variables and the path to the .mat file containing the true probability distribution P_* . The input parameters related to the problem are

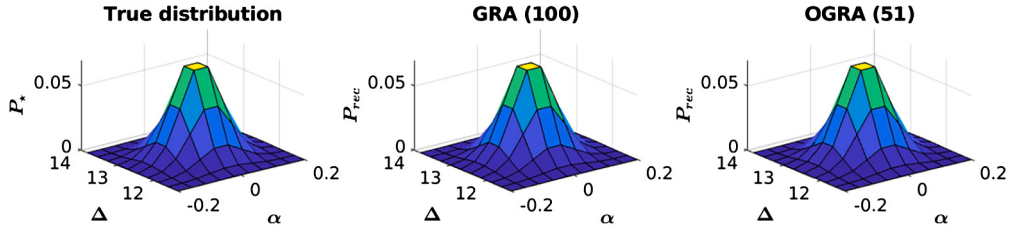


Fig. 2. The plot on the left shows the true Gaussian probability distribution for $K = 100$ uniform mesh points. The plots in the middle and on the right contain the reconstructed distributions for the control sets generated by GRA (containing 100 control fields) and OGRA (containing 51 control fields).

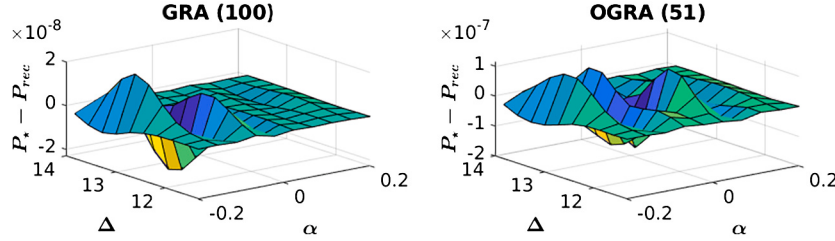


Fig. 3. Difference between the true probability distribution $P_\star$ and the distributions P_{rec} reconstructed using the control sets generated by GRA (left) and OGRA (right). In brackets are the number of control fields for each set.

- `um`: bound u_m for the absolute value of the control amplitudes;
- `tf`: maximum normalized control time t_f ;
- `Delta0`: frequency shift Δ_0 for the normalized resonance offset interval;
- `Delta1`: width Δ_1 of the normalized resonance offset interval;
- `Delta_interval`: interval boundaries for the normalized resonance offset Δ ;
- `alpha_interval`: interval boundaries for the control field inhomogeneity parameter α ;
- `nr_alphas`: number of grid points in the direction of α for the joint discrete parameter probability distribution of α and Δ ;
- `nr_Deltas`: number of grid points in the direction of Δ for the joint discrete parameter probability distribution of α and Δ ;
- `nr_spins`: number of spins in the system;
- `iterations`: (maximum) number of iterations performed by GRA and OGRA; for any full basis of the discrete parameter space, the obvious choice is the total number of discretization points, which is the product of `nr_alphas` and `nr_Deltas`;
- `Display_GRA` Display option to print information about the current iteration of GRA and OGRA in the command window; can be set to 'off' to display no output, 'iter' to show the current substep of GRA and OGRA, or 'iter-detailed' to also show the current optimization problem during the substeps of OGRA;
- `flag_orth`: flag variable that turns the orthonormalization of the remaining basis elements during OGRA on or off;
- `tol_OGRA_fit`: tolerance tol_2 for OGRA;
- `tol_OGRA_discr`: tolerance tol_1 for OGRA;
- `tol_svd`: tolerance for the SVD solver, used in the fitting step and (optionally) for the final identification problem;
- `solver`: optimization method used to solve the final identification problem (4); can be set to 'fmincon' to solve (4) using the second-order interior-point algorithm of MATLAB's `fmincon`-solver, or to 'svd' to solve a compact form of the problem (compare (22) in Section 6) using the SVD solver;

The .mat file has to contain the variable `P_star`, which is the vectorized true joint probability distribution $P_\star$. If the user is considering a true experimental (laboratory) setup, meaning that they perform real experiments for the different controls to obtain the experimental data and that the true probability distribution is truly unknown, they should replace "STEP 4" in the "main.m" file with a load command to fetch the real experimental data.

Finally, to run the code the user has to write on the MATLAB prompt the following

```
>> [ controls , bases , model , results ] = main
```

After the computations, the routine saves the results in the MATLAB workspace (as documented in the code) and plots the reconstructed probability distributions and their difference to the true one, as well as the condition numbers for different mesh sizes.

In particular, the results obtained by running "Test1" are the true and reconstructed probability distributions for the control fields generated by GRA and OGRA, shown in Fig. 2. In Fig. 3 we show the difference between the true and reconstructed distributions for the two control field sets. Additionally, the routine provides a table containing the exact condition numbers corresponding to GRA and OGRA. Since the solver for the discriminatory step problem is initialized with a random vector, there may be small variations in some results, without changing the overall outcome.

Examples of all figures produced by the different test problems are also provided in the folder "Results" that can be found in the corresponding "Test"-folder. There one can also find the .mat files containing the set of random controls for each test problem, loaded in the `main`.

5. Convergence analysis

In this section, we prove that the controls generated by GRA and OGRA make possible the identification of the unknown probability distributions of the parameters Δ and α , i.e., they make problem (3) uniquely solvable.

We start by recalling that problem (3) is equivalent to (4). Assuming that $P_\star$ can be written as $P_\star(\ell) = \sum_{j=1}^K \beta_{\star,j} \phi_j(\ell)$, we can write equation (4) in a compact form as follows:

$$\min_{\beta \in \mathbb{R}^K} \langle \beta_\star - \beta | W | \beta_\star - \beta \rangle, \quad (12)$$

where $W := \sum_k W(\mathbf{u}_k, t_k)$ is the sum of symmetric and positive semi-definite $K \times K$ -matrices whose elements are defined as:

$$[W(\mathbf{u}_k, t_k)]_{\ell,j} := \langle \gamma_\ell(\mathbf{u}_k, t_k) | \gamma_j(\mathbf{u}_k, t_k) \rangle \quad (13)$$

with

$$\gamma_j(\mathbf{u}_k, t_k) := \sum_{\ell} \phi_j(\ell) \mathbf{Y}_{\mathbf{u}_k, (\alpha_\ell, \Delta_\ell)}(t_k). \quad (14)$$

Since the set of vectors β is a convex subset of $\mathbb{R}^K$, we deduce that the problem is uniquely solvable if the matrix W is positive definite. In the case W has a non-trivial kernel, infinitely many solutions may exist which lead to wrong probability distributions different from the experimental one P_* . We stress that the non-triviality of the kernel depends completely on the choice of the controls $\mathbf{u}_k$ and the corresponding control times t_k .

