

Virtual Element Method for the Navier–Stokes Equation coupled with the Heat Equation

PAOLA F. ANTONIETTI

MOX-Laboratory for Modeling and Scientific Computing, Dipartimento di Matematica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133, Milano, Italy

GIUSEPPE VACCA

Department of Mathematics, University of Bari, Via Edoardo Orabona 4, 70125, Bari, Italy

AND

MARCO VERANI

MOX-Laboratory for Modeling and Scientific Computing, Dipartimento di Matematica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133, Milano, Italy

*Corresponding author: giuseppe.vacca@uniba.it

[Received on Date Month Year; revised on Date Month Year; accepted on Date Month Year]

We consider the Virtual Element discretization of the Navier–Stokes equations coupled with the heat equation where the viscosity depends on the temperature. We present the Virtual Element discretization of the coupled problem, show its well-posedness, and prove optimal error estimates. Numerical experiments which confirm the theoretical error bounds are also presented.

Keywords: Navier–Stokes equations, heat equation, Virtual Element methods.

1. Introduction

In this paper, we consider the stationary flow of a viscous incompressible fluid, in the case where the viscosity of the fluid depends on the temperature. Such a fluid-thermal coupling appears in many physical phenomena and it is involved in wide range industrial applications, e.g. cooling processes in steel industries, industrial furnaces, boilers, heat exchangers and nuclear power plants.

Due to its importance in many practical applications, the numerical approximation of coupled Navier-Stokes and heat equations have been widely studied. In [41] a generalized Boussinesq problem modelling the motion of a nonisothermal incompressible fluid subject to a heat source has been discretized by means of mixed finite element methods. A spectral discretization of the Navier-Stokes equation coupled with the heat equation has been proposed in [3], in the stationary case, and in [4] in the unsteady one. Finite Element approximation of the time dependent Boussinesq model with nonlinear viscosity depending on the temperature has been studied in [6]. Recently, the finite element approximation of the heat equation coupled with Stokes equations with nonlinear slip boundary conditions has been analyzed in [35]. Finite element methods for Darcy’s problem coupled with the heat equation has been studied in [23].

In this paper we propose to discretize the coupled thermo-fluid dynamic problem based on employing the Virtual Element method (VEM). The VEM belongs to the family of polytopal element methods, i.e., finite element methods that can support polygonal/polyhedral meshes and is a generalization of the Finite Element method that takes inspiration from modern Mimetic Finite Difference schemes [5, 18, 19]. In the context of fluid mechanics problems, the VEM has received substantial

attention in the last ten years, because it can support approximation spaces preserving the underlying physical models' fundamental properties (e.g., the divergence-free constraint for incompressible flow problems). More specifically, Virtual Element approximations of Stokes and Navier-Stokes equations have been proposed and analyzed in [9, 13, 15, 16, 44]; we refer, e.g., to [11, 20, 22], for the development of VEM for convection dominated problems. On top of the already mentioned fluid mechanics problems, we also mention that VEMs have received substantial attention also for a wide range of problems stemming from other Engineering and Applied Science applications; we refer to the recent book [10] where the current state of the art on the Virtual Element Method is presented, collecting contributions from many of the most active researchers in this field. We also mention that, in the context of fluid mechanics, other polytopal element methods have been also successfully developed in the last decades thanks to their intrinsic flexibility in handling complicated geometries; we mention, for example, polygonal FEMs, polygonal and Hybridizable Discontinuous Galerkin, Hybrid High Order methods [2, 7, 8, 27, 32, 33, 34, 39, 40].

The remaining part of the manuscript is organized as follows. Section 2 describes the coupled thermo-fluid dynamic problem together with the theoretical results. Section 3 is devoted to the Virtual Elements discretization. The existence of discrete solutions is established in 4 whereas the *a priori* convergence analysis is presented in Section 5. In Section 6 we present some numerical experiments to test the practical performance of the proposed VEM scheme. Finally, in Section 7 we draw some conclusions.

We close this introduction with some standard notation. Let $\Omega \subset \mathbb{R}^2$ be the computational domain, we denote with $\mathbf{x} = (x_1, x_2)$ the independent variable. With a usual notation the symbols ∇ and Δ denote the gradient and the Laplacian for scalar functions, while $\mathbf{\Delta}$, $\mathbf{\nabla}$, $\boldsymbol{\varepsilon}$ and div denote the vector Laplacian, the gradient, the symmetric gradient operator and the divergence operator, whereas div denotes the vector valued divergence operator for tensor fields. Throughout the paper, we will follow the usual notation for Sobolev spaces and norms [1]. Hence, for an open bounded domain ω , the norms in the spaces $W_p^s(\omega)$ and $L^p(\omega)$ are denoted by $\|\cdot\|_{W_p^s(\omega)}$ and $\|\cdot\|_{L^p(\omega)}$ respectively. Norm and seminorm in $H^s(\omega)$ are denoted respectively by $\|\cdot\|_{s,\omega}$ and $|\cdot|_{s,\omega}$, while $(\cdot, \cdot)_\omega$ and $\|\cdot\|_\omega$ denote the L^2 -inner product and the L^2 -norm (the subscript ω may be omitted when ω is the whole computational domain Ω).

2. Continuous problem

Let $\Omega \subseteq \mathbb{R}^2$ be a polygonal bounded domain with a Lipschitz-continuous boundary $\partial\Omega$. We consider the following coupled thermo-fluid dynamic problem

$$\left\{ \begin{array}{ll} -\text{div}(v(\vartheta)\boldsymbol{\varepsilon}(\mathbf{u})) + (\mathbf{\nabla}\mathbf{u})\mathbf{u} - \nabla p = \mathbf{f} & \text{in } \Omega, \\ \text{div } \mathbf{u} = 0 & \text{in } \Omega, \\ -\text{div}(\kappa \nabla \vartheta) + \mathbf{u} \cdot \nabla \vartheta = g & \text{in } \Omega, \\ \mathbf{u} = 0 & \text{on } \partial\Omega, \\ \vartheta = \vartheta_D & \text{on } \partial\Omega, \end{array} \right. \quad (2.1)$$

where the unknowns $\mathbf{u}$, p and ϑ represent the velocity, the pressure and the temperature of the fluid respectively; the function v is the temperature depending viscosity of the fluid, κ is the thermal conductivity, $\mathbf{f}$ and g stand for the external volume source terms and ϑ_D is the Dirichlet datum for the temperature.

Equation (2.1) models the stationary flow of a viscous incompressible fluid (governed by a Navier-Stokes equation), where the viscosity of the fluid depends on the temperature (governed by an elliptic equation).

For sake of simplicity we here consider Dirichlet homogeneous boundary conditions for the velocity field, different boundary conditions can be treated as well. Moreover, the analysis of the three dimensional case could be developed with similar arguments. From now on, we assume that the model data satisfy the following assumptions.

(A0) Data assumptions:

- the conductivity κ is positive and bounded, i.e. there exist $\kappa_*, \kappa^* > 0$ s.t.

$$0 < \kappa_* \leq \kappa(\mathbf{x}) \leq \kappa^* \quad \text{for all } \mathbf{x} \in \Omega;$$

- the viscosity ν is positive, bounded and Lipschitz continuous, i.e. there exist $\nu_*, \nu^*, \nu_{\text{lip}} > 0$ s.t.

$$\begin{aligned} 0 < \nu_* \leq \nu(\vartheta) \leq \nu^* & \quad \text{for all } \vartheta \in \mathbb{R}, \\ |\nu(\vartheta_1) - \nu(\vartheta_2)| \leq \nu_{\text{lip}} |\vartheta_1 - \vartheta_2| & \quad \text{for all } \vartheta_1, \vartheta_2 \in \mathbb{R}; \end{aligned}$$

- the external loads satisfy $\mathbf{f} \in [L^2(\Omega)]^2, g \in L^2(\Omega)$;
- the Dirichlet datum satisfy $\vartheta_D \in H^{1/2}(\partial\Omega)$.

Let us define the continuous spaces

$$\begin{aligned} \mathbf{V} &:= [H_0^1(\Omega)]^2, \quad Q := L_0^2(\Omega) = \left\{ q \in L^2(\Omega) \text{ s.t. } \int_{\Omega} q \, d\Omega = 0 \right\}, \\ \Sigma &:= H^1(\Omega), \quad \Sigma_0 := H_0^1(\Omega), \quad \Sigma_D := \{ \sigma \in H^1(\Omega) \text{ s.t. } \sigma|_{\partial\Omega} = \vartheta_D \}, \end{aligned}$$

endowed with natural norms, and the forms

$$a_V(\cdot; \cdot, \cdot): \Sigma \times \mathbf{V} \times \mathbf{V} \rightarrow \mathbb{R}, \quad a_V(\vartheta; \mathbf{u}, \mathbf{v}) := \int_{\Omega} \nu(\vartheta) \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v}) \, d\Omega, \quad (2.2)$$

$$a_T(\cdot, \cdot): \Sigma \times \Sigma \rightarrow \mathbb{R}, \quad a_T(\vartheta, \sigma) := \int_{\Omega} \kappa \nabla \sigma \cdot \nabla \vartheta \, d\Omega, \quad (2.3)$$

$$c_V(\cdot; \cdot, \cdot): \mathbf{V} \times \mathbf{V} \times \mathbf{V} \rightarrow \mathbb{R}, \quad c_V(\mathbf{w}; \mathbf{u}, \mathbf{v}) := \int_{\Omega} (\nabla \mathbf{u}) \mathbf{w} \cdot \mathbf{v} \, d\Omega, \quad (2.4)$$

$$c_T(\cdot; \cdot, \cdot): \mathbf{V} \times \Sigma \times \Sigma \rightarrow \mathbb{R}, \quad c_T(\mathbf{u}; \vartheta, \sigma) := \int_{\Omega} \mathbf{u} \cdot \nabla \vartheta \, \sigma \, d\Omega, \quad (2.5)$$

$$b(\cdot, \cdot): \mathbf{V} \times Q \rightarrow \mathbb{R}, \quad b(\mathbf{v}, q) := \int_{\Omega} q \operatorname{div} \mathbf{v} \, d\Omega. \quad (2.6)$$

Notice that under the assumptions (A0) the forms are well defined, furthermore the following hold

(P0) Stability properties of the continuous forms:

- for any $\vartheta \in \Sigma$ the bilinear form $a_V(\vartheta; \cdot, \cdot)$ is coercive and continuous, i.e.

$$a_V(\vartheta; \mathbf{v}, \mathbf{v}) \geq \nu_* |\mathbf{v}|_{1,\Omega}^2, \quad a_V(\vartheta; \mathbf{u}, \mathbf{v}) \leq \nu^* |\mathbf{u}|_{1,\Omega} |\mathbf{v}|_{1,\Omega}, \quad \text{for all } \mathbf{u}, \mathbf{v} \in \mathbf{V};$$

- the bilinear form $a_T(\cdot, \cdot)$ is coercive and continuous, i.e.

$$a_T(\sigma, \sigma) \geq \kappa_* |\sigma|_{1,\Omega}^2, \quad a_T(\vartheta, \sigma) \leq \kappa^* |\vartheta|_{1,\Omega} |\sigma|_{1,\Omega}, \quad \text{for all } \vartheta, \sigma \in \Sigma;$$

- the convective trilinear forms $c_V(\cdot; \cdot, \cdot)$ and $c_T(\cdot; \cdot, \cdot)$ are continuous with continuity constant C_{conv} , i.e.

$$\begin{aligned} c_V(\mathbf{w}; \mathbf{u}, \mathbf{v}) &\leq C_{\text{conv}} |\mathbf{w}|_{1,\Omega} |\mathbf{u}|_{1,\Omega} |\mathbf{v}|_{1,\Omega}, & \text{for all } \mathbf{w}, \mathbf{u}, \mathbf{v} \in \mathbf{V}, \\ c_T(\mathbf{u}; \vartheta, \sigma) &\leq C_{\text{conv}} |\mathbf{u}|_{1,\Omega} \|\vartheta\|_{1,\Omega} |\sigma|_{1,\Omega}, & \text{for all } \mathbf{u} \in \mathbf{V}, \vartheta \in W, \sigma \in \Sigma_0; \end{aligned}$$

- the bilinear form $b(\cdot, \cdot)$ is continuous and realizes the inf-sup condition with inf-sup constant $\beta > 0$, i.e.

$$\begin{aligned} b(\mathbf{v}, q) &\leq |\mathbf{v}|_{1,\Omega} \|q\|_{0,\Omega} & \text{for all } \mathbf{v} \in \mathbf{V} \text{ and } q \in Q, \\ \sup_{\mathbf{v} \in \mathbf{V}} \frac{b(\mathbf{v}, q)}{|\mathbf{v}|_{1,\Omega}} &\geq \beta \|q\|_{0,\Omega} & \text{for all } q \in Q. \end{aligned}$$

The variational formulation of Problem (2.1) reads as follows:

$$\left\{ \begin{array}{ll} \text{find } (\mathbf{u}, p, \vartheta) \in \mathbf{V} \times Q \times \Sigma_D, \text{ such that} & \\ a_V(\vartheta, \mathbf{u}, \mathbf{v}) + c_V(\mathbf{u}; \mathbf{u}, \mathbf{v}) + b(\mathbf{v}, p) = (\mathbf{f}, \mathbf{v}) & \text{for all } \mathbf{v} \in \mathbf{V}, \\ b(\mathbf{u}, q) = 0 & \text{for all } q \in Q, \\ a_T(\vartheta, \sigma) + c_T(\mathbf{u}; \vartheta, \sigma) = (g, \sigma) & \text{for all } \sigma \in \Sigma_0. \end{array} \right. \quad (2.7)$$

Let us introduce the kernel of the bilinear form $b(\cdot, \cdot)$ that corresponds to the functions in $\mathbf{V}$ with vanishing divergence, i.e.

$$\mathbf{Z} := \{\mathbf{v} \in \mathbf{V} \quad \text{s.t.} \quad \text{div } \mathbf{v} = 0\}.$$

Then, Problem (2.7) can be formulated in the equivalent kernel form:

$$\left\{ \begin{array}{ll} \text{find } (\mathbf{u}, \vartheta) \in \mathbf{Z} \times \Sigma_D, \text{ such that} & \\ a_V(\vartheta, \mathbf{u}, \mathbf{v}) + c_V(\mathbf{u}; \mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) & \text{for all } \mathbf{v} \in \mathbf{Z}, \\ a_T(\vartheta, \sigma) + c_T(\mathbf{u}; \vartheta, \sigma) = (g, \sigma) & \text{for all } \sigma \in \Sigma_0. \end{array} \right. \quad (2.8)$$

We can observe by a direct computation, that, for a fixed $\mathbf{w} \in \mathbf{Z}$, the bilinear forms $c_V(\mathbf{w}; \cdot, \cdot)$ and $c_T(\mathbf{w}; \cdot, \cdot)$ are skew symmetric, i.e.

$$\begin{aligned} c_V(\mathbf{w}; \mathbf{u}, \mathbf{v}) &= -c_V(\mathbf{w}; \mathbf{v}, \mathbf{u}) & \text{for all } \mathbf{u}, \mathbf{v} \in \mathbf{V}, \\ c_T(\mathbf{w}; \vartheta, \sigma) &= -c_T(\mathbf{w}; \sigma, \vartheta) & \text{for all } \vartheta, \sigma \in \Sigma. \end{aligned} \quad (2.9)$$

Therefore, for $\mathbf{w} \in \mathbf{Z}$, the forms $c_V(\mathbf{w}; \cdot, \cdot)$ and $c_T(\mathbf{w}; \cdot, \cdot)$ can be equivalently defined as follows

$$c_V^{\text{skew}}(\mathbf{w}; \mathbf{u}, \mathbf{v}) := \frac{1}{2} (c_V(\mathbf{w}; \mathbf{u}, \mathbf{v}) - c_V(\mathbf{w}; \mathbf{v}, \mathbf{u})) \quad \text{for all } \mathbf{u}, \mathbf{v} \in \mathbf{V}, \quad (2.10)$$

$$c_T^{\text{skew}}(\mathbf{w}; \vartheta, \sigma) := \frac{1}{2} (c_T(\mathbf{w}; \vartheta, \sigma) - c_T(\mathbf{w}; \sigma, \vartheta)) \quad \text{for all } \vartheta, \sigma \in \Sigma. \quad (2.11)$$

However, at the discrete level $c_V(\mathbf{w}; \cdot, \cdot)$ and $c_V^{\text{skew}}(\mathbf{w}; \cdot, \cdot)$ (resp. $c_T(\mathbf{w}; \cdot, \cdot)$ and $c_T^{\text{skew}}(\mathbf{w}; \cdot, \cdot)$) will lead to different bilinear forms, in general.