Using the notation (13)-(14), we can now also rewrite the subproblems of GRA in terms of the matrix W . The initialization problem (6) can be written as

$$\max_{\substack{\mathbf{u} \in \mathcal{U} \\ t \in [0, t_f]}} \| [W(\mathbf{u}, t)]_{1,1} \|^2. \quad (15)$$

The fitting step problem (7) is equivalent to

$$\min_{\beta \in \mathbb{R}^K} \langle \mathbf{v}_\beta | [W^k]_{[1:k+1, 1:k+1]} | \mathbf{v}_\beta \rangle, \quad (16)$$

where $W^k := \sum_{m=1}^k W(\mathbf{u}_m, t_m)$ and $\mathbf{v}_\beta := [\beta^\top, -1]^\top$. Finally, the discriminatory step problem (8) can be written as

$$\max_{\substack{\mathbf{u} \in \mathcal{U} \\ t \in [0, t_f]}} \langle \mathbf{v}_\beta | [W(\mathbf{u}, t)]_{[1:k+1, 1:k+1]} | \mathbf{v}_\beta \rangle. \quad (17)$$

A direct interpretation of these reformulated problems is that each control $\mathbf{u}_k$ generated by GRA at iteration k ensures that $\langle \mathbf{e}_k | W | \mathbf{e}_k \rangle > 0$. Iteratively, this implies that $\langle \mathbf{w} | W | \mathbf{w} \rangle > 0$ for any $\mathbf{w} \in \mathbb{R}^K$ which is equivalent to W being positive definite.

In more details, the first control $\mathbf{u}_1$ and the control time t_1 are chosen by the initialization (15) such that the first upper left entry of $W(\mathbf{u}_1, t_1)$ is positive. This guarantees that $\langle \mathbf{e}_1 | W | \mathbf{e}_1 \rangle > 0$ since

$$\langle \mathbf{e}_1 | W | \mathbf{e}_1 \rangle = \sum_{j=1}^K \langle \mathbf{e}_1 | W(\mathbf{u}_j, t_j) | \mathbf{e}_1 \rangle \geq [W(\mathbf{u}_1, t_1)]_{1,1} > 0,$$

where we used that $W(\mathbf{u}, t)$ is positive semi-definite for any $\mathbf{u}$ and t . Assume now that the upper left 2×2 -submatrix of $W^1 = W(\mathbf{u}_1, t_1)$ is not positive definite. Then it has a one-dimensional kernel spanned by a vector $\mathbf{v}_{\beta^1} := [\beta^1, -1]^\top \in \mathbb{R}^2$ (see [32, Lem. 5.3]). The corresponding scalar β^1 is clearly the unique solution to the fitting step problem (16) for $k = 1$. Now, the discriminatory step problem (17) attempts to find a control $\mathbf{u}_2 \in \mathcal{U}$ and a control time $t_2 \in [0, t_f]$ such that the vector $\mathbf{v}_{\beta^1}$ is not in the kernel of the upper left 2×2 -submatrix of $W(\mathbf{u}_2, t_2)$. If this is successful then the upper left 2×2 -submatrix of $W^2 = W(\mathbf{u}_1, t_1) + W(\mathbf{u}_2, t_2)$ is positive definite. This also implies that $\langle \mathbf{e}_2 | W | \mathbf{e}_2 \rangle > 0$. Repeating this procedure for $k = 2, \dots, K-1$, we obtain $\langle \mathbf{e}_k | W | \mathbf{e}_k \rangle > 0$ for all $k \in \{1, \dots, K\}$, which guarantees that W is positive definite. We summarize the arguments above in the following theorem.

Theorem 1. Let $\{(\mathbf{u}_k, t_k)\}_{k=1}^K$ be a set of controls and corresponding control times generated by GRA, such that $[W(\mathbf{u}_1, t_1)]_{1,1} > 0$. Let β^k be the solution to the fitting step problem (16) for $k = 1, \dots, K-1$, such that the vectors $\mathbf{v}_{\beta^k} = [(\beta^k)^\top, -1]^\top$ are not in the kernel of $[W(\mathbf{u}_{k+1}, t_{k+1})]_{[1:k+1, 1:k+1]}$. Then the matrix $W = \sum_k W(\mathbf{u}_k, t_k)$ is positive definite.

It remains to show that the discriminatory step can always find a control such that the vector $\mathbf{v}_{\beta^k}$ is not in the kernel of $[W(\mathbf{u}, t)]_{[1:k+1, 1:k+1]}$. In fact, it is sufficient to show that for any $k \in \{1, \dots, K\}$ there exists a control $\mathbf{u} \in \mathcal{U}$ and a $t \in [0, t_f]$ such that $\langle \mathbf{v}_{\beta^k} | [W(\mathbf{u}, t)]_{[1:k+1, 1:k+1]} | \mathbf{v}_{\beta^k} \rangle > 0$. We show in Theorem 2 that this is valid in the context of this paper.

Theorem 2. Let $k \in \{1, \dots, K-1\}$, $W^k_{[1:k, 1:k]}$ be positive definite, β^k the solution to the fitting-step problem (16), and $\mathbf{v}_{\beta^k} = [(\beta^k)^\top, -1]^\top$. Then any solution $(\mathbf{u}, t)$ to the discriminatory-step problem (17) satisfies

$$\langle \mathbf{v}_{\beta^k} | W_{[1:k+1, 1:k+1]}(\mathbf{u}, t) | \mathbf{v}_{\beta^k} \rangle = \| \mathbf{h}^{(k)}(\mathbf{e}^{k+1}, \mathbf{u}; t) - \mathbf{h}^{(k)}(\beta^k, \mathbf{u}; t) \|^2 > 0.$$

Proof. For brevity, we identify α with $1 + \alpha$ for the remainder of this proof. We start by writing

$$\mathbf{h}^{(K)}(\mathbf{e}^{k+1}, \mathbf{u}, t) - \mathbf{h}^{(k)}(\beta^k, \mathbf{u}, t) = \sum_{\ell=1}^K \left(\phi_{k+1}(\ell) - \sum_{j=1}^k \beta_j^k \phi_j(\ell) \right) \mathbf{Y}_{\mathbf{u}, (\Delta, \alpha)_\ell}(t).$$

Since the functions $\{\phi_1, \dots, \phi_K\}$ are linearly independent, it holds that

$$\exists \tilde{\ell} \in \{1, \dots, K\} : \phi_{diff}(\tilde{\ell}) := \phi_{k+1}(\tilde{\ell}) - \sum_{j=1}^k \beta_j^k \phi_j(\tilde{\ell}) \neq 0. \quad (18)$$

According to (1), we have for any $(\alpha, \Delta)_\ell$

$$\frac{d}{dt} \mathbf{X}(t) = [\Delta_\ell A + \alpha_\ell (u_x B_x + u_y B_y)] \mathbf{X}(t), \quad \mathbf{X}(0) = \mathbf{X}_0, \quad (19)$$

where

$$A = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \quad B_x = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix}, \quad B_y = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix}, \quad \mathbf{X}_0 = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}.$$

Now, consider the control $\tilde{\mathbf{u}} := [0, b]^\top$ and a corresponding control time $\tilde{t} \in [0, t_f]$, where both $b \in \mathbb{R} \setminus \{0\}$ and $\tilde{t}$ are to be chosen later. We have $\mathbf{Y}_{\mathbf{u}, (\Delta, \alpha)_\ell}(\tilde{t}) = C \mathbf{X}(\mathbf{u}, (\alpha, \Delta)_\ell; \tilde{t})$, where $\mathbf{X}(\mathbf{u}, (\alpha, \Delta)_\ell; \tilde{t})$ is the solution to (19) and $C = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix}$. Thus, we obtain

$$\mathbf{Y}_{\tilde{\mathbf{u}}, (\Delta, \alpha)_\ell}(\tilde{t}) = C e^{\tilde{t}(\Delta_\ell A + \alpha_\ell b B_y)} \mathbf{X}_0.$$

Since $\Delta_\ell A + \alpha_\ell b B_y$ is skew-symmetric, we can compute its exponential matrix explicitly. By setting $\tilde{A} := \tilde{t}(\Delta_\ell A + \alpha_\ell b B_y)$ and $x_\ell := \sqrt{\Delta_\ell^2 + \alpha_\ell^2 b^2}$, we have $e^{\tilde{A}} = I_3 + \frac{\sin(\tilde{t} x_\ell)}{\tilde{t} x_\ell} \tilde{A} + \frac{1 - \cos(\tilde{t} x_\ell)}{\tilde{t}^2 x_\ell^2} \tilde{A}^2$ (see, e.g., [38]). Since $C I_3 \mathbf{X}_0 = 0$ and

$$\tilde{A}^2 = \tilde{t}^2 \begin{bmatrix} -\Delta_\ell^2 - \alpha_\ell^2 b^2 & 0 & 0 \\ 0 & -\Delta_\ell^2 & \Delta_\ell \alpha_\ell b \\ 0 & \Delta_\ell \alpha_\ell b & -\alpha_\ell^2 b^2 \end{bmatrix},$$

we obtain

$$C \mathbf{X}(\tilde{\mathbf{u}}, \alpha_\ell; \tilde{t}) = C \left(\frac{\sin(\tilde{t} x_\ell)}{x_\ell} \begin{bmatrix} \alpha_\ell b \\ 0 \\ 0 \end{bmatrix} + \frac{1 - \cos(\tilde{t} x_\ell)}{x_\ell^2} \begin{bmatrix} 0 \\ \Delta_\ell \alpha_\ell b \\ -\alpha_\ell^2 b^2 \end{bmatrix} \right) = \begin{bmatrix} \frac{\sin(\tilde{t} x_\ell)}{x_\ell} \alpha_\ell b \\ \frac{1 - \cos(\tilde{t} x_\ell)}{x_\ell^2} \Delta_\ell \alpha_\ell b \\ \frac{1 - \cos(\tilde{t} x_\ell)}{x_\ell^2} \Delta_\ell \alpha_\ell b \end{bmatrix}.$$

Thus, we have

$$\mathbf{h}^{(K)}(\mathbf{e}^{k+1}, \tilde{\mathbf{u}}, \tilde{t}) - \mathbf{h}^{(k)}(\beta^k, \tilde{\mathbf{u}}, \tilde{t}) = b \sum_{\ell=1}^K \phi_{diff}(\ell) \alpha_\ell \begin{bmatrix} \frac{\sin(\tilde{t} x_\ell)}{x_\ell} \\ \frac{1 - \cos(\tilde{t} x_\ell)}{x_\ell^2} \\ \frac{1 - \cos(\tilde{t} x_\ell)}{x_\ell^2} \end{bmatrix} =: F(\tilde{t}).$$

Now, seeking a contradiction, assume that $\mathbf{h}^{(K)}(\mathbf{e}^{k+1}, \tilde{\mathbf{u}}, \tilde{t}) - \mathbf{h}^{(k)}(\beta^k, \tilde{\mathbf{u}}, \tilde{t}) = 0$ for all $\tilde{t} \in [0, t_f]$ and all $b \in \mathbb{R} \setminus \{0\}$. Since F is analytic in $\tilde{t}$, we obtain $F^{(k)}(\tilde{t}) = 0$ for all $k \in \mathbb{N}$ and all $\tilde{t} \in [0, t_f]$. For k odd, we have

$$F^{(k)}(\tilde{t}) = b \sum_{\ell=1}^K \phi_{diff}(\ell) \alpha_\ell (-1)^{\frac{k-1}{2}} \begin{bmatrix} x_\ell^{k-1} \cos(T x_\ell) \\ \Delta_\ell x_\ell^{k-2} \sin(T x_\ell) \end{bmatrix}. \quad (20)$$

Since $F^{(k)}(T) = 0$ for all k odd, the first component of $F^{(k)}(T)$ in (20), for different k odd, implies that

$$\underbrace{\begin{bmatrix} 1 & 1 & \dots & 1 \\ x_1^2 & x_2^2 & \dots & x_K^2 \\ x_1^4 & x_2^4 & \dots & x_K^4 \\ \vdots & \vdots & \ddots & \vdots \\ x_1^{\tilde{K}} & x_2^{\tilde{K}} & \dots & x_K^{\tilde{K}} \end{bmatrix}}_{=: D} \underbrace{\begin{bmatrix} \phi_{diff}(1) \alpha_1 \cos(T x_1) \\ \phi_{diff}(2) \alpha_2 \cos(T x_2) \\ \vdots \\ \phi_{diff}(K) \alpha_K \cos(T x_K) \end{bmatrix}}_{=: \Phi_{\tilde{t}}} = 0.$$