In the context of the analysis of incompressible flows, it is useful to introduce the concept of Helmholtz–Hodge projector (see for instance [38, Lemma 2.6] and [36, Theorem 3.3]). For every $\mathbf{w} \in [L^2(\Omega)]^2$ there exist $\mathbf{w}_0 \in H(\operatorname{div}; \Omega)$ and $\zeta \in H^1(\Omega)/\mathbb{R}$ such that

$$\mathbf{w} = \mathbf{w}_0 + \nabla \zeta, \quad (2.12)$$

where $\mathbf{w}_0$ is L^2 -orthogonal to the gradients, that is $(\mathbf{w}_0, \nabla \varphi) = 0$ for all $\varphi \in H^1(\Omega)$ (which implies, in particular, that $\mathbf{w}_0$ is solenoidal, i.e. $\operatorname{div} \mathbf{w}_0 = 0$). The orthogonal decomposition (2.12) is unique and is called Helmholtz–Hodge decomposition, and $\mathcal{P}(\mathbf{w}) := \mathbf{w}_0$ is the Helmholtz–Hodge projector of $\mathbf{w}$.

Combining the argument in [3, Theorem 2.2] and the definition of Helmholtz–Hodge projector, the following existence result holds.

Theorem 1 *Under the data assumptions (A0), Problem (2.7) admits at least a solution $(\mathbf{u}, p, \vartheta) \in \mathbf{V} \times Q \times \Sigma_{\mathbf{D}}$. Moreover the solution satisfies the bound*

$$\|\mathbf{u}\|_{1,\Omega}^2 + \|\vartheta\|_{1,\Omega}^2 \leq C_{\text{est}}^2 \left(|\mathcal{P}(\mathbf{f})|_{-1,\Omega}^2 + |g|_{-1,\Omega}^2 + \|\vartheta_{\mathbf{D}}\|_{1/2,\partial\Omega}^2 \right) \quad (2.13)$$

where the constant C_{est} depends on the domain Ω and on the constants κ_* , κ^* and v_* in the data assumptions (A0).

Assuming suitable bounds on the data of the problem and on the velocity solution it is possible to establish the following uniqueness result [3, Proposition 2.3]

Theorem 2 *Under the data assumptions (A0), there exist two positive constants C_{data} and C_{sol} depending on the domain Ω and on the constants κ_* , κ^* and v_* in the data assumptions (A0) s.t. if*

i) *the data of the problem satisfies*

$$C_{\text{data}}^2 \left(|\mathcal{P}(\mathbf{f})|_{-1,\Omega}^2 + |g|_{-1,\Omega}^2 + \|\vartheta_{\mathbf{D}}\|_{1/2,\partial\Omega}^2 \right) < 1, \quad (2.14)$$

ii) *Problem (2.7) admits a solution $(\mathbf{u}, p, \vartheta)$ with $\mathbf{u} \in W_q^1(\Omega)$ where $q > 2$ and the following bound holds*

$$C_{\text{sol}} v_{\text{lip}} \|\mathbf{u}\|_{[W_q^1(\Omega)]^2} < 1, \quad (2.15)$$

then this solution is unique.

3. Virtual Elements discretization

3.1. Notation and preliminaries

We now introduce some basic tools and notations useful in the construction and the theoretical analysis of Virtual Element Methods.

Let $\{\Omega_h\}_h$ be a sequence of decompositions of the domain $\Omega \subset \mathbb{R}^2$ into general polytopal elements E where $h := \sup_{E \in \Omega_h} h_E$. We suppose that $\{\Omega_h\}_h$ fulfils the following assumption.

(A1) Mesh assumption:

there exists a positive constant ρ such that for any $E \in \{\Omega_h\}_h$

- E is star-shaped with respect to a ball B_E of radius $\geq \rho h_E$;
- any edge e of E has length $\geq \rho h_E$.

We remark that the hypotheses above, though not too restrictive in many practical cases, could possibly be further relaxed, combining the present analysis with the studies in [21, 28, 30].

Using standard VEM notations, for $n \in \mathbb{N}$ and $s \in \mathbb{R}^+$ and $p = 1, \dots, +\infty$ let us introduce the spaces:

- $\mathbb{P}_n(\omega)$: the set of polynomials on $\omega \subset \Omega$ of degree $\leq n$ (with $\mathbb{P}_{-1}(\omega) = \{0\}$),
- $\mathbb{P}_n(\Omega_h) := \{q \in L^2(\Omega) \text{ s.t. } q|_E \in \mathbb{P}_n(E) \text{ for all } E \in \Omega_h\}$,
- $W_p^s(\Omega_h) := \{v \in L^2(\Omega) \text{ s.t. } v|_E \in W_p^s(E) \text{ for all } E \in \Omega_h\}$

equipped with the broken norm and seminorm

$$\begin{aligned} \|v\|_{W_p^s(\Omega_h)}^p &:= \sum_{E \in \Omega_h} \|v\|_{W_p^s(E)}^p, & |v|_{W_p^s(\Omega_h)}^p &:= \sum_{E \in \Omega_h} |v|_{W_p^s(E)}^p, & \text{if } 1 \leq p < \infty, \\ \|v\|_{W_p^s(\Omega_h)}^p &:= \max_{E \in \Omega_h} \|v\|_{W_p^s(E)}^p, & |v|_{W_p^s(\Omega_h)}^p &:= \max_{E \in \Omega_h} |v|_{W_p^s(E)}^p, & \text{if } p = \infty. \end{aligned}$$

Let $E \in \Omega_h$, we denote with h_E the diameter, with $|E|$ the area, with $\mathbf{x}_E = (x_{E,1}, x_{E,2})$ the centroid. A natural basis associated with the space $\mathbb{P}_n(E)$ is the set of normalized monomials

$$\mathbb{M}_n(E) := \{m_{\boldsymbol{\alpha}}, \text{ with } |\boldsymbol{\alpha}| \leq n\}$$

where, for any multi-index $\boldsymbol{\alpha} = (\alpha_1, \alpha_2) \in \mathbb{N}^2$

$$m_{\boldsymbol{\alpha}} := \prod_{i=1}^2 \left(\frac{x_i - x_{E,i}}{h_E} \right)^{\alpha_i} \quad \text{and} \quad |\boldsymbol{\alpha}| := \sum_{i=1}^2 \alpha_i.$$

Moreover for any $m \leq n$ we denote with

$$\widehat{\mathbb{P}}_{n \setminus m}(E) = \text{span} \{m_{\boldsymbol{\alpha}}, \text{ with } m+1 \leq |\boldsymbol{\alpha}| \leq n\}.$$

Furthermore, we introduce the following notations: let $\{\mathcal{X}^E\}_{E \in \Omega_h}$ be a family of forms $\mathcal{X}^E: \prod_{j=1}^{\ell} W_{p_j}^{s_j}(E) \rightarrow \mathbb{R}$, then we define

$$\mathcal{X}: \prod_{j=1}^{\ell} W_{p_j}^{s_j}(\Omega_h) \rightarrow \mathbb{R}, \quad \mathcal{X}(u_1, \dots, u_{\ell}) := \sum_{E \in \Omega_h} \mathcal{X}^E(u_1, \dots, u_{\ell}), \quad (3.1)$$

for any $u_j \in W_{p_j}^{s_j}(\Omega_h)$ and $j = 1, \dots, \ell$.

For any E , let us define the following polynomial projections:

- the **L^2 -projection** $\Pi_n^{0,E}: L^2(E) \rightarrow \mathbb{P}_n(E)$, given by

$$\int_E q_n (v - \Pi_n^{0,E} v) dE = 0 \quad \text{for all } v \in L^2(E) \text{ and } q_n \in \mathbb{P}_n(E), \quad (3.2)$$

with obvious extension for vector functions $\Pi_n^{0,E}: [L^2(E)]^2 \rightarrow [\mathbb{P}_n(E)]^2$ and tensor functions $\Pi_n^{0,E}: [L^2(E)]^{2 \times 2} \rightarrow [\mathbb{P}_n(E)]^{2 \times 2}$;

- the H^1 -seminorm projection $\Pi_n^{\nabla,E} : H^1(E) \rightarrow \mathbb{P}_n(E)$, defined by

$$\begin{cases} \int_E \nabla q_n \cdot \nabla (v - \Pi_n^{\nabla,E} v) \, dE = 0 & \text{for all } v \in H^1(E) \text{ and } q_n \in \mathbb{P}_n(E), \\ \int_{\partial E} (v - \Pi_n^{\nabla,E} v) \, ds = 0, \end{cases}$$

with extension for vector fields $\Pi_n^{\nabla,E} : [H^1(E)]^2 \rightarrow [\mathbb{P}_n(E)]^2$.

In the following the symbol $\lesssim$ will denote a bound up to a generic positive constant, independent of the mesh size h , but which may depend on Ω , on the “polynomial” order of the method k and on the regularity constant appearing in the mesh assumption **(A1)**.

We finally recall the following well know useful results:

- Poincaré inequality [42, Theorem 1.3.3]

$$\|\varphi\|_{1,\Omega} \lesssim |\varphi|_{1,\Omega} \quad \text{for any } \varphi \in H_0^1(\Omega); \quad (3.3)$$

- Sobolev embedding $H^1(\Omega) \subset L^p(\Omega)$ [42, Theorem 1.3.4]: let $2 \leq p < \infty$, then

$$\|\varphi\|_{L^p(\Omega)} \lesssim \|\varphi\|_{1,\Omega} \quad \text{for any } \varphi \in H^1(\Omega); \quad (3.4)$$

- Polynomial inverse estimate [29, Theorem 4.5.11]: let $1 \leq q, p \leq \infty$, then for any $E \in \Omega_h$

$$\|p_n\|_{L^q(E)} \lesssim h_E^{2/q-2/p} \|p_n\|_{L^p(E)} \quad \text{for any } p_n \in \mathbb{P}_n(E); \quad (3.5)$$

- Bramble-Hilbert Lemma [29, Lemma 4.3.8]: let $0 \leq t \leq s \leq n+1$, and $1 \leq q, p \leq \infty$ such that $s - 2/p > t - 2/q$, then for any $E \in \Omega_h$

$$|\varphi - \Pi_n^{0,E} \varphi|_{W_q^t(E)} \lesssim h_E^{s-t+2/q-2/p} |\varphi|_{W_p^s(E)} \quad \text{for any } \varphi \in W_p^s(E). \quad (3.6)$$

The present section is devoted to the construction of the proposed virtual elements scheme. In Subsections 3.2 and 3.3 we present the inf-sup stable divergence-free velocities-pressures pair of spaces and the temperatures spaces, respectively. In Subsection 3.4 we define the discrete computable forms. Finally in Subsection 3.5 we show the virtual elements discretization of Problems (2.7) and (2.8).

Let $k \geq 2$ be the “polynomial” order of the method. The lowest order case $k = 1$ can be treated as well using a slightly different approach [9]. We recall that, in standard finite element fashion, the VEM spaces are first defined elementwise and then assembled globally. In the following we will denote by E a general polygon having ℓ_e edges e , while $\mathbf{n}_E$ will denote the unit vector that is normal to ∂E and outward with respect to E .

3.2. Virtual Elements velocities space and pressure space

In the present section we outline an overview of the divergence-free Virtual Elements spaces for the Navier-Stokes equation [15, 16, 44].

We consider on each polygonal element $E \in \Omega_h$ the “enhanced” virtual space

$$\begin{aligned} \mathbf{V}_h(E) := \left\{ \mathbf{v}_h \in [C^0(\overline{E})]^2 \text{ s.t. } \right. & (i) \ \Delta \mathbf{v}_h + \nabla s \in \mathbf{x}^\perp \mathbb{P}_{k-1}(E), \text{ for some } s \in L_0^2(E), \\ & (ii) \ \operatorname{div} \mathbf{v}_h \in \mathbb{P}_{k-1}(E), \\ & (iii) \ \mathbf{v}_h|_e \in [\mathbb{P}_k(e)]^2 \ \forall e \in \partial E, \\ & \left. (iv) \ (\mathbf{v}_h - \Pi_k^{\nabla, E} \mathbf{v}_h, \mathbf{x}^\perp \widehat{p}_{k-1})_E = 0 \ \forall \widehat{p}_{k-1} \in \widehat{\mathbb{P}}_{k-1 \setminus k-3}(E) \right\}, \end{aligned} \quad (3.7)$$

where $\mathbf{x}^\perp = (x_2, -x_1)$. We here summarize the main properties of the space $\mathbf{V}_h(E)$ (we refer to [16] for a deeper analysis).

(P1) Polynomial inclusion: $[\mathbb{P}_k(E)]^2 \subseteq \mathbf{V}_h(E)$;

(P2) Degrees of freedom: the following linear operators $\mathbf{D}_\mathbf{V}$ constitute a set of DoFs for $\mathbf{V}_h(E)$:

D_V1 the values of $\mathbf{v}_h$ at the vertexes of the polygon E ,

D_V2 the values of $\mathbf{v}_h$ at $k-1$ distinct points of every edge $e \in \partial E$,

D_V3 the moments of $\mathbf{v}_h$

$$\frac{1}{|E|} \int_E \mathbf{v}_h \cdot \mathbf{m}^\perp m_\alpha \, dE \quad \text{for any } m_\alpha \in \mathbb{M}_{k-3}(E),$$

where $\mathbf{m}^\perp := \frac{1}{h_E} (x_2 - x_{2,E}, -x_1 + x_{1,E})$,

D_V4 the moments of $\operatorname{div} \mathbf{v}_h$

$$\frac{h_E}{|E|} \int_E \operatorname{div} \mathbf{v}_h m_\alpha \, dE \quad \text{for any } m_\alpha \in \mathbb{M}_{k-1}(E) \text{ with } |\alpha| > 0;$$

(P3) Polynomial projections: the DoFs $\mathbf{D}_\mathbf{V}$ allow us to compute the following linear operators:

$$\Pi_k^{0,E} : \mathbf{V}_h(E) \rightarrow [\mathbb{P}_k(E)]^2, \quad \Pi_{k-1}^{0,E} : \nabla \mathbf{V}_h(E) \rightarrow [\mathbb{P}_{k-1}(E)]^{2 \times 2}.$$

The global velocity space $\mathbf{V}_h$ is defined by gluing the local spaces with the obvious associated sets of global DoFs:

$$\mathbf{V}_h := \{ \mathbf{v}_h \in \mathbf{V} \text{ s.t. } \mathbf{v}_h|_E \in \mathbf{V}_h(E) \text{ for all } E \in \Omega_h \}. \quad (3.8)$$

The discrete pressure space Q_h is given by the piecewise polynomial functions of degree $k-1$, i.e.

$$Q_h := \{ q_h \in Q \text{ s.t. } q_h|_E \in \mathbb{P}_{k-1}(E) \text{ for all } E \in \Omega_h \}. \quad (3.9)$$

The couple of spaces $(\mathbf{V}_h, Q_h)$ is inf-sup stable [15], and we denote with $\widehat{\beta} > 0$ the inf-sup stability constant, i.e. $\widehat{\beta}$ is such that

$$\sup_{\mathbf{v}_h \in \mathbf{V}_h} \frac{b(\mathbf{v}_h, q_h)}{\|\mathbf{v}_h\|_{1,\Omega}} \geq \widehat{\beta} \|q_h\|_{0,E} \quad \text{for any } q_h \in Q_h. \quad (3.10)$$

Let us introduce the discrete kernel

$$\mathbf{Z}_h := \{ \mathbf{v}_h \in \mathbf{V}_h \text{ s.t. } b(\mathbf{v}_h, q_h) = 0 \text{ for all } q_h \in Q_h \} \quad (3.11)$$

then recalling (ii) in (3.7) and (3.9), the following kernel inclusion holds

$$\mathbf{Z}_h \subseteq \mathbf{Z}, \quad (3.12)$$

i.e. the functions in the discrete kernel are exactly divergence-free.