Notice that D is a Vandermonde matrix (see, e.g., [39]). Now, let $\tilde{K} = 2(K-1)$, meaning that $D \in \mathbb{R}^{(\frac{\tilde{K}}{2}+1) \times K}$ is a square matrix. Then, the determinant of D is given exactly by

$$\det(D) = \prod_{1 \leq i < j \leq K} (x_j^2 - x_i^2).$$

This implies that two rows of D are linearly independent if and only if $|x_i| \neq |x_j|$. Hence, $\det(D_x) = \det(D_y) \neq 0$ (and therefore $\phi_i = 0$) if and only if $|x_i| \neq |x_j|$ for $i \neq j$. Recalling that $x_\ell = \sqrt{\Delta_\ell^2 + \alpha_\ell^2 b^2}$, $|x_i| \neq |x_j|$ is equivalent to $\Delta_i^2 + \alpha_i^2 b^2 \neq \Delta_j^2 + \alpha_j^2 b^2$. For $i \neq j$ we also have $\alpha_i \neq \alpha_j$ and/or $\Delta_i \neq \Delta_j$ by definition. Since $\alpha_\ell \in [0.8, 1.2]$ and $\Delta_\ell \in \Delta_0 + 2\pi[-0.2, 0.2]$ with $\Delta_0 \geq 0.4\pi$, we obtain $\alpha_i^2 \neq \alpha_j^2$ and/or $\Delta_i^2 \neq \Delta_j^2$ for $i \neq j$. Thus, there exists $b \in \mathbb{R} \setminus 0$ such that $\Delta_i^2 + \alpha_i^2 b^2 \neq \Delta_j^2 + \alpha_j^2 b^2$ for all $i, j \in \{1, \dots, K\}$ with $i \neq j$. In conclusion, we have $|x_i| \neq |x_j|$ for $i \neq j$, which implies that $\phi_i = 0$ and therefore $\phi_{diff}(\ell) \alpha_\ell \cos(\tilde{x}_\ell) = 0$ for all $\ell \in \{0, \dots, K\}$ and all $\tilde{x} \in [0, t_f]$. However, we also have $\phi_{diff}(\tilde{\ell}) \neq 0$ by (18), $\alpha_{\tilde{\ell}} > 0$ and $x_{\tilde{\ell}} > 0$. Thus, there exists $\tilde{x} \in [0, t_f]$ such that $\phi_{diff}(\tilde{\ell}) \alpha_{\tilde{\ell}} \cos(\tilde{x}_{\tilde{\ell}}) \neq 0$, which is a contradiction. $\square$

Analogously to the proof of Theorem 2, one can show that any solution $(\mathbf{u}_1, t_1)$ to the initialization problem (15) satisfies $[W(\mathbf{u}_1, t_1)]_{1,1} > 0$. We conclude our analysis by the following theorem.

Theorem 3. Let $(\mathbf{u}_k, t_k)$, $k = 1, \dots, K$, be a set of controls and corresponding control times generated by GRA. Then problem (4) is uniquely solvable by $\beta = \beta_\star$.

Proof. Let β^k be the solution to the fitting step problem (16) for $k = 1, \dots, K-1$. By Theorem 2, the vector $\mathbf{v}_{\beta^k} = [(\beta^k)^T, -1]^T$ is not in the kernel of $[W(\mathbf{u}_{k+1}, t_{k+1})]_{[1:k+1, 1:k+1]}$ for all $k \in \{1, \dots, K-1\}$. Thus, we obtain by Theorem 1 that the matrix $W = \sum_k W(\mathbf{u}_k, t_k)$ is positive definite. Hence, problem (12) is uniquely solvable by $\beta = \beta_\star$. By equivalency of problems (12) and (4), we obtain the result. $\square$

Notice that, in the notation above, OGRA simply reorders rows and columns of the matrix W^k while attempting to find and correct its kernel. In fact, the second improvement in lines 8-9 in OGRA skips the discriminatory step only if there exists a row and column of W^k with index ℓ_{k+1} such that, by swapping ϕ_{k+1} and $\phi_{\ell_{k+1}}$, the matrix $W^k_{[1:k+1, 1:k+1]}$ is positive definite. Thus, if tol_1 is sufficiently small, one can also prove convergence of OGRA analogously to GRA.

6. Numerical results

We test GRA and OGRA on the setting described in Sec. 2. We choose a maximum control time of 160 ms, which corresponds to a normalized time $t_f = 16$. The shift of the parameter Δ is set to $\Delta_0 = 4\pi$ and the width of its interval to $4\pi\Delta_1$, with $\Delta_1 = 0.2$. We consider two different probability distributions $P_\star$, a simple Gaussian one (see panel on the left in Fig. 2) and a step distribution with three peaks (see panel on the left in Fig. 4). They are discretized by a uniform mesh of 100 points (10 points in each direction). Similarly, we discretize the set of linearly independent functions $\{\phi_j\}_{j=1}^K$ by setting $K = 100$ and $\phi_j = \mathbf{e}_j \in \mathbb{R}^{100}$ the j -th canonical vector in $\mathbb{R}^{100}$. Finally, we fix the tolerances for OGRA to be $\text{tol}_1 = 10^{-14}$ and $\text{tol}_2 = 10^{-4}$.

Now, let us briefly discuss how we solve the sub-steps of the algorithms numerically. The initialization and discriminatory step problems are solved by a second-order trust-region method. For the fitting step, we use the equivalent compact form (16). The corresponding first-order optimality system is given by

$$[W^k]_{[1:k, 1:k]} \beta = [W^k]_{[1:k, k+1]}. \quad (21)$$

Since the matrix $[W^k]_{[1:k, 1:k]}$ is symmetric and positive definite, any solution to Eq. (21) is a global solution to Eq. (16). Hence, we solve the fitting step problem by solving the linear system (21) using a solver based on the SVD. This solver first computes the SVD of $[W^k]_{[1:k, 1:k]}$, i.e., two orthogonal matrices $U, V \in \mathbb{R}^{k \times k}$ and a diagonal matrix $\Sigma \in \mathbb{R}^{k \times k}$ such that $U \Sigma V^T = [W^k]_{[1:k, 1:k]}$. To make the method more robust against numerical instabilities, it then removes all singular values that are smaller than a given tolerance, and the corresponding columns of U and V . Finally, it computes β by setting $\tilde{\beta} = V^T \beta$ and solving $\tilde{\Sigma} \tilde{\beta} = U^T [W^k]_{[1:k, k+1]}$.

After running the algorithms, we reconstruct $P_\star$ by solving problem (4). Notice that, using the notation (13)-(14), the gradient of the cost function in (4) is given by $W\beta - \sum_k \Gamma(\mathbf{u}_k, t_k)^T \mathbf{Y}_{\mathbf{u}_k}^{\text{exp}}(t_k)$, where the columns of Γ are given by the $\gamma_j(\mathbf{u}_k, t_k)$ defined in Eq. (14). We can also immediately see that the Hessian of the cost function in Eq. (4) is exactly W , which is guaranteed to be positive definite by our analysis in Sec. 5. Hence, the global solution to Eq. (4) is given by the (unique) solution to

$$W\beta = \sum_k \Gamma(\mathbf{u}_k, t_k)^T \mathbf{Y}_{\mathbf{u}_k}^{\text{exp}}(t_k). \quad (22)$$

However, in order to ensure that the coefficients of the computed solution correspond to a probability distribution (i.e., belong to $\hat{\mathbb{R}}^K$), we add the necessary constraints and solve Eq. (4) with the second-order interior point algorithm of MATLAB's *fmincon*-solver. Nonetheless, the code includes an option to solve directly Eq. (22) using a SVD solver (see Section 4).