3.3. Virtual Elements space for the temperatures

In this section, we briefly introduce the H^1 -conforming virtual space for the temperatures, that consists on the so-called nodal “enhanced” virtual space [5]. We thus consider on each element $E \in \Omega_h$ the space

$$\begin{aligned} \Sigma_h(E) := \left\{ \sigma_h \in C^0(\overline{E}) \text{ s.t. } (i) \Delta \sigma_h \in \mathbb{P}_k(E), \right. \\ (ii) \sigma_h|_e \in \mathbb{P}_k(e) \ \forall e \in \partial E, \\ \left. (iii) (\sigma_h - \Pi_k^{\nabla, E} \sigma_h, \widehat{p}_k)_E = 0 \ \forall \widehat{p}_k \in \widehat{\mathbb{P}}_{k \setminus k-2}(E) \right\}. \end{aligned} \quad (3.13)$$

We here summarize the main properties of the space $\Sigma_h(E)$ (see [5]).

(P4) Polynomial inclusion: $\mathbb{P}_k(E) \subseteq \Sigma_h(E)$;

(P5) Degrees of freedom: the following linear operators $\mathbf{D}_\Sigma$ constitute a set of DoFs for $\Sigma_h(E)$:

D_{Σ1} the values of σ_h at the vertexes of the polygon E ,

D_{Σ2} the values of σ_h at $k-1$ distinct points of every edge $e \in \partial E$,

D_{Σ3} the moments of σ_h

$$\frac{1}{|E|} \int_E \sigma_h m_\alpha \, dE \quad \text{for any } m_\alpha \in \mathbb{M}_{k-2}(E);$$

(P6) Polynomial projections: the DoFs $\mathbf{D}_\Sigma$ allow us to compute the following linear operators:

$$\Pi_k^{0,E} : \Sigma_h(E) \rightarrow \mathbb{P}_k(E), \quad \Pi_{k-1}^{0,E} : \nabla \Sigma_h(E) \rightarrow [\mathbb{P}_{k-1}(E)]^2.$$

The global velocity space Σ_h is defined by gluing the local spaces with the obvious associated sets of global DoFs:

$$\Sigma_h := \{ \sigma_h \in \Sigma \text{ s.t. } \sigma_h|_E \in \Sigma_h(E) \text{ for all } E \in \Omega_h \},$$

with its homogeneous boundary counterpart

$$\Sigma_{0,h} := \{ \sigma_h \in \Sigma_h \text{ s.t. } \sigma_h|_{\partial\Omega} = 0 \}. \quad (3.14)$$

In order to treat non-homogeneous Dirichlet data, accordingly to the data assumptions **(A0)**, let us consider the “quasi”-interpolant $\vartheta_{D,h}$ of ϑ_D defined by:

$$\vartheta_{D,h} \in C^0(\partial\Omega) \quad \text{s.t.} \quad \vartheta_{D,h}|_e \in \mathbb{P}_k(e) \quad \text{for any edge } e \subset \partial\Omega. \quad (3.15)$$

Obviously, in case $\vartheta_D \in C^0(\partial\Omega)$ the usual nodal interpolant can be employed to build $\vartheta_{D,h}$.

Then the discrete counterpart of the space Σ_D is defined as follows:

$$\Sigma_{D,h} := \{ \sigma_h \in \Sigma_h \text{ s.t. } \sigma_h|_{\partial\Omega} = \vartheta_{D,h} \}. \quad (3.16)$$

3.4. Virtual Element forms

The next step in the construction of our method is the definition of discrete versions of the continuous second order forms in (2.2) and (2.3), of the convective forms in (2.4) and (2.5) and the approximation of the right-hand side terms. It is clear that for an arbitrary function in $\mathbf{V}_h$ or Σ_h the forms are not computable since the discrete functions are not known in closed form. Therefore, following the usual procedure in the VEM setting, we need to construct discrete forms that are computable by means of the only knowledge of the DoFs.

In the light of properties **(P3)** and **(P6)** we define the computable discrete forms.

- Discrete second order forms (cf. (2.2) and (2.3))

$$a_{V,h}^E(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) := \int_E \mathbf{v}(\Pi_k^{0,E} \vartheta_h) \Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{u}_h) : \Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{v}_h) dE + s_V^E(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h), \quad (3.17)$$

$$a_{T,h}^E(\vartheta_h, \sigma_h) := \int_E \kappa \Pi_{k-1}^{0,E} \nabla \vartheta_h \cdot \Pi_{k-1}^{0,E} \nabla \sigma_h dE + s_T^E(\vartheta_h, \sigma_h), \quad (3.18)$$

where the VEM stabilizing terms are given by

$$s_V^E(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) := \mathbf{v}(\Pi_0^{0,E} \vartheta_h) S_V^E((I - \Pi_k^{0,E}) \mathbf{u}_h, (I - \Pi_k^{0,E}) \mathbf{v}_h), \quad (3.19)$$

$$s_T^E(\vartheta_h, \sigma_h) := (\Pi_0^{0,E} \kappa) S_T^E((I - \Pi_k^{0,E}) \vartheta_h, (I - \Pi_k^{0,E}) \sigma_h), \quad (3.20)$$

and where $S_V^E(\cdot, \cdot) : \mathbf{V}_h(E) \times \mathbf{V}_h(E) \rightarrow \mathbb{R}$ and $S_T^E(\cdot, \cdot) : \Sigma_h(E) \times \Sigma_h(E) \rightarrow \mathbb{R}$ are computable symmetric discrete forms satisfying

$$\begin{aligned} |\mathbf{v}_h|_{1,E}^2 &\lesssim S_V^E(\mathbf{v}_h, \mathbf{v}_h) \lesssim |\mathbf{v}_h|_{1,E}^2 \quad \text{for all } \mathbf{v} \in \mathbf{V}_h(E) \cap \ker \Pi_k^{0,E}, \\ |\sigma_h|_{1,E}^2 &\lesssim S_T^E(\sigma_h, \sigma_h) \lesssim |\sigma_h|_{1,E}^2 \quad \text{for all } \sigma_h \in \Sigma_h(E) \cap \ker \Pi_k^{0,E}. \end{aligned} \quad (3.21)$$

Many examples can be found in the VEM literature [12, 18]. In the present paper we consider the so-called *dof i-dof i* stabilization defined as follows: let $\vec{\mathbf{u}}_h, \vec{\mathbf{v}}_h$ and $\vec{\vartheta}_h, \vec{\sigma}_h$ denote the real valued vectors containing the values of the local degrees of freedom associated to $\mathbf{u}_h, \mathbf{v}_h$ in the space $\mathbf{V}_h$ and to ϑ_h and σ_h in the space Σ_h respectively, then

$$S_V^E(\mathbf{u}_h, \mathbf{v}_h) := \vec{\mathbf{u}}_h \cdot \vec{\mathbf{v}}_h, \quad S_T^E(\vartheta_h, \sigma_h) := \vec{\vartheta}_h \cdot \vec{\sigma}_h.$$

- Discrete convective forms (cf. (2.4) and (2.5))

$$c_{V,h}^E(\mathbf{w}_h; \mathbf{u}_h, \mathbf{v}_h) := \int_E [\Pi_{k-1}^{0,E} (\nabla \mathbf{u}_h) \Pi_k^{0,E} \mathbf{w}_h] \cdot \Pi_k^{0,E} \mathbf{v}_h dE, \quad (3.22)$$

$$c_{T,h}^E(\mathbf{u}_h; \vartheta_h, \sigma_h) := \int_E (\Pi_k^{0,E} \mathbf{u}_h \cdot \Pi_{k-1}^{0,E} \nabla \vartheta_h) \Pi_k^{0,E} \sigma_h dE, \quad (3.23)$$

and their skew-symmetric formulation (cf. (2.10) and (2.11))

$$c_{V,h}^{\text{skew},E}(\mathbf{w}_h; \mathbf{u}_h, \mathbf{v}_h) := \frac{1}{2} (c_{V,h}^E(\mathbf{w}_h; \mathbf{u}_h, \mathbf{v}_h) - c_{V,h}^E(\mathbf{w}_h; \mathbf{v}_h, \mathbf{u}_h)), \quad (3.24)$$

$$c_{T,h}^{\text{skew},E}(\mathbf{w}_h; \vartheta_h, \sigma_h) := \frac{1}{2} (c_{T,h}^E(\mathbf{w}_h; \vartheta_h, \sigma_h) - c_{T,h}^E(\mathbf{w}_h; \sigma_h, \vartheta_h)). \quad (3.25)$$

- Discrete external forces

$$\mathbf{f}_{h|E} := \Pi_k^{0,E} \mathbf{f}, \quad g_{h|E} := \Pi_k^{0,E} g, \quad \text{for all } E \in \Omega_h. \quad (3.26)$$

The global forms, defined accordingly to the notation (3.1), satisfy the following.

(P7) Stability properties of the discrete forms:

- for any $\vartheta \in \Sigma$ the bilinear form $a_{V,h}(\vartheta; \cdot, \cdot)$ is coercive and continuous, i.e.

$$a_{V,h}(\vartheta; \mathbf{v}_h, \mathbf{v}_h) \geq \alpha_* \mathbf{v}_* |\mathbf{v}_h|_{1,\Omega}^2, \quad a_{V,h}(\vartheta; \mathbf{u}_h, \mathbf{v}_h) \leq \alpha^* \mathbf{v}^* |\mathbf{u}_h|_{1,\Omega} |\mathbf{v}_h|_{1,\Omega},$$

for all $\mathbf{u}_h, \mathbf{v}_h \in \mathbf{V}_h$;

- the bilinear form $a_{T,h}(\cdot, \cdot)$ is coercive and continuous, i.e.

$$a_{T,h}(\sigma_h, \sigma_h) \geq \beta_* \kappa_* |\sigma_h|_{1,\Omega}^2, \quad a_T(\vartheta_h, \sigma_h) \leq \beta^* \kappa^* |\vartheta_h|_{1,\Omega} |\sigma_h|_{1,\Omega},$$

for all $\vartheta_h, \sigma_h \in \Sigma_h$;

- the convective trilinear forms $c_{V,h}^{\text{skew}}(\cdot; \cdot, \cdot)$ and $c_{T,h}^{\text{skew}}(\cdot; \cdot, \cdot)$ are continuous, i.e.

$$\begin{aligned} c_{V,h}^{\text{skew}}(\mathbf{w}; \mathbf{u}, \mathbf{v}) &\leq \widehat{C}_{\text{conv}} |\mathbf{w}|_{1,\Omega} |\mathbf{u}|_{1,\Omega} |\mathbf{v}|_{1,\Omega}, & \text{for all } \mathbf{w}, \mathbf{u}, \mathbf{v} \in \mathbf{V}, \\ c_{T,h}^{\text{skew}}(\mathbf{u}; \vartheta, \sigma) &\leq \widehat{C}_{\text{conv}} |\mathbf{u}|_{1,\Omega} \|\vartheta\|_{1,\Omega} |\sigma|_{1,\Omega}, & \text{for all } \mathbf{u} \in \mathbf{V}, \vartheta \in \Sigma, \sigma \in \Sigma_0, \end{aligned}$$

where α_* , α^* , β_* , β^* , $\widehat{C}_{\text{conv}}$ are positive constants depending on the domain Ω , on the “polynomial” order of the method k and on the regularity constant appearing in the mesh assumption (A1).

Remark 1 Notice that the discrete convective forms $c_{V,h}^{\text{skew}}(\cdot; \cdot, \cdot)$ $c_{T,h}^{\text{skew}}(\cdot; \cdot, \cdot)$ in (3.24) and (3.25) represent the VEM counterpart of (2.10) and (2.11) respectively, in particular are forced to be skew-symmetric. Indeed, even if $\mathbf{w}_h \in \mathbf{Z}_h$, the forms $c_{V,h}^E(\mathbf{w}_h; \cdot, \cdot)$ $c_{T,h}^E(\mathbf{w}_h; \cdot, \cdot)$ in (3.22) and (3.23), as opposed to the continuous case, in principle are not skew-symmetric, since the projections can spoil the divergence-free property of the discrete functions. Different choices for the approximation of the trilinear forms have been proposed in [11, 20] for convection dominated problems.

3.5. Virtual Element problem

Having in mind the spaces (3.8), (3.9), (3.16) and (3.14), the discrete forms (3.17), (3.18), (3.24), (3.25), the form (2.6), the discrete loading terms (3.26), the virtual element discretization of Problem (2.7) is given by

$$\left\{ \begin{array}{ll} \text{find } (\mathbf{u}_h, p_h, \vartheta_h) \in \mathbf{V}_h \times Q_h \times \Sigma_{D,h}, \text{ such that} & \\ a_{V,h}(\vartheta_h, \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) + b(\mathbf{v}_h, p_h) = (\mathbf{f}_h, \mathbf{v}_h) & \forall \mathbf{v}_h \in \mathbf{V}_h, \\ b(\mathbf{u}_h, q_h) = 0 & \forall q_h \in Q_h, \\ a_{T,h}(\vartheta_h, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h, \sigma_h) = (g_h, \sigma_h) & \forall \sigma_h \in \Sigma_{0,h}. \end{array} \right. \quad (3.27)$$

Recalling the definition of the discrete kernel $\mathbf{Z}_h$ in (3.11), the previous problem can be also written in the kernel formulation, i.e.

$$\left\{ \begin{array}{ll} \text{find } (\mathbf{u}_h, \vartheta_h) \in \mathbf{Z}_h \times \Sigma_{D,h}, \text{ such that} & \\ a_{V,h}(\vartheta_h, \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) = (\mathbf{f}_h, \mathbf{v}_h) & \forall \mathbf{v}_h \in \mathbf{Z}_h, \\ a_{T,h}(\vartheta_h, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h, \sigma_h) = (g_h, \sigma_h) & \forall \sigma_h \in \Sigma_{0,h}. \end{array} \right. \quad (3.28)$$

4. Existence of discrete solutions

In the present section we establish the existence of discrete solutions.

Proposition 3 *Under the data assumptions (A0), there exists C_{exi} depending on the domain Ω and the shape regularity constant in assumption (A1) s.t. if*

$$\|\vartheta_{D,h}\|_{1/2,\partial\Omega} \leq C_{\text{exi}} \min\{\nu_*, \kappa_*\} \quad (4.1)$$

then Problem (3.27) admits at least a solution $(\mathbf{u}_h, p_h, \vartheta_h) \in \mathbf{Z}_h \times Q_h \times \Sigma_{D,h}$. Moreover, the following bound holds

$$|\mathbf{u}_h|_{1,\Omega}^2 + \|\vartheta_h\|_{1,\Omega}^2 \leq \widehat{C}_{\text{est}}^2 \left(\|\mathbf{f}_h\|_{\mathbf{Z}_h^*}^2 + |g|_{-1,\Omega}^2 + \|\vartheta_D\|_{1/2,\partial\Omega}^2 \right) \quad (4.2)$$

where the constant $\widehat{C}_{\text{est}}$ depends on the domain Ω and on the constants κ_* , κ^* and ν_* in the data assumptions (A0) and the shape regularity constant in assumption (A1).

Proof. Let $\vartheta_h^\partial \in \Sigma_{D,h}$ be the solution of the following auxiliary problem

$$\begin{cases} \text{find } \vartheta_h^\partial \in \Sigma_{D,h}, \text{ such that} \\ \int_{\Omega} \nabla \vartheta_h^\partial \cdot \nabla \sigma_h \, d\Omega = 0 \quad \forall \sigma_h \in \Sigma_{0,h}. \end{cases}$$

Then, by Lax-Milgram Lemma, ϑ_h^∂ satisfies the bound

$$\|\vartheta_h^\partial\|_{1,\Omega} \leq C_{\Omega} \|\vartheta_{D,h}\|_{1/2,\partial\Omega}. \quad (4.3)$$

Notice that Problem (3.28) can be equivalently formulated as follows

$$\begin{cases} \text{find } (\mathbf{u}_h, \vartheta_h^o) \in \mathbf{Z}_h \times \Sigma_{0,h}, \text{ such that} \\ a_{V,h}(\vartheta_h^o + \vartheta_h^\partial, \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) = (\mathbf{f}_h, \mathbf{v}_h) \quad \forall \mathbf{v}_h \in \mathbf{Z}_h, \\ a_{T,h}(\vartheta_h^o + \vartheta_h^\partial, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h^o + \vartheta_h^\partial, \sigma_h) = (g_h, \sigma_h) \quad \forall \sigma_h \in \Sigma_{0,h}. \end{cases} \quad (4.4)$$

We now prove that Problem (4.4) admits solutions. Let $\beth := \mathbf{Z}_h \times \Sigma_{0,h}$ and, for any $(\mathbf{u}_h, \vartheta_h^o) \in \beth$ let $\mathcal{A}(\mathbf{u}_h, \vartheta_h^o) : \beth \rightarrow \mathbb{R}$ given by

$$\begin{aligned} \mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{v}_h, \sigma_h) &:= a_{V,h}(\vartheta_h^o + \vartheta_h^\partial, \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) - (\mathbf{f}_h, \mathbf{v}_h) \\ &\quad + a_{T,h}(\vartheta_h^o + \vartheta_h^\partial, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h^o + \vartheta_h^\partial, \sigma_h) - (g_h, \sigma_h) \end{aligned} \quad (4.5)$$

for all $(\mathbf{v}_h, \sigma_h) \in \beth$. Then employing property (P7) and the skew-symmetry of the convective forms $c_{V,h}^{\text{skew}}$ and $c_{T,h}^{\text{skew}}$ (cf. (3.24) and (3.25)) we obtain

$$\begin{aligned} \mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{u}_h, \vartheta_h^o) &\geq \alpha_* \nu_* |\mathbf{u}_h|_{1,\Omega}^2 + \beta_* \kappa_* |\vartheta_h^o|_{1,\Omega}^2 - \|\mathbf{f}_h\|_{\mathbf{Z}_h^*} |\mathbf{u}_h|_{1,\Omega} \\ &\quad - |g_h|_{-1,\Omega} |\vartheta_h^o|_{-1,\Omega} + a_{T,h}(\vartheta_h^\partial, \vartheta_h^o) + c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h^\partial, \vartheta_h^o). \end{aligned} \quad (4.6)$$

From property **(P7)** and (4.3) it holds that

$$\begin{aligned} a_{T,h}(\vartheta_h^\partial, \vartheta_h^o) &\geq -a_{T,h}(\vartheta_h^\partial, \vartheta_h^\partial)^{1/2} a_{T,h}(\vartheta_h^o, \vartheta_h^o)^{1/2} \geq -\beta^* \kappa^* C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega} |\vartheta_h^o|_{1,\Omega}, \\ c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h^\partial, \vartheta_h^o) &\geq -\widehat{C}_{\text{conv}} \|\vartheta_h^\partial\|_{1,\Omega} |\mathbf{u}_h|_{1,\Omega} |\vartheta_h^o|_{1,\Omega} \geq -\frac{\widehat{C}_{\text{conv}}}{2} C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega} (|\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2). \end{aligned} \quad (4.7)$$

Therefore inserting (4.7) in (4.6), we infer

$$\begin{aligned} \mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{u}_h, \vartheta_h^o) &\geq \left(\min\{\alpha_*, \beta_*\} \min\{v_*, \kappa_*\} - \frac{1}{2} \widehat{C}_{\text{conv}} C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega} \right) (|\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2) + \\ &\quad - \left((\|\mathbf{f}_h\|_{\mathbf{Z}_h^*}^2 + |g_h|_{-1,\Omega}^2)^{1/2} + \beta^* \kappa^* C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega} \right) (|\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2)^{1/2}. \end{aligned} \quad (4.8)$$

Let us set

$$C_{\text{exi}} \leq \frac{\min\{\alpha_*, \beta_*\}}{\widehat{C}_{\text{conv}} C_\Omega},$$

then from (4.1) and (4.8) we get

$$\begin{aligned} \mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{u}_h, \vartheta_h^o) &\geq \frac{1}{2} \min\{\alpha_*, \beta_*\} \min\{v_*, \kappa_*\} (|\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2) + \\ &\quad - \left((\|\mathbf{f}_h\|_{\mathbf{Z}_h^*}^2 + |g_h|_{-1,\Omega}^2)^{1/2} + \beta^* \kappa^* C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega} \right) (|\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2)^{1/2}. \end{aligned} \quad (4.9)$$

Let

$$\rho := 2 \frac{(\|\mathbf{f}_h\|_{\mathbf{Z}_h^*}^2 + |g_h|_{-1,\Omega}^2)^{1/2} + \beta^* \kappa^* C_\Omega \|\vartheta_{D,h}\|_{1/2,\Omega}}{\min\{\alpha_*, \beta_*\} \min\{v_*, \kappa_*\}} \quad (4.10)$$

and $\mathcal{S} := \{(\mathbf{v}_h, \sigma_h) \in \mathfrak{D} \mid |\mathbf{u}_h|_{1,\Omega}^2 + |\vartheta_h^o|_{1,\Omega}^2 \leq \rho^2\}$ then (4.9) implies

$$\mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{u}_h, \vartheta_h^o) \geq 0 \quad \text{for any } (\mathbf{u}_h, \vartheta_h^o) \in \partial \mathcal{S}.$$

Employing the fixed-point Theorem [37, Chap. IV, Corollary 1.1] there exists $(\mathbf{u}_h, \vartheta_h^o) \in \mathcal{S}$ s.t. $\mathcal{A}(\mathbf{u}_h, \vartheta_h^o)(\mathbf{v}_h, \sigma_h) = 0$ for all $(\mathbf{v}_h, \sigma_h) \in \mathfrak{D}$, i.e. Problem (4.4) admits solution. Let $\vartheta_h = \vartheta_h^o + \vartheta_h^\partial$, then $(\mathbf{u}_h, \vartheta_h) \in \mathbf{Z}_h \times \Sigma_{0,h}$ is a solution of Problem (3.28). Bound (4.2) follows from (4.10) and (4.3) and the definition of $\mathcal{S}$.

Finally, the existence of solutions $(\mathbf{u}_h, p_h, \vartheta_h^o) \in \mathbf{V}_h \times Q_h \times \Sigma_{0,h}$ of the equivalent Problem (3.27) is a direct consequence of (3.10) (see for instance [24]). $\square$

Remark 2 Notice that, differently from the continuous problem, in order to establish the existence result for the discrete problem (3.27) we need that the conductivity κ and the viscosity v are sufficiently large with respect to the Dirichlet datum ϑ_D (cf. bound (4.1)). A careful inspection of the proof of Theorem 1 (see [3, Theorem 2.2]) shows that the differences between the continuous and the discrete case are:

1. the different construction of the (skew-symmetric) convective forms $c_T(\mathbf{u}; \cdot, \cdot)$ and $c_{T,h}^{\text{skew}}(\mathbf{u}_h; \cdot, \cdot)$;
2. the presence of the projection terms within the trilinear forms;

3. *the existence by the Hopf Lemma of an extension ϑ^∂ of the boundary datum ϑ_D s.t. $\|\vartheta^\partial\|_{1,\Omega} \lesssim \|\vartheta_D\|_{1/2,\partial\Omega}$ and $\|\vartheta^\partial\|_{L^q(\Omega)}$ arbitrary small (for $2 < q < +\infty$).*

Item 2 in the list above represents the main obstacle in the proof and makes condition (4.1) a peculiarity of VEM w.r.t. standard inf-sup stable FEM.

We now investigate whether Problem (3.28) admits a unique solution. We introduce the analysis with the following lemma.

Lemma 1 *Let $\sigma_{h,1}, \sigma_{h,2} \in \Sigma_{D,h}$, let $\mathbf{v}_h \in \mathbf{V}_h$ such that $\mathbf{v}_h \in [W_q^1(\Omega_h)]^2$ with $q > 2$. There exists a positive constant $\widehat{C}_{\text{lip}}$ only depending on Ω , k , and the shape regularity constant in (A1) such that for any $\mathbf{w}_h \in \mathbf{V}_h$*

$$|a_{V,h}(\sigma_{h,1}; \mathbf{v}_h, \mathbf{w}_h) - a_{V,h}(\sigma_{h,2}; \mathbf{v}_h, \mathbf{w}_h)| \leq \widehat{C}_{\text{lip}} v_{\text{lip}} |\sigma_{h,1} - \sigma_{h,2}|_{1,\Omega} |\mathbf{v}_h|_{[W_q^1(\Omega)]^2} |\mathbf{w}_h|_{1,\Omega}. \quad (4.11)$$

Proof. From definition of $a_{V,h}(\cdot, \cdot)$ (cf. (3.17)) we have

$$\begin{aligned} |a_{V,h}(\sigma_{h,1}; \mathbf{v}_h, \mathbf{w}_h) - a_{V,h}(\sigma_{h,2}; \mathbf{v}_h, \mathbf{w}_h)| &\leq \sum_{E \in \Omega_h} |s_V^E(\sigma_{h,1}; \mathbf{v}_h, \mathbf{w}_h) - s_V^E(\sigma_{h,2}; \mathbf{v}_h, \mathbf{w}_h)| + \\ &+ \sum_{E \in \Omega_h} \int_E |v(\Pi_k^{0,E} \sigma_{h,1}) - v(\Pi_k^{0,E} \sigma_{h,2})| |\Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{v}_h)| |\Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{w}_h)| dE := \eta_1 + \eta_2. \end{aligned} \quad (4.12)$$

Concerning the term η_1 , from (3.19) and (3.21), and recalling property (P0), we infer

$$\begin{aligned} \eta_1 &\leq \sum_{E \in \Omega_h} |v(\Pi_0^{0,E} \sigma_{h,1}) - v(\Pi_0^{0,E} \sigma_{h,2})| |S_V^E((I - \Pi_k^{0,E}) \mathbf{v}_h, (I - \Pi_k^{0,E}) \mathbf{w}_h)| \\ &\lesssim v_{\text{lip}} \sum_{E \in \Omega_h} |\Pi_0^{0,E}(\sigma_{h,1} - \sigma_{h,2})| |(I - \Pi_k^{0,E}) \mathbf{v}_h|_{1,E} |(I - \Pi_k^{0,E}) \mathbf{w}_h|_{1,E}. \end{aligned}$$

Then employing the inverse estimate (3.5) and the Bramble-Hilbert Lemma (3.6)

$$\eta_1 \lesssim v_{\text{lip}} \sum_{E \in \Omega_h} h_E^{-2/p} \|\Pi_0^{0,E}(\sigma_{h,1} - \sigma_{h,2})\|_{L^p(E)} h_E^{1-2/q} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2} |\mathbf{w}_h|_{1,E}$$

where p is such that $1/p + 1/q = 1/2$. Therefore, since $\Pi_0^{0,E}$ is continuous with respect to the L^p -norm, from the Hölder inequality for sequences we infer

$$\begin{aligned} \eta_1 &\lesssim v_{\text{lip}} \sum_{E \in \Omega_h} \|\sigma_{h,1} - \sigma_{h,2}\|_{L^p(E)} |\mathbf{v}_h|_{[W_q^1(E)]^2} |\mathbf{w}_h|_{1,E} \\ &\lesssim v_{\text{lip}} \left(\sum_{E \in \Omega_h} \|\sigma_{h,1} - \sigma_{h,2}\|_{L^p(E)}^p \right)^{1/p} \left(\sum_{E \in \Omega_h} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2}^q \right)^{1/q} \left(\sum_{E \in \Omega_h} |\mathbf{w}_h|_{1,E}^2 \right)^{1/2} \\ &\lesssim v_{\text{lip}} \|\sigma_{h,1} - \sigma_{h,2}\|_{L^p(\Omega)} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2} |\mathbf{w}_h|_{1,\Omega}. \end{aligned}$$

Then, since $p \geq 2$, from the Sobolev embedding (3.4) we obtain

$$\eta_1 \lesssim v_{\text{lip}} |\sigma_{h,1} - \sigma_{h,2}|_{1,\Omega} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2} |\mathbf{w}_h|_{1,\Omega}. \quad (4.13)$$

For what concerns the term η_2 in (4.12), using similar argument to that used for η_1 , we have

$$\begin{aligned} \eta_2 &\leq \sum_{E \in \Omega_h} \int_E |\nu(\Pi_k^{0,E} \sigma_{h,1}) - \nu(\Pi_k^{0,E} \sigma_{h,2})| |\Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{v}_h)| |\Pi_{k-1}^{0,E} \boldsymbol{\varepsilon}(\mathbf{w}_h)| dE \\ &\lesssim \nu_{\text{lip}} \sum_{E \in \Omega_h} \|\Pi_k^{0,E}(\sigma_{h,1} - \sigma_{h,2})\|_{L^p(E)} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2} |\mathbf{w}_h|_{1,E} \\ &\lesssim \nu_{\text{lip}} |\sigma_{h,1} - \sigma_{h,2}|_{1,\Omega} |\mathbf{v}_h|_{[W_q^1(\Omega_h)]^2} |\mathbf{w}_h|_{1,\Omega}. \end{aligned} \quad (4.14)$$

The thesis follows combining (4.13) and (4.14) in (4.12). $\square$

We now state the following result concerning the uniqueness of the discrete solution, which is obtained under suitable assumptions (cf. (4.15)-(4.17) and Remark 3).

Proposition 4 *Under the data assumptions (A0), there exist two positive constants $\widehat{C}_{\text{data}}$ and $\widehat{C}_{\text{sol}}$ depending on the domain Ω , on k , on the constants κ_* , κ^* and ν_* in the data assumptions (A0) and on the shape regularity constant in assumption (A1) s.t. if*

i) *the data of the problem satisfies*

$$\widehat{C}_{\text{data}}^2 \left(|\mathbf{f}_h|_{\mathbf{Z}_h^*}^2 + |g_h|_{1,\Omega}^2 + \|\vartheta_{\text{D},h}\|_{1/2,\partial\Omega}^2 \right) < 1, \quad (4.15)$$

ii) *Problem (3.27) admits a solution $(\mathbf{u}_h, p_h, \vartheta_h^o)$ with $\mathbf{u}_h \in [W_q^1(\Omega_h)]^2$ where $q > 2$ and the following bound holds*

$$\widehat{C}_{\text{sol}} \nu_{\text{lip}} |\mathbf{u}_h|_{[W_q^1(\Omega_h)]^2} < 1, \quad (4.16)$$

then, referring to (4.2) and (4.11), this solution is unique, provided that

$$K := \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\alpha_* \nu_* \widehat{C}_{\text{data}}} \left(1 + \frac{\widehat{C}_{\text{lip}}}{\beta_* \kappa_* \widehat{C}_{\text{sol}}} \right) < 1. \quad (4.17)$$

Proof. We limit to sketch the proof since it follows the guidelines of the proof of [3, Proposition 2.3]. Let $(\mathbf{u}_{h,i}, \vartheta_{h,i}) \in \mathbf{Z}_h \times \Sigma_{\text{D},h}$ with $i = 1, 2$ two solutions of Problem (3.28) and let

$$\boldsymbol{\delta}_h := \mathbf{u}_{h,1} - \mathbf{u}_{h,2} \in \mathbf{Z}_h, \quad \zeta_h := \vartheta_{h,1} - \vartheta_{h,2} \in \Sigma_{0,h}.$$

Then from (3.28), employing the skew-symmetry of the convective forms $c_{V,h}^{\text{skew}}$ and $c_{T,h}^{\text{skew}}$, we infer

$$a_{T,h}(\zeta_h, \zeta_h) = -c_{T,h}^{\text{skew}}(\boldsymbol{\delta}_h; \vartheta_{h,1}, \zeta_h), \quad (4.18)$$

and

$$a_{V,h}(\vartheta_{h,2}, \boldsymbol{\delta}_h, \boldsymbol{\delta}_h) = -c_{V,h}^{\text{skew}}(\boldsymbol{\delta}_h; \mathbf{u}_{h,1}, \boldsymbol{\delta}_h) - (a_{V,h}(\vartheta_{h,1}, \mathbf{u}_{h,1}, \boldsymbol{\delta}_h) - a_{V,h}(\vartheta_{h,2}, \mathbf{u}_{h,1}, \boldsymbol{\delta}_h)). \quad (4.19)$$