Now, we run both GRA and OGRA on the canonical set $\{\phi_j\}_{j=1}^{100}$ of hat functions. In contrast to [31], we do not include any additional random vectors in the canonical set for OGRA and also do not remove any elements from the set during OGRA (but still reorder them). The reason for this is that we experienced for the problem of this paper that additional random elements do not improve the results and removing elements from the canonical set does not reduce the number of controls, but is more likely to make the final identification problem numerically unstable. While GRA computes 100 controls, OGRA only designs 51 by skipping 48 discriminatory steps. We then choose $P_\star$ as the Gaussian distribution in Fig. 2 (left) and compute the corresponding experimental realizations $\{\mathbf{Y}_{\mathbf{u}_k}^{\text{exp}}(t_k)\}_{k=1}^{\tilde{K}}$ for the two resulting sets of control fields, with $\tilde{K} = 100$ for GRA and $\tilde{K} = 51$ for OGRA. Reconstructing $P_\star$ as described above, we obtain the coefficient vectors β_{rec} and thereby the distributions $P_{\text{rec}} = \sum_{j=1}^{100} \beta_{\text{rec},j} \phi_j$ corresponding to GRA and OGRA, shown in Fig. 2. Looking at the errors with respect to the true distribution $P_\star$ shown in Fig. 3, we observe that GRA outperforms OGRA by one order of magnitude. However, the difference is so small that it is not visible in the reconstructed distributions. Similar results are obtained for a step distribution with three peaks in Fig. 4.

To investigate the dependence of the results on the choice of parameters, we repeat the experiment for different maximum control times, widths of the Δ -interval and $K = 400$ mesh points. First, we take a look at the number of control fields generated by OGRA in Table 5. We observe that the number of generated control fields is increasing with decreasing maximum control time and decreasing width of the Δ -interval. We also observe that the ratio between the number of GRA controls (which is equal to the number of mesh points K) and the number of OGRA controls is decreasing with an increasing number of mesh points. To validate this point, we plot the number of controls for both algorithms, for different total numbers of mesh points in Fig. 5.

An explanation of this behavior is given by the condition number of the corresponding matrices W , defined in Eq. (13), representing the compact form (12) of the final identification problem. The condition numbers corresponding to GRA and OGRA for the settings in Table 5 are shown in Tables 6 and 7. Based on our theoretical results for GRA and OGRA proving that $\tilde{K} = K$ controls are sufficient, we also add a set of fully random controls (randomized within the given bounds u_m

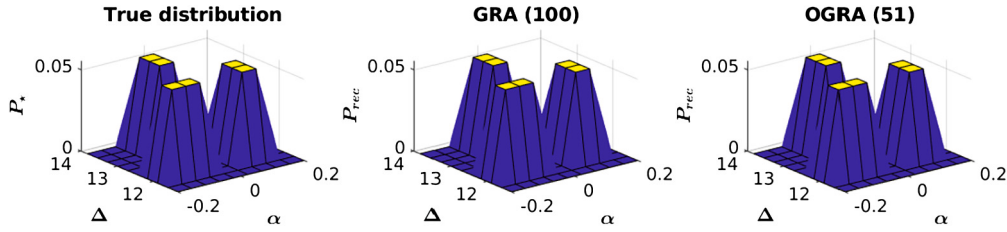


Fig. 4. Same as Fig. 2 but for a step distribution with three peaks. In brackets are the number of control fields for each set.

Table 5

Number of controls computed by OGRA for a control bound $u_m = 10$, and different numbers of discretization points K , maximum control times t_f and widths $4\pi\Delta_1$ of the Δ -interval. Bold numbers indicate that the number of OGRA controls is less than 60% of the number of GRA controls. Notice that GRA always generates K controls.

Δ_1	t_f	$K = 100$				$K = 400$			
		8	16	24	32	8	16	24	32
0.1		70	52	50	50	259	299	276	240
0.2		58	51	50	50	305	288	223	220
0.4		56	50	50	50	326	256	211	200
0.8		51	50	50	50	292	210	200	200
1.6		50	50	51	50	275	205	200	200

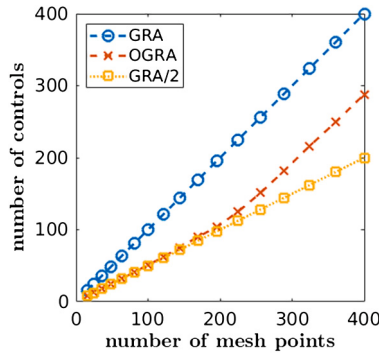


Fig. 5. Number of controls for GRA (dashed circles) and OGRA (solid crosses) for different total numbers of mesh points. To highlight the ratio between the amount of controls, we also plot half the amount of GRA controls (dotted squares).

and t_f) that has the same number of controls as GRA (i.e., $\tilde{K} = 100$ and $\tilde{K} = 400$, respectively). We observe that the condition number shows the same correlation with respect to the maximum control time, width of the Δ -interval and number of mesh points, as the number of OGRA controls. In particular, the condition number of OGRA is below $1e15$ for all settings where OGRA computed less than 60% of the number of GRA controls.

Regarding the condition numbers, GRA and random controls show the same behavior as OGRA. The reason can be found by taking a closer look at the entries of the matrix W . It can be shown that the difference between two adjacent rows or columns of W is bounded in norm by u_m , t_f and the mesh size for the probability distribution, i.e., $\alpha_{\ell+1} - \alpha_{\ell}$ and $\Delta_{\ell+1} - \Delta_{\ell}$. The interested reader can find more details about this result in Appendix A. We conclude that, if the control bound, the maximum control time, or the mesh size (or equivalently the width of the Δ -interval) is too small, the difference between two adjacent rows/columns of W can become numerically equal to zero, implying that W has a nontrivial kernel.

In order to investigate the impact of this numerical instability on the reconstructed results, we consider again the setting of the beginning of this section (i.e., $\Delta_1 = 0.2$ and $t_f = 16$), but for $K = 400$ mesh points. The results for a Gaussian and a step distribution with three peaks are

plotted in Figs. 6 and 7, respectively. We observe that all three control field sets are able to fully reconstruct the step distribution and, at least partially, the Gaussian distribution. This is because the admissible set of solutions for the final identification problem (4) is restricted to $\mathbb{R}^K$. Thus, a bad condition number does not necessarily imply that it is impossible to (at least partially) reconstruct the true probability distribution. However, a good condition number guarantees stability of the numerical solver and improves the accuracy of the results. In this context, notice that if we would sufficiently increase either the control bound u_m , or the maximum control time t_f , both GRA and OGRA would show better condition numbers and be able to perfectly reconstruct also the Gaussian distribution in Fig. 6.

We observe also that, if one knows the number of sufficient control functions $\tilde{K} = K$, then even completely random control fields can be able to perform similarly to GRA and OGRA controls. However, while OGRA finds automatically $\tilde{K}$ (reduces the number of control fields to a sufficient amount), there is no indicator for a sufficient amount of random controls in general. Additionally, the corresponding condition numbers are in many cases worse than for GRA and OGRA, as seen in Tables 6 and 7, meaning that they are more likely to show numerical instabilities. Thus, the recommended strategy is clearly OGRA, since it is able to reduce the number of control fields by up to 50%, while accurately reconstructing the probability distributions.

Lastly, we remark that making the tolerance tol_2 smaller can generally lead to even fewer controls being computed by OGRA. However, this in turn can lead to less accurate results in the reconstructed solution, meaning the user has to decide for themselves if such a trade-off is desirable.

7. Conclusion

In conclusion, we introduce SPIRED, a Greedy reconstruction algorithm to estimate spin distribution in NMR. We show that this approach can be used to jointly find the distribution of two Hamiltonian parameters, namely the offset term and the magnetic field inhomogeneity. We discuss the accuracy and limitations of this method through experimentally relevant numerical simulations. We provide and describe the codes allowing to reproduce the results of this paper. A proof of the algorithm convergence is also given.

This paper opens the way to a series of interesting questions in quantum control. A first step is to apply this algorithm to other areas in which an ensemble of quantum systems is used. Among others, we mention Bose Einstein Condensates in an optical lattice [40,11] or molecular rotational dynamics in gas phase [10,41]. The greedy reconstruction algorithm can in principle be applied to these examples, but specific constraints related to the corresponding experimental setups would be to take into account and would require adaptations of the SPIRED code. A final stage concerns the experimental implementation of this approach which seems realistic in the near future in view of the current state of the art.

CRedit authorship contribution statement

Simon Buchwald: Writing – original draft. **Gabriele Ciaramella:** Writing – original draft. **Julien Salomon:** Writing – original draft. **Dominique Sugny:** Writing – original draft.

Table 6

Condition number of W for different control sets, maximum control times t_f and widths $4\pi\Delta_1$ of the Δ -interval. The total number of mesh points is $K = 100$ and the bound on the control is $u_m = 10$. Bold numbers indicate that the condition number is smaller than $1e15$.

$\Delta_1 \backslash t_f$	GRA				OGRA				random control set			
	8	16	24	32	8	16	24	32	8	16	24	32
0.1	3e16	5e07	7e03	6e03	3e15	3e08	7e05	7e06	1e18	1e09	1e06	3e06
0.2	7e09	4e06	3e03	1e03	5e09	1e08	1e08	1e06	7e12	1e09	4e04	9e03
0.4	1e11	5e03	1e03	1e03	2e11	3e06	1e07	1e07	1e16	7e04	9e03	1e03
0.8	1e06	2e03	1e03	1e03	3e07	8e05	2e05	2e05	7e10	3e03	2e03	1e03
1.6	2e04	9e02	1e03	8e02	2e06	4e07	1e05	7e07	5e09	1e04	2e03	1e03

Table 7

Same as Table 6 but for a total number of mesh points $K = 400$.

$\Delta_1 \backslash t_f$	GRA				OGRA				random control set			
	8	16	24	32	8	16	24	32	8	16	24	32
0.1	2e20	1e19	1e20	2e15	3e19	2e19	9e19	3e15	4e19	2e19	2e19	3e18
0.2	1e19	1e19	1e14	4e13	6e19	5e19	2e14	1e14	2e19	2e19	6e19	3e18
0.4	8e19	2e19	1e14	1e04	5e19	8e18	1e14	4e07	6e19	4e19	8e15	2e06
0.8	3e19	2e13	1e04	9e03	3e19	2e14	1e07	5e07	5e19	4e18	1e08	1e05
1.6	6e20	1e10	2e04	6e03	1e20	4e11	8e07	2e08	3e19	6e18	1e08	2e04

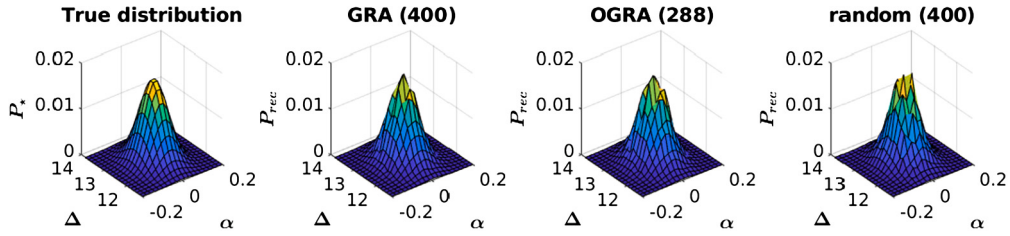


Fig. 6. Same as Fig. 2 but for $K = 400$ and including the reconstructed distribution for 400 random control fields. In brackets are the number of control fields for each set.

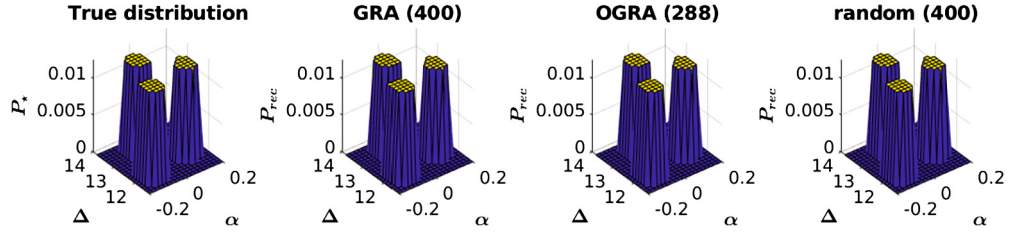


Fig. 7. Same as Fig. 6 but for a step distribution with three peaks. In brackets are the number of control fields for each set.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

Data will be made available on request.