Recalling properties (P7), from (4.18) we derive

$$\beta_* \kappa_* |\zeta_h|_{1,\Omega}^2 \leq \widehat{C}_{\text{conv}} |\boldsymbol{\delta}_h|_{1,\Omega} \|\vartheta_{h,1}\|_{1,\Omega} |\zeta_h|_{1,\Omega}.$$

Thus recalling that $(\mathbf{u}_{h,1}, \mathbf{v}_{h,1})$ satisfies bound (4.2), from (4.15)

$$|\zeta_h|_{1,\Omega} \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}} \beta_* \kappa_*} |\boldsymbol{\delta}_h|_{1,\Omega}. \quad (4.20)$$

In a similar way, employing properties (P7) and Lemma 1, from (4.19) we infer

$$\alpha_* \mathbf{v}_* |\boldsymbol{\delta}_h|_{1,\Omega}^2 \leq \widehat{C}_{\text{conv}} |\mathbf{u}_{h,1}|_{1,\Omega} |\boldsymbol{\delta}_h|_{1,\Omega}^2 + \widehat{C}_{\text{lip}} \mathbf{v}_{\text{lip}} |\mathbf{u}_{h,1}|_{[W_q^1(\Omega_h)]^2} |\zeta_h|_{1,\Omega} |\boldsymbol{\delta}_h|_{1,\Omega}.$$

Using again (4.2) and (4.15), from (4.16) and (4.20) we derive

$$\alpha_* \mathbf{v}_* |\boldsymbol{\delta}_h|_{1,\Omega}^2 \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} |\boldsymbol{\delta}_h|_{1,\Omega}^2 + \frac{\widehat{C}_{\text{lip}}}{\widehat{C}_{\text{sol}}} \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}} \beta_* \kappa_*} |\boldsymbol{\delta}_h|_{1,\Omega}^2,$$

that is

$$|\boldsymbol{\delta}_h|_{1,\Omega}^2 \leq K |\boldsymbol{\delta}_h|_{1,\Omega}^2.$$

Now the proof follows from (4.17). $\square$

Remark 3 (On (4.16)) *The validity of (4.16) requires the existence of a sufficiently small positive constant $\widehat{C}_{\text{sol}}$, which, in combination with the validity of (4.17), asks for the existence of a sufficiently large constant $\widehat{C}_{\text{data}}$. This latter requirement turns out to be equivalent to assuming sufficiently small data (cf. (4.15)).*

5. Convergence results: a priori analysis

In present section we derive the rate of convergence for the proposed VEM scheme (3.27). We preliminary recall some optimal approximation properties for the discrete temperatures spaces $\Sigma_{0,h}$, $\Sigma_{D,h}$ and the velocities space $\mathbf{V}_h$ that can be found in [17, 31] and [16] respectively.

Lemma 2 (Approximation property of $\Sigma_{0,h}$, $\Sigma_{D,h}$) *Under the assumption (A1) for any $\boldsymbol{\tau} \in \Sigma_0 \cap H^{s+1}(\Omega_h)$ (resp. $\Sigma_D \cap H^{s+1}(\Omega_h)$) there exists $\boldsymbol{\tau}_I \in \Sigma_{0,h}$ (resp. $\Sigma_{D,h}$) such that for all $E \in \Omega_h$ it holds*

$$\|\boldsymbol{\tau} - \boldsymbol{\tau}_I\|_{0,E} + h_E |\boldsymbol{\tau} - \boldsymbol{\tau}_I|_{1,E} \lesssim h_E^{s+1} |\boldsymbol{\tau}|_{s+1,E}$$

where $0 < s \leq k$.

Lemma 3 (Approximation property of $\mathbf{V}_h$) *Under the assumption (A1) for any $\mathbf{v} \in \mathbf{V} \cap [H^{s+1}(\Omega_h)]^2$ there exists $\mathbf{v}_I \in \mathbf{V}_h$ such that for all $E \in \Omega_h$ it holds*

$$\|\mathbf{v} - \mathbf{v}_I\|_{0,E} + h_E |\mathbf{v} - \mathbf{v}_I|_{1,E} \lesssim h_E^{s+1} |\mathbf{v}|_{s+1,E}$$

where $0 < s \leq k$. Furthermore if $\mathbf{v} \in \mathbf{Z}$ it holds $\mathbf{v}_I \in \mathbf{Z}_h$.

In order to derive the error estimates, we state the following assumption on the solutions and on the data of Problem (2.1).

(A2) Regularity assumptions:

- the solution $(\mathbf{u}, p, \vartheta)$ satisfy $\mathbf{u} \in [H^{s+1}(\Omega_h)]^2$, $p \in [H^s(\Omega_h)]$, $\vartheta \in [H^{s+1}(\Omega_h)]$,
- the viscosity ν , the conductivity κ satisfy $\nu(\vartheta) \in W_\infty^s(\Omega_h)$, $\kappa \in W_\infty^s(\Omega_h)$,
- the external loads $\mathbf{f}$ and g satisfy $\mathbf{f} \in [H^{s+1}(\Omega_h)]^2$, $g \in [H^{s+1}(\Omega_h)]$.

for some $0 < s \leq k$.

In the following $\mathcal{R}(\cdot)$ will denote a generic constant independent of h , but depending on the norms of $\mathbf{u}$, p , ϑ , ν , κ , $\mathbf{f}$ and g in the aforementioned spaces, on Ω , on the “polynomial” order k of the method and on the regularity constant appearing in the mesh assumption **(A1)**.

Proposition 5 *Under the assumptions **(A0)**, **(A1)** and under the assumptions of Theorem 2 and Proposition 4, let $(\mathbf{u}, p, \vartheta)$ be the solution of (2.7) and $(\mathbf{u}_h, p_h, \vartheta_h)$ be the solution of (3.27). Assume that ν_* and κ_* in assumptions **(A0)** are such that*

$$\alpha_* \nu_* - \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} - \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}}} \frac{C_{\text{lip}}}{C_{\text{sol}} \beta_* \kappa_*} > 0. \quad (5.1)$$

Assuming moreover **(A2)** the following error estimates hold:

$$\begin{aligned} \|\mathbf{u}_h - \mathbf{u}\|_{1,\Omega} + \|\vartheta - \vartheta_h\|_{1,\Omega} &\leq \mathcal{R}(\mathbf{u}, \vartheta, \nu, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}, \\ \|p - p_h\|_{0,\Omega} &\leq \mathcal{R}(\mathbf{u}, p, \vartheta, \nu, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}. \end{aligned}$$

Proof. We preliminary observe that under the assumptions of Theorem 2 (resp. Proposition 4) and recalling (2.13) (resp. (4.2)) the following hold

$$\|\mathbf{u}\|_{1,\Omega}^2 + \|\vartheta\|_{1,\Omega}^2 \leq \frac{C_{\text{est}}^2}{C_{\text{data}}^2}, \quad \|\mathbf{u}_h\|_{1,\Omega}^2 + \|\vartheta_h\|_{1,\Omega}^2 \leq \frac{\widehat{C}_{\text{est}}^2}{\widehat{C}_{\text{data}}^2}. \quad (5.2)$$

The proof follows the following steps.

Step 1. Interpolation error estimates.

Let us introduce the following error quantities

$$\begin{aligned} \mathbf{e}_I &:= \mathbf{u} - \mathbf{u}_I, & \sigma_I &:= \vartheta - \vartheta_I, & \rho_I &:= p - p_I, \\ \mathbf{e}_h &:= \mathbf{u}_I - \mathbf{u}_h, & \sigma_h &:= \vartheta_I - \vartheta_h, & \rho_h &:= p_I - p_h, \end{aligned}$$

where ϑ_I and $\mathbf{u}_I$ are the interpolant function of ϑ and $\mathbf{u}$ of Lemma 2 and Lemma 3 respectively and $p_I \in Q_h$ is defined elementwise by $p_I|_E = \Pi_{k-1}^{0,E} p$ for any $E \in \Omega_h$ (cf. (3.2)). Therefore from Lemma 3, Lemma 2 and Bramble-Hilbert Lemma (3.6), we have

$$\|\mathbf{e}_I\|_{1,\Omega} \leq \mathcal{R}(\mathbf{u}) h^s, \quad \|\sigma_I\|_{1,\Omega} \leq \mathcal{R}(\vartheta) h^s, \quad \|\rho_I\|_{1,\Omega} \leq \mathcal{R}(p) h^s. \quad (5.3)$$

Notice now that since $\mathbf{u} \in \mathbf{Z}$, from Lemma 3 it holds that $\mathbf{e}_h \in \mathbf{Z}_h$. Whereas from Lemma 2, $\sigma_h \in \Sigma_{0,h}$.

Step 2. Error equation for the temperature.

Employing Property **(P7)** and manipulating the second equation in (2.8) and (3.28) we obtain

$$\begin{aligned}
\beta_* \kappa_* |\sigma_h|_{1,\Omega}^2 &\leq a_{T,h}(\sigma_h, \sigma_h) = a_{T,h}(\vartheta_I, \sigma_h) - a_{T,h}(\vartheta_h, \sigma_h) \\
&= (a_{T,h}(\vartheta_I, \sigma_h) - a_T(\vartheta, \sigma_h)) + (a_T(\vartheta, \sigma_h) - a_{T,h}(\vartheta_h, \sigma_h)) \\
&= (a_{T,h}(\vartheta_I, \sigma_h) - a_T(\vartheta, \sigma_h)) + \left(c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h, \sigma_h) - c_T^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) \right) + (g - g_h, \sigma_h) \\
&=: \eta_{a,T} + \eta_{c,T} + \eta_g.
\end{aligned} \tag{5.4}$$

We now estimate the three terms above.

• estimate of $\eta_{a,T}$: we split $\eta_{a,T}$ into local contributions $\eta_{a,T}^E$. Recalling definitions (2.3) and (3.18) and the property of the L^2 -projection we have

$$\begin{aligned}
\eta_{a,T}^E &:= \int_E \kappa \Pi_{k-1}^{0,E} \nabla \vartheta_I \cdot \Pi_{k-1}^{0,E} \nabla \sigma_h \, dE + s_T^E(\vartheta_I, \sigma_h) - \int_E \kappa \nabla \vartheta \cdot \nabla \sigma_h \, dE \\
&= \int_E \kappa (\Pi_{k-1}^{0,E} \nabla \vartheta_I - \nabla \vartheta) \cdot \Pi_{k-1}^{0,E} \nabla \sigma_h \, dE - \int_E \kappa \nabla \vartheta \cdot (I - \Pi_{k-1}^{0,E}) \nabla \sigma_h \, dE + s_T^E(\vartheta_I, \sigma_h) \\
&= \int_E \kappa (\Pi_{k-1}^{0,E} \nabla \vartheta_I - \nabla \vartheta) \cdot \Pi_{k-1}^{0,E} \nabla \sigma_h \, dE - \int_E (I - \Pi_{k-1}^{0,E}) (\kappa \nabla \vartheta) \cdot \nabla \sigma_h \, dE + s_T^E(\vartheta_I, \sigma_h).
\end{aligned}$$

Then from Property **(P0)**, (3.20) and (3.21), employing Bramble-Hilbert Lemma (3.6), we obtain

$$\begin{aligned}
\eta_{a,T}^E &\leq \left(\kappa^* \|\Pi_{k-1}^{0,E} \nabla \vartheta_I - \nabla \vartheta\|_{0,E} + \|(I - \Pi_{k-1}^{0,E})(\kappa \nabla \vartheta)\|_{0,E} + \kappa^* |(I - \Pi_k^{0,E}) \vartheta_I|_{1,E} \right) |\sigma_h|_{1,E} \\
&\leq \kappa^* \left(|\sigma_I|_{1,E} + |(I - \Pi_k^{0,E}) \vartheta|_{1,E} + \frac{1}{\kappa^*} \|(I - \Pi_{k-1}^{0,E})(\kappa \nabla \vartheta)\|_{0,E} + \right) |\sigma_h|_{1,E} \\
&\lesssim \kappa^* h_E^s \left(|\sigma_I|_{1,E} + |\vartheta|_{s+1,E} + \frac{\|\kappa\|_{W_\infty^s(E)}}{\kappa^*} \|\vartheta\|_{s+1,E} \right) |\sigma_h|_{1,E}.
\end{aligned}$$

Therefore summing the local contributions and employing (5.3) we obtain

$$\eta_{a,T} \leq \mathcal{R}(\vartheta, \kappa) h^s |\sigma_h|_{1,\Omega}. \tag{5.5}$$

• estimate of $\eta_{c,T}$: employing the skew-symmetry of $c_{T,h}^{\text{skew}}(\cdot; \cdot, \cdot)$ (cf. (3.25)) we have

$$\begin{aligned}
\eta_{c,T} &= c_{T,h}^{\text{skew}}(\mathbf{u}_h; \vartheta_h, \sigma_h) - c_{T,h}^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) - c_T^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) \\
&= -c_{T,h}^{\text{skew}}(\mathbf{u}_h; \sigma_I, \sigma_h) - c_{T,h}^{\text{skew}}(\mathbf{u} - \mathbf{u}_h; \vartheta, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) - c_T^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h).
\end{aligned} \tag{5.6}$$

From Property **(P7)**, bounds (5.2) and (5.3) we have

$$\begin{aligned}
c_{T,h}^{\text{skew}}(\mathbf{u}_h; \sigma_I, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u} - \mathbf{u}_h; \vartheta, \sigma_h) &\leq \widehat{C}_{\text{conv}} (|\mathbf{u}_h|_{1,\Omega} \|\sigma_I\|_{1,\Omega} + |\mathbf{u} - \mathbf{u}_h|_{1,\Omega} \|\vartheta\|_{1,\Omega}) |\sigma_h|_{1,\Omega} \\
&\leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} \|\sigma_I\|_{1,\Omega} |\sigma_h|_{1,\Omega} + \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}}} (|\mathbf{e}_I|_{1,\Omega} + |\mathbf{e}_h|_{1,\Omega}) |\sigma_h|_{1,\Omega} \\
&\leq \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}}} |\mathbf{e}_h|_{1,\Omega} |\sigma_h|_{1,\Omega} + \mathcal{R}(\mathbf{u}, \vartheta) h^s |\sigma_h|_{1,\Omega}.
\end{aligned} \tag{5.7}$$

Furthermore employing [16, Lemma 4.3]

$$c_{T,h}^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) - c_T^{\text{skew}}(\mathbf{u}; \vartheta, \sigma_h) \leq \mathcal{R}(\mathbf{u}, \vartheta) h^s |\sigma_h|_{1,\Omega}. \quad (5.8)$$

Therefore collecting in (5.6), bounds (5.7) and (5.8) we obtain

$$\eta_{c,T} \leq \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}}} |\mathbf{e}_h|_{1,\Omega} |\sigma_h|_{1,\Omega} + \mathcal{R}(\mathbf{u}, \vartheta) h^s |\sigma_h|_{1,\Omega}. \quad (5.9)$$

• estimate of η_g : employing the definition of L^2 -projection (3.2) and Bramble-Hilbert Lemma (3.6) we infer

$$\eta_g \leq \mathcal{R}(g) h^{s+2} |\sigma_h|_{1,\Omega}. \quad (5.10)$$

Finally combining (5.5), (5.9) and (5.10) in (5.4) we get

$$|\sigma_h|_{1,\Omega} \leq \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}} \beta_* \kappa_*} |\mathbf{e}_h|_{1,\Omega} + \mathcal{R}(\mathbf{u}, \vartheta, \kappa) h^s + \mathcal{R}(g) h^{s+2}. \quad (5.11)$$

Step 3. Error equation for the velocity.