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Appendix A. Numerical stability of the matrix W

We study in this section the numerical stability of W . Notice that for constant controls the solution to the dynamical equation (1) can be written as

$$X_{u,(\Delta,\alpha)}(t) = e^{t(\Delta A + \alpha(u_x B_x + u_y B_y))} X_0.$$

Recall that for two matrices X and Y , we have

$$\|e^Y - e^X\| \leq \|Y - X\| e^{\|X\|}.$$

Now, consider two parameter pairs $(\alpha_\ell, \Delta_\ell)$ and $(\alpha_{\ell+1}, \Delta_{\ell+1})$, and define $D_\ell := t(\Delta_\ell A + \alpha_\ell(u_x B_x + u_y B_y))$ and $D_{\ell+1} := t(\Delta_{\ell+1} A + \alpha_{\ell+1}(u_x B_x + u_y B_y))$. Since $\|X_0\| = 1$, $|u_x| \leq u_m$ and $|u_y| \leq u_m$, we obtain

$$\begin{aligned} \|X_{u,(\alpha_\ell, \Delta_\ell)}(t) - X_{u,(\alpha_{\ell+1}, \Delta_{\ell+1})}(t)\| &\leq e^{\|D_\ell\|} e^{\|D_{\ell+1}\|} \|t(\Delta_\ell A + \alpha_\ell(u_x B_x + u_y B_y)) \\ &\quad - t(\Delta_{\ell+1} A + \alpha_{\ell+1}(u_x B_x + u_y B_y))\| \\ &\leq e^{\|D_\ell\|} e^{\|D_{\ell+1}\|} t_f \left(|(\Delta_\ell - \Delta_{\ell+1})| \|A\| \right. \\ &\quad \left. + |\alpha_\ell - \alpha_{\ell+1}| u_m (\|B_x\| + \|B_y\|) \right). \end{aligned}$$

Since the exponential mapping is continuous, we have $e^{\|D_{\ell+1}\|} \rightarrow e^{\|D_{\ell}\|}$ for $\Delta_{\ell+1} \rightarrow \Delta_{\ell}$ and $\alpha_{\ell+1} \rightarrow \alpha_{\ell}$. Thus, the norm of the difference between the two solutions $\mathbf{X}_{\mathbf{u},(\alpha_{\ell},\Delta_{\ell})(t_f)}$ and $\mathbf{X}_{\mathbf{u},(\alpha_{\ell+1},\Delta_{\ell+1})(t)}$ is bounded by the differences $|\Delta_{\ell} - \Delta_{\ell+1}|$, $|\alpha_{\ell} - \alpha_{\ell+1}|$, the bound to the control u_m and the maximum control time t_f . Recalling (14) and that $\phi_j = \mathbf{e}_j$ in our example, the matrix entries of W are given by

$$W_{\ell,j} = \sum_k \langle \mathbf{Y}_{\mathbf{u}_k,(\alpha_{\ell},\Delta_{\ell})(t_k)} | \mathbf{Y}_{\mathbf{u}_k,(\alpha_j,\Delta_j)(t_k)} \rangle.$$

References

- [1] D. D'Alessandro, *Introduction to Quantum Control and Dynamics*, Chapman & Hall/CRC, Boca Raton, 2007.
- [2] U. Boscain, M. Sigalotti, D. Sugny, Introduction to the Pontryagin maximum principle for quantum optimal control, *PRX Quantum* 2 (2021) 030203, <https://doi.org/10.1103/PRXQuantum.2.030203>.
- [3] C. Altafini, F. Ticozzi, Modeling and control of quantum systems: an introduction, *IEEE Trans. Autom. Control* 57 (2012) 1898.
- [4] D. Dong, I.A. Petersen, Quantum control theory and applications: a survey, *IET Control Theory Appl.* 4 (2010) 2651.
- [5] A. Borzi, G. Ciarabella, M. Sprengel, *Formulation and Numerical Solution of Quantum Control Problems*, SIAM, Philadelphia, PA, 2017.
- [6] M. Levitt, *Spin Dynamics: Basics of Nuclear Magnetic Resonance*, Wiley, 2013.
- [7] S.J. Glaser, U. Boscain, T. Calarco, C.P. Koch, W. Köckenberger, R. Kosloff, I. Kuprov, B. Luy, S. Schirmer, T. Schulte-Herbrüggen, D. Sugny, F.K. Wilhelm, Training Schrödinger's cat: quantum optimal control, *Eur. Phys. J. D* 69 (12) (2015) 279, <https://doi.org/10.1140/epjd/e2015-60464-1>.
- [8] M. Lapert, Y. Zhang, M.A. Janich, S.J. Glaser, D. Sugny, Exploring the physical limits of saturation contrast in magnetic resonance imaging, *Sci. Rep.* 2 (2012) 589, <https://doi.org/10.1038/srep00589>.
- [9] C. Brif, R. Chakrabarti, R. Rabitz, Control of quantum phenomena: past, present and future, *New J. Phys.* 12 (2010) 075008.
- [10] C.P. Koch, U. Lemeshko, D. Sugny, Quantum control of molecular rotation, *Rev. Mod. Phys.* 91 (2019) 035005, <https://doi.org/10.1103/RevModPhys.91.035005>.
- [11] N. Dupont, G. Chatelain, L. Gabardos, M. Arnal, J. Billy, B. Peaudecerf, D. Sugny, D. Guéry-Odelin, Quantum state control of a Bose-Einstein condensate in an optical lattice, *PRX Quantum* 2 (2021) 040303, <https://doi.org/10.1103/PRXQuantum.2.040303>, <https://link.aps.org/doi/10.1103/PRXQuantum.2.040303>.
- [12] M. Lapert, G. Ferrini, D. Sugny, Optimal control of quantum superpositions in a bosonic Josephson junction, *Phys. Rev. A* 85 (2012) 023611, <https://doi.org/10.1103/PhysRevA.85.023611>.
- [13] A. Acín, I. Bloch, H. Buhman, T. Calarco, C. Eichler, J. Eisert, D. Esteve, N. Gisin, S.J. Glaser, F. Jelezko, S. Kuhr, M. Lewenstein, M.F. Riedel, P.O. Schmidt, R. Thew, A. Wallraff, I. Walmsley, F.K. Wilhelm, The quantum technologies roadmap: a European community view, *New J. Phys.* 20 (8) (2018) 080201, <https://doi.org/10.1088/1367-2630/aad1ea>.
- [14] C.P. Koch, U. Boscain, T. Calarco, G. Dirr, S. Filipp, S.J. Glaser, R. Kosloff, S. Montangero, T. Schulte-Herbrüggen, D. Sugny, F.K. Wilhelm, Quantum optimal control in quantum technologies. Strategic report on current status, visions and goals for research in Europe, *EPJ Quantum Technol.* 9 (2022) 19, <https://doi.org/10.1140/epjqt/s40507-022-00138-x>.
- [15] J.M. Geremia, H. Rabitz, Optimal identification of Hamiltonian information by closed-loop laser control of quantum systems, *Phys. Rev. Lett.* 89 (2002) 263902, <https://doi.org/10.1103/PhysRevLett.89.263902>, <https://link.aps.org/doi/10.1103/PhysRevLett.89.263902>.
- [16] C.W. Helstrom, Quantum detection and estimation theory, *J. Stat. Phys.* 1 (1969) 231–252, <https://doi.org/10.1007/BF01007479>.
- [17] V. Giovannetti, S. Lloyd, L. Maccone, Quantum-enhanced measurements: beating the standard quantum limit, *Science* 306 (5700) (2004) 1330–1336, <https://doi.org/10.1126/science.1104149>.
- [18] C.L. Degen, F. Reinhard, P. Cappellaro, Quantum sensing, *Rev. Mod. Phys.* 89 (2017) 035002, <https://doi.org/10.1103/RevModPhys.89.035002>.
- [19] Y. Maday, J. Salomon, A greedy algorithm for the identification of quantum systems, in: *Proceedings of the 48th IEEE Conference on Decision and Control*, 2009, Held jointly with the 28th Chinese Control Conference (CDC/CCC 2009), IEEE Conference on Decision and Control, 2009, pp. 375–379.
- [20] S. Conolly, D. Nishimura, A. Macovski, Optimal control solutions to the magnetic resonance selective excitation problem, *IEEE Trans. Med. Imaging* 5 (2) (1986) 106–115, <https://doi.org/10.1109/TMI.1986.4307754>.
- [21] T.-M. Zhang, R.-B. Wu, F.-H. Zhang, T.-J. Tarn, G.-L. Long, Minimum-time selective control of homonuclear spins, *IEEE Trans. Control Syst. Technol.* 23 (5) (2015) 2018–2025, <https://doi.org/10.1109/TCST.2015.2390191>.
- [22] Q. Ansel, S.J. Glaser, D. Sugny, Selective and robust time-optimal rotations of spin systems, *J. Phys. A, Math. Theor.* 54 (8) (2021) 085204, <https://doi.org/10.1088/1751-8121/abdba1>, <https://dx.doi.org/10.1088/1751-8121/abdba1>.
- [23] L. Van Damme, Q. Ansel, S.J. Glaser, D. Sugny, Time-optimal selective pulses of two uncoupled spin-1/2 particles, *Phys. Rev. A* 98 (2018) 043421, <https://doi.org/10.1103/PhysRevA.98.043421>.
- [24] J. Liu, H. Yuan, Quantum parameter estimation with optimal control, *Phys. Rev. A* 96 (2017) 012117, <https://doi.org/10.1103/PhysRevA.96.012117>.
- [25] H. Yuan, C.F. Fung, Quantum parameter estimation with general dynamics, *npj Quantum Inf.* 3 (2017) 14.
- [26] C. Lin, Y. Ma, D. Sels, Optimal control for quantum metrology via Pontryagin's principle, *Phys. Rev. A* 103 (2021) 052607, <https://doi.org/10.1103/PhysRevA.103.052607>.
- [27] J. Liu, M. Zhang, H. Chen, L. Wang, H. Yuan, Optimal scheme for quantum metrology, *Adv. Quantum Technol.* 5 (1) (2022) 2100080, <https://doi.org/10.1002/qute.202100080>.
- [28] C. Lin, Y. Ma, D. Sels, Application of Pontryagin's maximum principle to quantum metrology in dissipative systems, *Phys. Rev. A* 105 (2022) 042621, <https://doi.org/10.1103/PhysRevA.105.042621>.
- [29] D. Ma, V. Gulani, N. Seiberlich, Magnetic resonance fingerprinting, *Nature* 495 (2013) 187.
- [30] Q. Ansel, M. Tesch, S.J. Glaser, D. Sugny, Optimizing fingerprinting experiments for parameter identification: application to spin systems, *Phys. Rev. A* 96 (2017) 053419, <https://doi.org/10.1103/PhysRevA.96.053419>.
- [31] S. Buchwald, G. Ciarabella, J. Salomon, D. Sugny, Greedy reconstruction algorithm for the identification of spin distribution, *Phys. Rev. A* 104 (2021) 063112.
- [32] S. Buchwald, G. Ciarabella, J. Salomon, Analysis of a greedy reconstruction algorithm, *SIAM J. Control Optim.* 59 (6) (2021) 4511–4537, <https://doi.org/10.1137/20M1373384>.
- [33] K. Kobzar, T.E. Skinner, N. Khaneja, S.J. Glaser, B. Luy, Exploring the limits of broadband excitation and inversion: ii. rf-power optimized pulses, *J. Magn. Reson.* 194 (1) (2008) 58–66, <https://doi.org/10.1016/j.jmr.2008.05.023>.
- [34] M. Lapert, Y. Zhang, M. Braun, S.J. Glaser, D. Sugny, Singular extremals for the time-optimal control of dissipative spin $\frac{1}{2}$ particles, *Phys. Rev. Lett.* 104 (2010) 083001, <https://doi.org/10.1103/PhysRevLett.104.083001>.
- [35] N. Khaneja, R. Brockett, S.J. Glaser, *Phys. Rev. A* 63 (2001) 032308.
- [36] B. Bonnard, O. Cots, S.J. Glaser, M. Lapert, D. Sugny, Y. Zhang, Geometric optimal control of the contrast imaging problem in nuclear magnetic resonance, *IEEE Trans. Autom. Control* 57 (8) (2012) 1957–1969, <https://doi.org/10.1109/TAC.2012.2195859>.
- [37] T.E. Skinner, T.O. Reiss, B. Luy, N. Khaneja, S.J. Glaser, Tailoring the optimal control cost function to a desired output: application to minimizing phase errors in short broadband excitation pulses, *J. Magn. Reson.* 172 (1) (2005) 17–23.
- [38] Rodrigues, Des lois géométriques qui régissent les déplacements d'un système solide dans l'espace, et de la variation des coordonnées provenant de ces déplacements considérés indépendamment des causes qui peuvent les produire, *J. Math. Pures Appl.* (1840) 380–440.
- [39] K. Lundengård, Generalized Vandermonde matrices and determinants in electromagnetic compatibility, 2017.
- [40] C.A. Weidner, D.Z. Anderson, Experimental demonstration of shaken-lattice interferometry, *Phys. Rev. Lett.* 120 (2018) 263201, <https://doi.org/10.1103/PhysRevLett.120.263201>.
- [41] D. Sugny, A. Keller, O. Atabek, D. Daems, C.M. Dion, S. Guérin, H.R. Jauslin, Laser control for the optimal evolution of pure quantum states, *Phys. Rev. A* 71 (2005) 063402, <https://doi.org/10.1103/PhysRevA.71.063402>.