We now analyse the velocity error equation. Combining the first equation in (2.8) and (3.28) together with Property (P7), we have

$$\begin{aligned} \alpha_* \mathbf{v}_* |\mathbf{e}_h|_{1,\Omega}^2 &\leq a_{V,h}(\vartheta_h; \mathbf{e}_h, \mathbf{e}_h) = a_{V,h}(\vartheta_h; \mathbf{u}_I, \mathbf{e}_h) - a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{e}_h) \\ &= (a_{V,h}(\vartheta_h; \mathbf{u}_I, \mathbf{e}_h) - a_V(\vartheta; \mathbf{u}, \mathbf{e}_h)) + (a_V(\vartheta; \mathbf{u}, \mathbf{e}_h) - a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{e}_h)) \\ &= (a_{V,h}(\vartheta_h; \mathbf{u}_I, \mathbf{e}_h) - a_V(\vartheta; \mathbf{u}, \mathbf{e}_h)) + (c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{e}_h) - c_V^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{e}_h)) + (\mathbf{f} - \mathbf{f}_h, \mathbf{e}_h) \\ &=: \eta_{a,V} + \eta_{c,V} + \eta_f. \end{aligned} \quad (5.12)$$

We now estimate each term above.

• estimate of $\eta_{a,V}$: we split the term into local contributions $\eta_{a,V}^E$. Recalling definitions (2.2) and (3.17) we have

$$\begin{aligned} \eta_{a,V}^E &:= \int_E \mathbf{v}(\Pi_k^{0,E} \vartheta_h) \Pi_{k-1}^{0,E} \nabla \mathbf{u}_I \cdot \Pi_{k-1}^{0,E} \nabla \mathbf{e}_h \, dE + s_V^E(\vartheta_h; \mathbf{u}_I, \mathbf{e}_h) - \int_E \mathbf{v}(\vartheta) \nabla \mathbf{u} \cdot \nabla \mathbf{e}_h \, dE \\ &= \int_E \mathbf{v}(\Pi_k^{0,E} \vartheta_h) (\Pi_{k-1}^{0,E} \nabla \mathbf{u}_I - \nabla \mathbf{u}) \cdot \Pi_{k-1}^{0,E} \nabla \mathbf{e}_h \, dE + s_V^E(\vartheta_h; \mathbf{u}_I, \mathbf{e}_h) + \\ &\quad - \int_E \mathbf{v}(\vartheta) \nabla \mathbf{u} \cdot (I - \Pi_{k-1}^{0,E}) \nabla \mathbf{e}_h \, dE + \int_E (\mathbf{v}(\Pi_k^{0,E} \vartheta_h) - \mathbf{v}(\vartheta)) \nabla \mathbf{u} \cdot \Pi_{k-1}^{0,E} \nabla \mathbf{e}_h \, dE \\ &=: \zeta_1^E + \zeta_2^E + \zeta_3^E + \zeta_4^E. \end{aligned} \quad (5.13)$$

Recalling Property (P0), the stability bounds (3.19) and (3.21), and the definition of L^2 -projection (3.2), the first three sub-terms can be bounded as follows

$$\begin{aligned} \sum_{i=1}^3 \zeta_i^E &\leq (\mathbf{v}^* \|\Pi_{k-1}^{0,E} \nabla \mathbf{u}_I - \nabla \mathbf{u}\|_{0,E} + \|(I - \Pi_{k-1}^{0,E})(\mathbf{v}(\vartheta) \nabla \mathbf{u})\|_{0,E} + \mathbf{v}^* \|(I - \Pi_k^{0,E}) \mathbf{u}_I\|_{1,E}) |\mathbf{e}_h|_{1,E} \\ &\leq \mathbf{v}^* \left(|\mathbf{e}_I|_{1,E} + \|(I - \Pi_k^{0,E}) \mathbf{u}\|_{1,E} + \frac{1}{\mathbf{v}^*} \|(I - \Pi_k^{0,E})(\mathbf{v}(\vartheta) \nabla \mathbf{u})\|_{0,E} \right) |\mathbf{e}_h|_{1,E} \\ &\lesssim \mathbf{v}^* h_E^s \left(|\mathbf{e}_I|_{1,E} + \|\mathbf{u}\|_{s+1,E} + \frac{\|\mathbf{v}(\vartheta)\|_{W_\infty^s(E)}}{\mathbf{v}^*} \|\mathbf{u}\|_{s+1,E} \right) |\mathbf{e}_h|_{1,E}. \end{aligned}$$

Therefore

$$\sum_{E \in \Omega_h} \sum_{i=1}^3 \zeta_i^E \leq \mathcal{R}(\mathbf{u}, \mathbf{v}) h^s |\mathbf{e}_h|_{1,\Omega}. \quad (5.14)$$

For the term ζ_4 , employing similar computations to that in (4.14) (recalling that p is such that $1/p + 1/q = 1/2$) we infer

$$\begin{aligned} \zeta_4^E &\leq \nu_{\text{lip}} \|\Pi_k^{0,E} \vartheta_h - \vartheta\|_{L^p(E)} |\mathbf{u}|_{[W_q^1(E)]^2} |\mathbf{e}_h|_{1,E} \\ &\leq \nu_{\text{lip}} \left(\|\Pi_k^{0,E} (\vartheta_h - \vartheta)\|_{L^p(E)} + \|(I - \Pi_k^{0,E}) \vartheta\|_{L^p(E)} \right) |\mathbf{u}|_{[W_q^1(E)]^2} |\mathbf{e}_h|_{1,E} \\ &\lesssim \nu_{\text{lip}} \left(\|\vartheta_h - \vartheta\|_{L^p(E)} + \|\vartheta\|_{W_{\tilde{p}}^s(E)} h_E^s \right) |\mathbf{u}|_{[W_q^1(E)]^2} |\mathbf{e}_h|_{1,E} \\ &\lesssim \nu_{\text{lip}} \left(\|\sigma_h\|_{L^p(E)} + \|\sigma_I\|_{L^p(E)} + \|\vartheta\|_{W_{\tilde{p}}^s(E)} h_E^s \right) |\mathbf{u}|_{[W_q^1(E)]^2} |\mathbf{e}_h|_{1,E} \end{aligned}$$

where we also used the continuity of the L^2 -projection w.r.t. the L^p -norm and Bramble-Hilbert (3.6). From the Hölder inequality for sequences and the Sobolev embedding (3.4), we infer

$$\sum_{E \in \Omega_h} \zeta_4^E \lesssim \nu_{\text{lip}} (|\sigma_h|_{1,\Omega} + \|\sigma_I\|_{1,\Omega} + \mathcal{R}(\vartheta) h^s) |\mathbf{u}|_{[W_q^1(\Omega_h)]^2} |\mathbf{e}_h|_{1,\Omega}.$$

Therefore from (5.3) and (2.15)

$$\sum_{E \in \Omega_h} \zeta_4^E \leq \frac{C_{\text{lip}}}{C_{\text{sol}}} |\sigma_h|_{1,\Omega} |\mathbf{e}_h|_{1,\Omega} + \mathcal{R}(\vartheta) h^s |\mathbf{e}_h|_{1,\Omega} \quad (5.15)$$

for a suitable positive constant C_{lip} . Collecting (5.13) and (5.14) in (5.15) we obtain

$$\eta_{a,V} \leq \frac{C_{\text{lip}}}{C_{\text{sol}}} |\sigma_h|_{1,\Omega} |\mathbf{e}_h|_{1,\Omega} + \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}) h^s |\mathbf{e}_h|_{1,\Omega}. \quad (5.16)$$

- estimate of $\eta_{c,V}$: direct application of [16, Lemma 4.3 and Lemma 4.4] yields

$$\eta_{c,V} \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} |\mathbf{e}_h|_{1,\Omega}^2 + \mathcal{R}(\mathbf{u}) h^s |\mathbf{e}_h|_{1,\Omega}. \quad (5.17)$$

- estimate of η_f : using the same argument used above for η_g we obtain

$$\eta_f \leq \mathcal{R}(\mathbf{f}) h^{s+2} |\mathbf{e}_h|_{1,\Omega}. \quad (5.18)$$

Collecting (5.16), (5.17) and (5.18) in the error equation (5.12) we obtain

$$\alpha_* \nu_* |\mathbf{e}_I|_{1,\Omega} \leq \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}) h^s + \mathcal{R}(\mathbf{f}) h^{s+2} + \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} |\mathbf{e}_h|_{1,\Omega} + \frac{C_{\text{lip}}}{C_{\text{sol}}} |\sigma_h|_{1,\Omega}. \quad (5.19)$$

Step 4. Error estimates for the velocity and the temperature

Now employing the estimate (5.11) in (5.19) we infer

$$\left(\alpha_* \mathbf{v}_* - \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} - \frac{\widehat{C}_{\text{conv}} C_{\text{est}}}{C_{\text{data}}} \frac{C_{\text{lip}}}{C_{\text{sol}} \beta_* \kappa_*} \right) |\mathbf{e}_I|_{1,\Omega} \leq \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}.$$

Therefore from (5.1) we get

$$|\mathbf{e}_I|_{1,\Omega} \leq \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}. \quad (5.20)$$

We finally collect the estimates (5.20), (5.11) and the error bounds (5.3) obtaining

$$\|\mathbf{u} - \mathbf{u}_h\|_{1,\Omega} + \|\vartheta - \vartheta_h\|_{1,\Omega} \leq \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}. \quad (5.21)$$

Step 5. Error estimate for the pressure.

In the last step we briefly study the convergence error for the pressures. From the first equation in (2.7) and (3.27), simple computations yield

$$\begin{aligned} b(\mathbf{v}_h, \rho_h) &= b(\mathbf{v}_h, p_I) - b(\mathbf{v}_h, p_h) = b(\mathbf{v}_h, p) - b(\mathbf{v}_h, p_h) - b(\mathbf{v}_h, p_I) \\ &= (a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) - a_V(\vartheta; \mathbf{u}, \mathbf{v}_h)) + (c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) - c_V^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h)) + \\ &\quad + (\mathbf{f} - \mathbf{f}_h, \mathbf{v}_h) - b(\mathbf{v}_h, \rho_I) \\ &:= \gamma_{a,V} + \gamma_{c,V} + \gamma_{\mathbf{f}} + \gamma_b. \end{aligned} \quad (5.22)$$

- estimate of $\gamma_{a,V}$: the estimate of $\gamma_{a,V}$ follows the same steps for the bound of $\eta_{a,V}$

$$\gamma_{a,V} \leq \mathbf{v}^* \|\mathbf{u} - \mathbf{u}_h\|_{1,\Omega} \|\mathbf{v}_h\|_{1,\Omega} + \frac{C_{\text{lip}}}{C_{\text{sol}}} \|\vartheta - \vartheta_h\|_{1,\Omega} \|\mathbf{v}_h\|_{1,\Omega} + \mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}) h^s \|\mathbf{v}_h\|_{1,\Omega}$$

therefore from (5.21) we have

$$\gamma_{a,V} \leq (\mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}) \|\mathbf{v}_h\|_{1,\Omega}. \quad (5.23)$$

- estimate of $\gamma_{c,V}$: simple computations yield

$$\begin{aligned} \gamma_{c,V} &= (c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) - c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h)) + (c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h) - c_V^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h)) \\ &= (c_{V,h}^{\text{skew}}(\mathbf{u}_h - \mathbf{u}; \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}_h - \mathbf{u}, \mathbf{v}_h)) + (c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h) - c_V^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h)). \end{aligned} \quad (5.24)$$

Employing Property (P7) and bounds (5.2) and estimate (5.21), we have

$$\begin{aligned} c_{V,h}^{\text{skew}}(\mathbf{u}_h - \mathbf{u}; \mathbf{u}_h, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}_h - \mathbf{u}, \mathbf{v}_h) &\leq \widehat{C}_{\text{conv}} (\|\mathbf{u}\|_{1,\Omega} + \|\mathbf{u}_h\|_{1,\Omega}) \|\mathbf{u} - \mathbf{u}_h\|_{1,\Omega} \|\mathbf{v}_h\|_{1,\Omega} \\ &\leq \widehat{C}_{\text{conv}} \left(\frac{C_{\text{est}}}{C_{\text{data}}} + \frac{\widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} \right) \|\mathbf{u} - \mathbf{u}_h\|_{1,\Omega} \|\mathbf{v}_h\|_{1,\Omega} \\ &\leq (\mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}) \|\mathbf{v}_h\|_{1,\Omega}. \end{aligned} \quad (5.25)$$

Whereas applying [16, Lemma 4.3] we have

$$c_{V,h}^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h) - c_V^{\text{skew}}(\mathbf{u}; \mathbf{u}, \mathbf{v}_h) \leq \mathcal{R}(\mathbf{u}) h^s |\mathbf{v}_h|_{1,\Omega}. \quad (5.26)$$

Therefore (5.24), (5.25) and (5.26) implies

$$\gamma_{c,v} \leq (\mathcal{R}(\mathbf{u}, \vartheta, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}) |\mathbf{v}_h|_{1,\Omega}. \quad (5.27)$$

• estimate of $\gamma_f + \gamma_b$: the term γ_f can be bounded as η_f . The term γ_b can be bounded, recalling Property (P0), by Bramble-Hilbert Lemma 3.6. Then we have

$$\gamma_f + \gamma_b \leq (\mathcal{R}(\mathbf{f}) h^{s+2} + \mathcal{R}(p) h^s) |\mathbf{v}_h|_{1,\Omega}. \quad (5.28)$$

Employing the inf-sup stability (3.10), from (5.23)–(5.28), we obtain

$$\hat{\beta} \|\rho_h\|_{0,E} \leq \sup_{\mathbf{v}_h \in \mathbf{V}_h} \frac{b(\mathbf{v}_h, \rho_h)}{|\mathbf{v}_h|_{1,\Omega}} \leq \mathcal{R}(\mathbf{u}, \vartheta, p, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}.$$

The pressure error estimates follows from the previous bound and (5.3):

$$\|p - p_h\|_{0,E} \leq \mathcal{R}(\mathbf{u}, \vartheta, p, \mathbf{v}, \kappa) h^s + \mathcal{R}(\mathbf{f}, g) h^{s+2}.$$

□

Remark 4 *Note that the proposed VEM scheme (3.28) has the following favorable property that extends to the context of coupled problems the convergence result obtained for the Navier–Stokes equation: the error components partly decouple. In fact the velocity and the temperature errors do not depend directly on the discrete pressures, but only indirectly through the approximation of the loading and convection terms and such dependence on the full load is much weaker with respect to standard mixed schemes. In some situations the partial decoupling of the errors induces a positive effect on the velocity/temperature approximation (see for instance Test 2).*

6. Numerical results

In this section we present two numerical experiments to test the practical performance of the proposed VEM scheme (3.27) and the possible advantages related to the divergence-free velocity solutions. In the numerical tests we set $k = 2$.

6.1. Fixed point iteration

We describe the linearization strategy based on a fixed-point iteration adopted to solve the non-linear coupled problem.

LINEAR FIXED POINT ITERATION

Starting from $(\mathbf{u}_h^0, p_h^0) = (\mathbf{0}, 0)$, for $n \geq 0$, until convergence solve

- HEAT equation

$$\begin{cases} \text{find } \vartheta_h^{n+1} \in \Sigma_{D,h}, \text{ such that} \\ a_{T,h}(\vartheta_h^{n+1}, \sigma_h) + c_{T,h}^{\text{skew}}(\mathbf{u}_h^n; \vartheta_h^{n+1}, \sigma_h) = (g_h, \sigma_h) \end{cases} \quad \forall \sigma_h \in \Sigma_{0,h},$$

- `OSEEN` equation

$$\begin{cases} \text{find } (\mathbf{u}_h^{n+1}, p_h^{n+1}) \in \mathbf{V}_h \times Q_h, \text{ such that} \\ a_{V,h}(\vartheta_h^{n+1}; \mathbf{u}_h^{n+1}, \mathbf{v}_h) + c_{V,h}^{\text{skew}}(\mathbf{u}_h^n, \mathbf{u}_h^{n+1}, \mathbf{v}_h) + b(\mathbf{v}_h, p_h^{n+1}) = (\mathbf{f}_h, \mathbf{v}_h) & \forall \mathbf{v}_h \in \mathbf{V}_h, \\ b(\mathbf{u}_h^{n+1}, q_h) = 0 & \forall q_h \in Q_h. \end{cases}$$

Notice that at each step of the fixed-point iteration the solution $\mathbf{u}_h^n$ is still divergence-free, therefore the linearization procedure does not affect the divergence-free property of the final discrete solution. At each iteration we need to solve two linear systems with dimension

$$\begin{aligned} \text{N_DOFs}_{\text{HEAT}} &= N_V + (k-1)N_E + \frac{k(k-1)}{2}N_P, \\ \text{N_DOFs}_{\text{OSEEN}} &= 2N_V + 2(k-1)N_E + \left(\frac{k(3k-1)}{2}\right)N_P + 1, \end{aligned}$$

where N_V , N_E , N_P denotes the number of internal vertices, internal edges and polygons in Ω_h respectively. The dimension of the linear system for the `OSEEN` equation can be significantly reduced considering the so-called reduced spaces (see [16]). We now derive the convergence of the fixed-point iteration.

Proposition 6 *Under the assumptions (A0), (A1) and under the assumptions of Proposition 4, let $(\mathbf{u}_h, p_h, \vartheta_h)$ be the solution of (3.27). Then the sequence $\{(\mathbf{u}_h^n, p_h^n, \vartheta_h^n)\}_n$ generated by the `LINEAR FIXED POINT ITERATION` satisfies*

$$\lim_{n \rightarrow \infty} \|\mathbf{u}_h - \mathbf{u}_h^n\|_{1,\Omega} = 0 \quad \lim_{n \rightarrow \infty} \|\vartheta_h - \vartheta_h^n\|_{1,\Omega} = 0 \quad \lim_{n \rightarrow \infty} \|p_h - p_h^n\|_{0,\Omega} = 0.$$

Proof. We only sketch the proof since it can be derived using similar technique to that in the proof of Proposition 4. For any $n \geq 0$, let us introduce the error quantities

$$\boldsymbol{\delta}_h^n := \mathbf{u}_h^n - \mathbf{u}_h, \quad \sigma_h^n := \vartheta_h^n - \vartheta_h, \quad \rho_h^n := p_h^n - p_h.$$

Note that $\boldsymbol{\delta}_h^n \in \mathbf{Z}_h$ and $\sigma_h^n \in \Sigma_{0,h}$. Simple computations yield

$$a_{T,h}(\sigma_h^{n+1}, \sigma_h^{n+1}) = -c_{T,h}^{\text{skew}}(\boldsymbol{\delta}_h^n, \vartheta_h, \sigma_h^{n+1}).$$

Therefore from Property (P7) and bound (5.2) we obtain.

$$\beta_* \kappa_* \|\sigma_h^{n+1}\|_{1,\Omega} \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} \|\boldsymbol{\delta}_h^n\|_{1,\Omega}. \quad (6.1)$$

For the velocity we can derive the following equation

$$a_{V,h}(\vartheta_h^{n+1}; \boldsymbol{\delta}_h^{n+1}, \boldsymbol{\delta}_h^{n+1}) = -c_{V,h}^{\text{skew}}(\boldsymbol{\delta}_h^n, \mathbf{u}_h, \boldsymbol{\delta}_h^{n+1}) - \left(a_{V,h}(\vartheta_h^{n+1}; \mathbf{u}_h, \boldsymbol{\delta}_h^{n+1}) - a_{V,h}(\vartheta_h; \mathbf{u}_h, \boldsymbol{\delta}_h^{n+1}) \right).$$

Then, employing properties **(P7)** and Lemma 1 we infer

$$\alpha_* \mathbf{v}_* |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega}^2 \leq \widehat{C}_{\text{conv}} |\mathbf{u}_h|_{1,\Omega} |\boldsymbol{\delta}_h^n|_{1,\Omega} |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} + \widehat{C}_{\text{lip}} \mathbf{v}_{\text{lip}} |\mathbf{u}_h|_{[W_q^1(\Omega_h)]^2} |\boldsymbol{\sigma}_h^{n+1}|_{1,\Omega} |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega}.$$

Using (5.2) and (4.16), from (6.1) we derive

$$\alpha_* \mathbf{v}_* |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} |\boldsymbol{\delta}_h^n|_{1,\Omega} + \frac{\widehat{C}_{\text{lip}}}{\widehat{C}_{\text{sol}}} \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}} \beta_* \kappa_*} |\boldsymbol{\delta}_h^n|_{1,\Omega},$$

that, recalling (4.17), corresponds to

$$|\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} \leq K |\boldsymbol{\delta}_h^n|_{1,\Omega}.$$

Therefore, since $\mathbf{u}_h^0 = \mathbf{0}$ from (5.2), we have

$$|\boldsymbol{\delta}_h^n|_{1,\Omega} \leq \frac{\widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} K^n, \quad |\boldsymbol{\sigma}_h^n|_{1,\Omega} \leq \frac{\widehat{C}_{\text{conv}} \widehat{C}_{\text{est}}^2}{\widehat{C}_{\text{data}}^2 \beta_* \kappa_*} K^{n-1}. \quad (6.2)$$

Concerning the pressure we derive

$$\begin{aligned} b(\mathbf{v}_h, \rho_h^{n+1}) &= a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) - a_{V,h}(\vartheta_h^{n+1}; \mathbf{u}_h^{n+1}, \mathbf{v}_h) + \\ &\quad + c_{V,h}^{\text{skew}}(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) - c_{V,h}^{\text{skew}}(\mathbf{u}_h^n, \mathbf{u}_h^{n+1}, \mathbf{v}_h) \\ &= (a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) - a_{V,h}(\vartheta_h^{n+1}; \mathbf{u}_h, \mathbf{v}_h)) - a_{V,h}(\vartheta_h^{n+1}; \boldsymbol{\delta}_h^{n+1}, \mathbf{v}_h) \\ &\quad - c_{V,h}^{\text{skew}}(\boldsymbol{\delta}_h^n; \mathbf{u}_h, \mathbf{v}_h) - c_{V,h}^{\text{skew}}(\mathbf{u}_h^n; \boldsymbol{\delta}_h^{n+1}, \mathbf{v}_h). \end{aligned} \quad (6.3)$$

Lemma 1 implies

$$a_{V,h}(\vartheta_h; \mathbf{u}_h, \mathbf{v}_h) - a_{V,h}(\vartheta_h^{n+1}; \mathbf{u}_h, \mathbf{v}_h) \leq \widehat{C}_{\text{lip}} \mathbf{v}_{\text{lip}} |\boldsymbol{\sigma}_h^{n+1}|_{1,\Omega} |\mathbf{u}_h|_{[W_q^1(\Omega)]^2} |\mathbf{v}_h|_{1,\Omega}, \quad (6.4)$$

whereas Property **(P7)** yields

$$\begin{aligned} -a_{V,h}(\vartheta_h^{n+1}; \boldsymbol{\delta}_h^{n+1}, \mathbf{v}_h) &\leq \alpha^* \mathbf{v}^* |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} |\mathbf{v}_h|_{1,\Omega} \\ -c_{V,h}^{\text{skew}}(\boldsymbol{\delta}_h^n; \mathbf{u}_h, \mathbf{v}_h) - c_{V,h}^{\text{skew}}(\mathbf{u}_h^n; \boldsymbol{\delta}_h^{n+1}, \mathbf{v}_h) &\leq \widehat{C}_{\text{conv}} \left(|\boldsymbol{\delta}_h^n|_{1,\Omega} |\mathbf{u}_h|_{1,\Omega} + |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} |\mathbf{u}_h^n|_{1,\Omega} \right) |\mathbf{v}_h|_{1,\Omega}. \end{aligned} \quad (6.5)$$

Then recalling (4.16) and (5.2) from (6.3), (6.4) and (6.5), for any $\mathbf{v}_h \in \mathbf{V}_h$ we have

$$\frac{b(\mathbf{v}_h, \rho_h^{n+1})}{|\mathbf{v}_h|_{1,\Omega}} \leq \frac{\widehat{C}_{\text{lip}}}{\widehat{C}_{\text{sol}}} |\boldsymbol{\sigma}_h^{n+1}|_{1,\Omega} + \left(\alpha^* \mathbf{v}^* + \widehat{C}_{\text{conv}} \frac{\widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} \right) |\boldsymbol{\delta}_h^{n+1}|_{1,\Omega} + \widehat{C}_{\text{conv}} \frac{\widehat{C}_{\text{est}}}{\widehat{C}_{\text{data}}} |\boldsymbol{\delta}_h^n|_{1,\Omega}.$$

Therefore from the inf-sup stability (3.10) and (6.2) we have

$$\widehat{\beta} \|\rho_h^{n+1}\|_{0,\Omega} \leq \widehat{C}_{\text{conv}} \frac{\widehat{C}_{\text{est}}^2}{\widehat{C}_{\text{data}}^2} \left(1 + \frac{\widehat{C}_{\text{lip}}}{\widehat{C}_{\text{sol}} \beta_* \kappa_*} \right) K^n + O(K^{n+1}).$$

Now the thesis follows from (4.17). $\square$

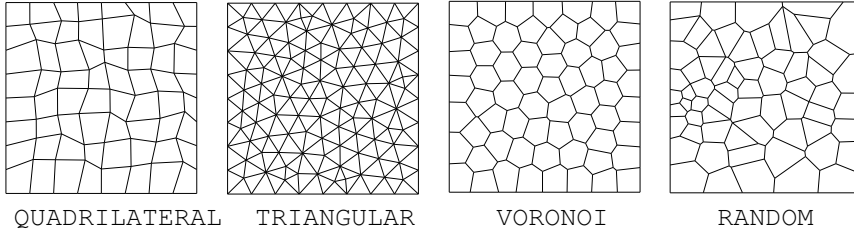


Figure 1. Test 1. Example of the adopted polygonal meshes.

Test 1. Error convergence In this test we examine the practical performance and the convergence properties of the proposed scheme (3.27) in the light of Proposition 5.

In order to compute the VEM error between the exact solution $\mathbf{u}_{\text{ex}}$ and the VEM solution $\mathbf{u}_h$ we consider the computable error quantities

$$\begin{aligned} \text{err}(\mathbf{u}_h, H^1) &:= \text{sqrt} \left(\sum_{E \in \Omega_h} \|\nabla \mathbf{u}_{\text{ex}} - \Pi_1^{0,E} \nabla \mathbf{u}_h\|_{0,E}^2 \right) / \|\nabla \mathbf{u}_{\text{ex}}\|_0, \\ \text{err}(\mathbf{u}_h, L^2) &:= \text{sqrt} \left(\sum_{E \in \Omega_h} \|\mathbf{u}_{\text{ex}} - \Pi_2^{0,E} \mathbf{u}_h\|_{0,E} \right) / \|\mathbf{u}_{\text{ex}}\|_0. \end{aligned}$$

Analogous error quantities are considered to measure the error between the exact solution ϑ_{ex} and the VEM solution ϑ_h . For the pressure error we take $\text{err}(p_h, L^2) := \|p - p_h\|_{0,\Omega}$.

In the present test we consider Problem (2.1) on the unit square $\Omega = (0, 1)^2$, the viscosity is $\nu(\vartheta) = 2 + \vartheta^2$, the conductivity is $\kappa(x, y) = 1$, the load terms $\mathbf{f}$ and g and the Dirichlet boundary conditions are chosen in accordance with the analytical solution

$$\begin{aligned} \mathbf{u}_{\text{ex}}(x, y) &= e^x \begin{pmatrix} \sin y + y \cos y - x \sin y \\ -x \cos y - y \sin y - \cos y \end{pmatrix}, \\ p_{\text{ex}}(x, y) &= \sin(\pi x) \cos(4\pi y), \\ \vartheta_{\text{ex}}(x, y) &= \sin(4\pi x) \sin(\pi y). \end{aligned}$$

Notice that the viscosity ν satisfies Property (P0) whenever $|\vartheta| \leq \vartheta_{\text{max}}$. The domain Ω is partitioned with the following sequences of polygonal meshes: QUADRILATERAL distorted meshes, TRIANGULAR meshes, CVT (Centroidal Voronoi Tessellations) meshes, RANDOM Voronoi meshes (see Fig. 1). For the generation of the Voronoi meshes we used the code `Polymesher` [43]. For each family of meshes we take the sequence with diameter $h = 2^{-2}, 2^{-3}, 2^{-4}, 2^{-5}, 2^{-6}$.

In Table 1 we show the dimensions of the linear systems solved at each iteration of the fixed-point procedure and number of iterations needed for the adopted sequences of meshes. We iterate the fixed-point procedure with the tolerance $\text{TOLL} = 1\text{e-}7$.

In Fig. 2 we display the errors $\text{err}(\mathbf{u}_h, H^1)$, $\text{err}(\mathbf{u}_h, L^2)$, $\text{err}(\vartheta_h, H^1)$, $\text{err}(\vartheta_h, L^2)$ and $\text{err}(p_h, L^2)$ for the sequences of meshes aforementioned. We notice that the theoretical predictions of Proposition 5 are confirmed, moreover we observe that the method is robust with respect to the mesh distortion.

QUADRILATERAL MESHES			
1/h	HEAT	OSEEN	N_IT
4	49	147	6
8	225	643	6
16	961	2691	6
32	3969	11011	5
64	16129	44547	6

VORONOI MESHES			
1/h	HEAT	OSEEN	N_IT
4	63	175	6
8	315	823	6
16	1399	3567	6
32	5847	14767	5
64	23877	60043	6

TRIANGULAR MESHES			
1/h	HEAT	OSEEN	N_IT
4	129	403	6
8	605	1847	6
16	2547	7705	5
32	10331	31121	5
64	41985	126211	6

RANDOM MESHES			
1/h	HEAT	OSEEN	N_IT
4	69	187	6
8	307	807	6
16	1367	3503	6
32	5659	14391	5
64	22921	58131	6

TABLE 1 *Test 1. Dimensions of the linear systems solved at each iteration of the fixed-point procedure and number of iterations.*

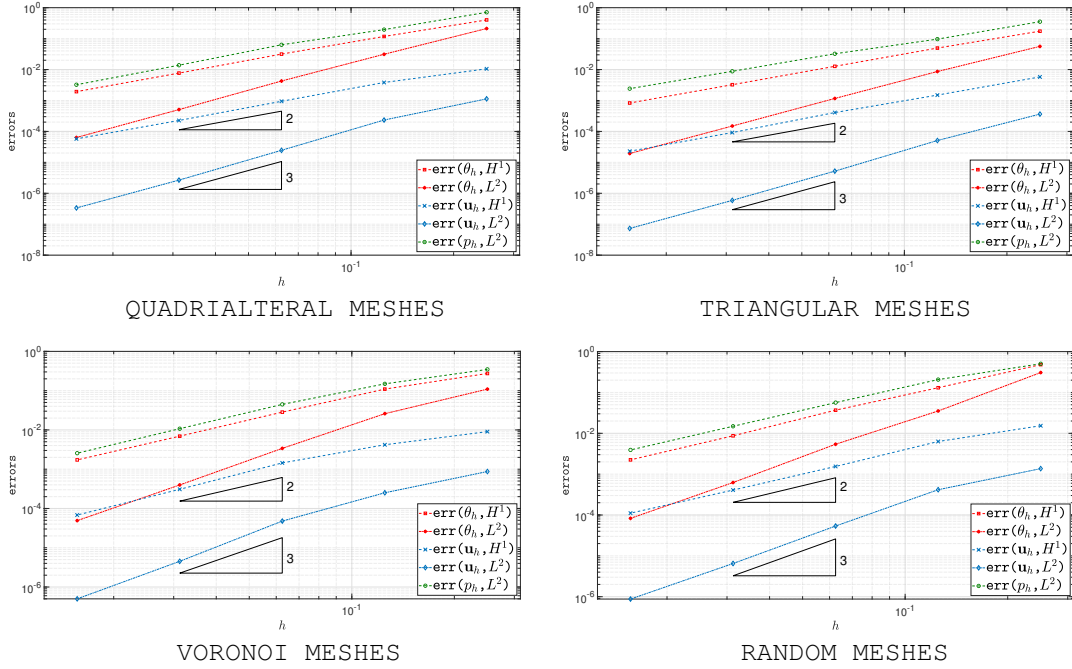


Figure 2. Test 1. Convergence histories of the VEM (3.27) for the adopted families of polygonal meshes.

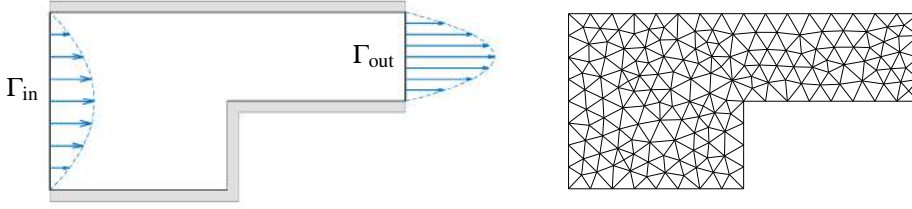


Figure 3. Test 2. Domain and boundary conditions description (left) and mesh with diameter $h = 2^{-2}$ (right).

Test 2. Convection-dominated transport of a passive scalar. The scope of the present test is to show that divergence-free discrete velocity solutions might also be of advantage in coupled problems. We investigate the behavior of the proposed method for the test proposed in [38, Example 6.6]. We consider the following coupled problem on the domain $\Omega = (0, 4) \times (0, 2) \setminus [2, 4] \times [0, 1]$ with a flow field which is governed by the Stokes equation

$$\begin{cases} -\nu \operatorname{div}(\boldsymbol{\varepsilon}(\mathbf{u})) - \nabla p = \mathbf{0} & \text{in } \Omega, \\ \operatorname{div} \mathbf{u} = 0 & \text{in } \Omega, \\ -\kappa \Delta \vartheta + \mathbf{u} \cdot \nabla \vartheta = 0 & \text{in } \Omega, \end{cases} \quad (6.6)$$

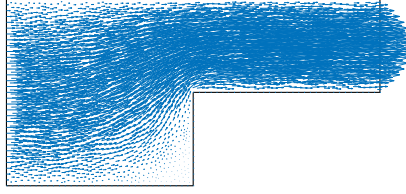
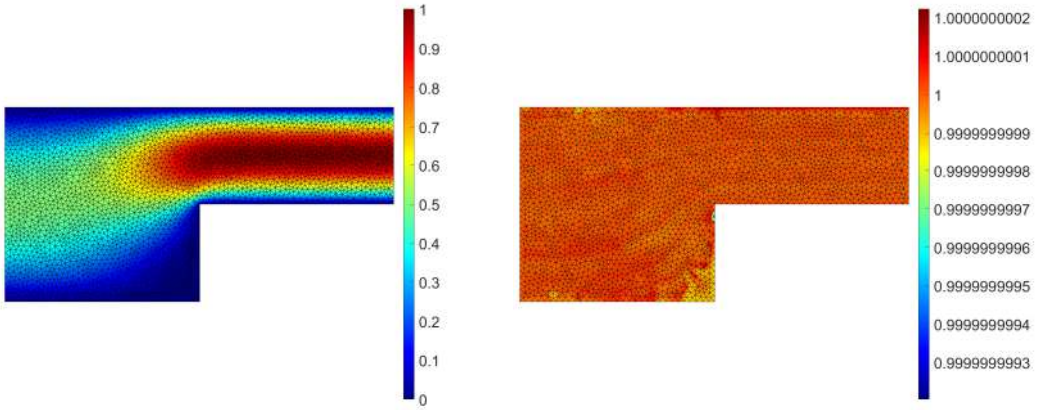
coupled with the following boundary conditions (see Fig. 3 left)

$$\begin{cases} \mathbf{u} = (0.5y(2-y), 0)^T & \text{on } \Gamma_{in}, \\ \mathbf{u} = (4(y-1)(2-y), 0)^T & \text{on } \Gamma_{out}, \\ \mathbf{u} = \mathbf{0} & \text{on } \partial\Omega \setminus (\Gamma_{in} \cup \Gamma_{out}), \end{cases} \quad \begin{cases} \vartheta = 1 & \text{on } \Gamma_{in}, \\ \kappa \nabla \vartheta \cdot \mathbf{n} = 0 & \text{on } \partial\Omega \setminus \Gamma_{in}, \end{cases} \quad (6.7)$$

the viscosity is $\nu = 1e-2$, the conductivity is $\kappa = 1e-6$. The exact temperature solution is $\vartheta_{ex} = 1$.

For the sake of the comparison with the results reported in [38, Example 6.6], the domain Ω is partitioned with a sequence of triangular meshes with diameter $h = 2^{-1}, 2^{-2}, 2^{-3}, 2^{-4}, 2^{-5}$ (see Fig. 3 right).

The continuous convective form associated with the spaces satisfying the boundary conditions (6.7) is not skew-symmetric, therefore in the discrete scheme we consider the convective form (3.23) in the place of (3.25). Note that the heat equation for the temperature in (6.6) is convection-dominated. However, it is worth stressing that, differently from [38], we do not employ any stabilization technique to stabilize the discrete problem (we refer to [11, 20, 22] for the analysis of the stabilized method in the VEM context both for the convection-dominated elliptic equations and Oseen equations). Fig. 4, Fig. 5 and Table 2 present the results of the numerical simulations. As exhibited in [38, Example 6.6] the violation of the divergence constraint, as it appears in the classical mixed finite element, pollutes the discrete temperature solution with strong spurious oscillations. By contrast, the proposed VEM scheme with divergence-free velocity solution computes the temperature with high accuracy as depicted in Fig. 5 and Table 2.

Figure 4. Test 2. Quiver of the numerical velocity $\mathbf{u}_h$ with $h = 2^{-4}$.Figure 5. Test 2. Absolute value of the velocity $|\mathbf{u}_h|$ (left) and temperature ϑ_h (right) with $h = 2^{-4}$.

$1/h$	2	4	8	16	32
$\min(\vartheta_h - 1)$	$-7.88\text{e-}11$	$-2.78\text{e-}10$	$-2.87\text{e-}08$	$-4.47\text{e-}09$	$-3.23\text{e-}09$
$\max(\vartheta_h - 1)$	$+4.50\text{e-}11$	$+7.51\text{e-}11$	$+2.42\text{e-}08$	$+2.90\text{e-}09$	$+5.34\text{e-}09$

TABLE 2. Test 2. Minimal and maximal of $\vartheta - 1$ (DOFs values).

7. Conclusions

We presented a Virtual Element discretization of a thermo-fluid dynamic problem modeling the stationary flow of a viscous incompressible fluid (governed by a Navier-Stokes equation), where the viscosity of the fluid depends on the temperature (governed by an elliptic equation). We have shown that the discrete problem is well-posed and prove optimal error estimates. Numerical experiments which confirm the theoretical error bounds are also presented. Further developments may include the introduction of uncertainty in the thermal conductivity κ , based on the ideas of [25, 26] as well as the three-dimensional extension, based on employing the three-dimensional VEM for Navier-Stokes proposed in [14].

Funding

P.F.A. and M.V. have been partially funded by the research grants PRIN2017 n. 201744KLJL “*Virtual Element Methods: Analysis and Applications*” and PRIN2020 n. 20204LN5N5 “*Advanced polyhedral discretisations of heterogeneous PDEs for multiphysics problems*”, funded by the Italian Ministry of Universities and Research (MUR). P.F.A., G.V. and M.V. are members of INDAM-GNCS.

REFERENCES

1. R. A. Adams. *Sobolev spaces*, volume 65 of *Pure and Applied Mathematics*. Academic Press, New York-London, 1975.
2. J. Aghili and D. A. Di Pietro. An Advection-Robust Hybrid High-Order Method for the Oseen Problem. *J. Sci. Comput.*, 77(3):1310–1338, 2018.
3. R. Agroum, S. M. Aouadi, C. Bernardi, and J. Satouri. Spectral discretization of the Navier-Stokes equation coupled with the heat equation. *ESAIM Math. Model. Numer. Anal.*, 49(3):621–639, 2015.
4. R. Agroum, C. Bernardi, and J. Satouri. Spectral discretization of the time-dependent Navier-Stokes problem coupled with the heat equation. *Appl. Math. Comput.*, 268(1):59–82, 2015.
5. B. Ahmad, A. Alsaedi, F. Brezzi, L. D. Marini, and A. Russo. Equivalent projectors for virtual element methods. *Comput. Math. Appl.*, 66(3):376–391, 2013.
6. R. Aldbaissy, F. Hecht, G. Mansour, and T. Sayah. A full discretisation of the time-dependent Boussinesq (buoyancy) model with nonlinear viscosity. *Calcolo*, 55(4):Article Number 44, 2018.
7. P. Antonietti, L. Mascotto, M. Verani, and S. Zonca. Stability analysis of polytopic Discontinuous Galerkin approximations of the Stokes problem with applications to fluid-structure interaction problems. *J. Sci. Comput.*, 90(1):Article number 23, 2022.
8. P. Antonietti, M. Verani, C. Vergara, and S. Zonca. Numerical solution of fluid-structure interaction problems by means of a high order Discontinuous Galerkin method on polygonal grids. *Finite Elements in Analysis and Design*, 159:1 – 14, 2019.
9. P. F. Antonietti, L. Beirão da Veiga, D. Mora, and M. Verani. A stream virtual element formulation of the Stokes problem on polygonal meshes. *SIAM J. Numer. Anal.*, 52(1):386–404, 2014.
10. P. F. Antonietti, L. Beirão da Veiga, and G. Manzini. *The Virtual Element Method and its Applications*. Springer International Publishing, 2022.
11. L. Beirão da Veiga, F. Dassi, C. Lovadina, and G. Vacca. SUPG-stabilized Virtual Elements for diffusion-convection problems: a robustness analysis. *ESAIM Math. Model. Numer. Anal.*, 55:2233–2558, 2021.
12. L. Beirão da Veiga, F. Dassi, and A. Russo. High-order virtual element method on polyhedral meshes. *Comput. Math. Appl.*, 74(5):1110–1122, 2017.
13. L. Beirão da Veiga, F. Dassi, and G. Vacca. The Stokes complex for Virtual Elements in three dimensions. *Math. Models Methods Appl. Sci.*, 30(3):477–512, 2018.
14. L. Beirão da Veiga, F. Dassi, and G. Vacca. The Stokes complex for virtual elements in three dimensions. *Math. Models Methods Appl. Sci.*, 30(3):477–512, 2020.
15. L. Beirão da Veiga, C. Lovadina, and G. Vacca. Divergence free virtual elements for the Stokes problem on polygonal meshes. *ESAIM Math. Model. Numer. Anal.*, 51(2):509–535, 2017.
16. L. Beirão da Veiga, C. Lovadina, and G. Vacca. Virtual elements for the Navier-Stokes problem on polygonal meshes. *SIAM J. Numer. Anal.*, 56(3):1210–1242, 2018.
17. L. Beirão da Veiga and G. Vacca. Sharper error estimates for virtual elements and a bubble-enriched version. *SIAM J. Numer. Anal.*, 60(4):1853–1878, 2022.
18. L. Beirão da Veiga, F. Brezzi, A. Cangiani, G. Manzini, L. D. Marini, and A. Russo. Basic principles of virtual element methods. *Math. Models Methods Appl. Sci.*, 23(1):199–214, 2013.
19. L. Beirão da Veiga, F. Brezzi, L. D. Marini, and A. Russo. The Hitchhiker’s Guide to the Virtual Element Method. *Math. Models Methods Appl. Sci.*, 24(8):1541–1573, 2014.

20. L. Beirão da Veiga, F. Dassi, and G. Vacca. Vorticity-stabilized Virtual Elements for the Oseen Equation. *Math. Models Methods Appl. Sci.*, 31(14):3009–3052, 2021.
21. L. Beirão da Veiga, C. Lovadina, and A. Russo. Stability analysis for the virtual element method. *Math. Mod.and Meth. in Appl. Sci.*, 27(13):2557–2594, 2017.
22. M. F. Benedetto, S. Berrone, A. Borio, S. Pieraccini, and S. Scialò. Order preserving SUPG stabilization for the virtual element formulation of advection–diffusion problems. *Comput. Methods Appl. Mech. Engrg.*, 311:18–40, 2016.
23. C. Bernardi, S. Dib, V. Girault, F. Hecht, F. Murat, and T. Sayah. Finite element methods for Darcy’s problem coupled with the heat equation. *Numer. Math.*, 139(2):315–348, 2018.
24. D. Boffi, F. Brezzi, and M. Fortin. *Mixed finite element methods and applications*, volume 44 of *Springer Series in Computational Mathematics*. Springer, Heidelberg, 2013.
25. F. Bonizzoni and F. Nobile. Perturbation analysis for the Darcy problem with log-normal permeability. *SIAM/ASA J. Uncertain. Quantif.*, 2(1):223–244, 2014.
26. F. Bonizzoni and F. Nobile. Regularity and sparse approximation of the recursive first moment equations for the lognormal Darcy problem. *Comput. Math. Appl.*, 80(12):2925–2947, 2020.
27. L. Botti, D. A. Di Pietro, and J. Droniou. A Hybrid High-Order discretisation of the Brinkman problem robust in the Darcy and Stokes limits. *Comput. Methods Appl. Mech. Engrg.*, 341:278–310, 2018.
28. S. C. Brenner, Q. Guan, and L. Y. Sung. Some estimates for virtual element methods. *Comput. Methods Appl. Math.*, 17(4):553–574, 2017.
29. S. C. Brenner and L. R. Scott. *The Mathematical Theory of Finite Element Methods*, volume 15 of *Texts in Applied Mathematics*. Springer, New York, third edition, 2008.
30. S. C. Brenner and L. Y. Sung. Virtual element methods on meshes with small edges or faces. *Math. Models Methods Appl. Sci.*, 28(7):1291–1336, 2018.
31. A. Cangiani, E.H. Georgoulis, T. Pryer, and O. Sutton. A posteriori error estimates for the virtual element method. *Numer. Math.*, 137(4):857–893, 2017.
32. D. Castañón Quiroz and D. A. Di Pietro. A Hybrid High-Order method for the incompressible Navier–Stokes problem robust for large irrotational body forces. *Comput. Math. Appl.*, 79(9):2655–2677, 2020.
33. B. Cockburn, G. Fu, and W. Qiu. A note on the devising of superconvergent HDG methods for Stokes flow by M -decompositions. *IMA J. Numer. Anal.*, 37(2):730–749, 2017.
34. D. A. Di Pietro and S. Krell. A hybrid high-order method for the steady incompressible Navier-Stokes problem. *J. Sci. Comput.*, 74(3):1677–1705, 2018.
35. J. K. Djoko, V. S. Konlack, and M. Mbheou. Stokes equations under nonlinear slip boundary conditions coupled with the heat equation: A priori error analysis. *Numer. Methods Partial Differential Equations*, 36(1):86–117, 2020.
36. N. Gauger, A. Linke, and P. Schroeder. On high-order pressure-robust space discretisations, their advantages for incompressible high Reynolds number generalised Beltrami flows and beyond. *SMAI J. Comput. Math.*, 5:88–129, 2019.
37. V. Girault and P.-A. Raviart. *Finite element methods for Navier-Stokes equations*, volume 5 of *Springer Series in Computational Mathematics*. Springer-Verlag, Berlin, 1986. Theory and algorithms.
38. V. John, A. Linke, C. Merdon, M. Neilan, and L. G. Rebholz. On the divergence constraint in mixed finite element methods for incompressible flows. *SIAM Rev.*, 59(3):492–544, 2017.
39. K. Lipnikov, D. Vassilev, and I. Yotov. Discontinuous Galerkin and mimetic finite difference methods for coupled Stokes-Darcy flows on polygonal and polyhedral grids. *Numer. Math.*, 126(2):321–360, 2014.
40. S. Natarajan. On the application of polygonal finite element method for Stokes flow – A comparison between equal order and different order approximation. *Comput. Aided Geom. Design*, 77:Article number:101813, 2020.
41. R. Oyarzúa, T. Qin, and D. Schötzau. An exactly divergence-free finite element method for a generalized Boussinesq problem. *IMA J. Numer. Anal.*, 34(3):1104–1135, 2014.
42. A. Quarteroni and A. Valli. *Numerical approximation of partial differential equations*, volume 23 of *Springer Series in Computational Mathematics*. Springer-Verlag, Berlin, 1994.

43. C. Talischi, G. H. Paulino, A. Pereira, and I. F.M . Menezes. Polymesher: a general-purpose mesh generator for polygonal elements written in matlab. *Struct. Multidisc Optimiz.*, 45(3):309–328, 2012.
44. G. Vacca. An H^1 -conforming virtual element for Darcy and Brinkman equations. *Math. Models Methods Appl. Sci.*, 28(1):159–194, 2018.